

Nonlinear Physics: Theory and Experiment IV

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Recurrence in the KdV Equation?

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Korteweg-de Vries equation (KdV)

$$u_t + uu_x + \delta^2 u_{xxx} = 0$$

a particular continuum limit of an FPU lattice

$$\frac{d^2}{dt^2} y = (y_{n-1} - 2y_n + y_{n+1}) [1 - \alpha (y_{n-1} - y_{n+1})]$$

Repeated near-recurrences are observed in FPU lattices with

$$y_n(0) = \sin\left(\frac{n\pi}{N}\right), \quad \frac{d}{dt} y_n(0) = 0.$$

[FPU, 1955]

Look for (near) recurrence in KDV by simulation
[Zabusky & Kruskal,1965]

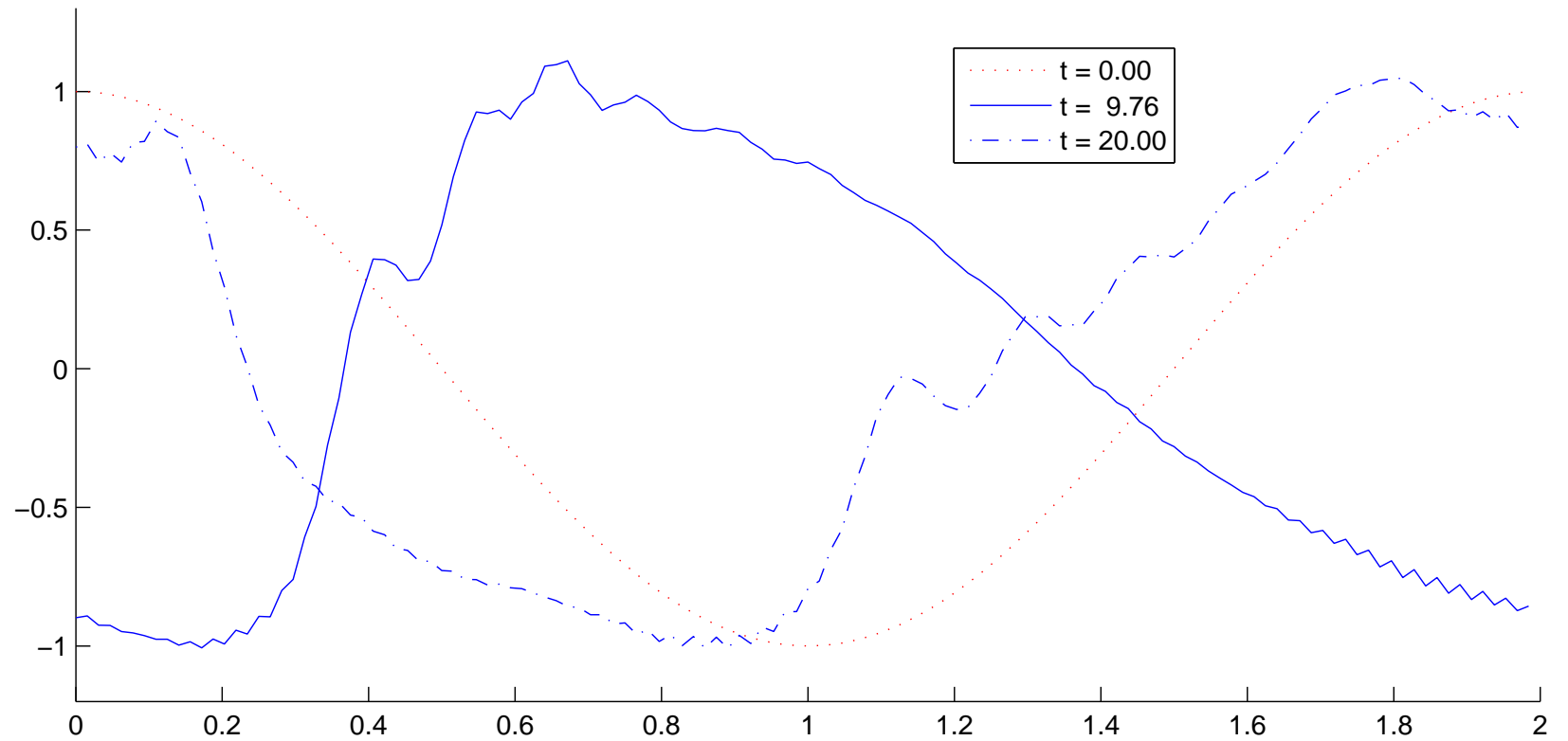
$$u_t + uu_x + \delta^2 u_{xxx} = 0$$

$$\delta = .022^2 \ll 1 \quad u(x, 0) = \sin(\pi x)$$

$$\frac{d}{dt}u_n = \frac{(u_{n-1} + u_n + u_{n+1})(u_{n-1} - u_{n+1})}{3h} + \delta^2 \frac{u_{n-2} - 2u_{n-1} + 2u_{n+1} - u_{n+2}}{2h^3}$$

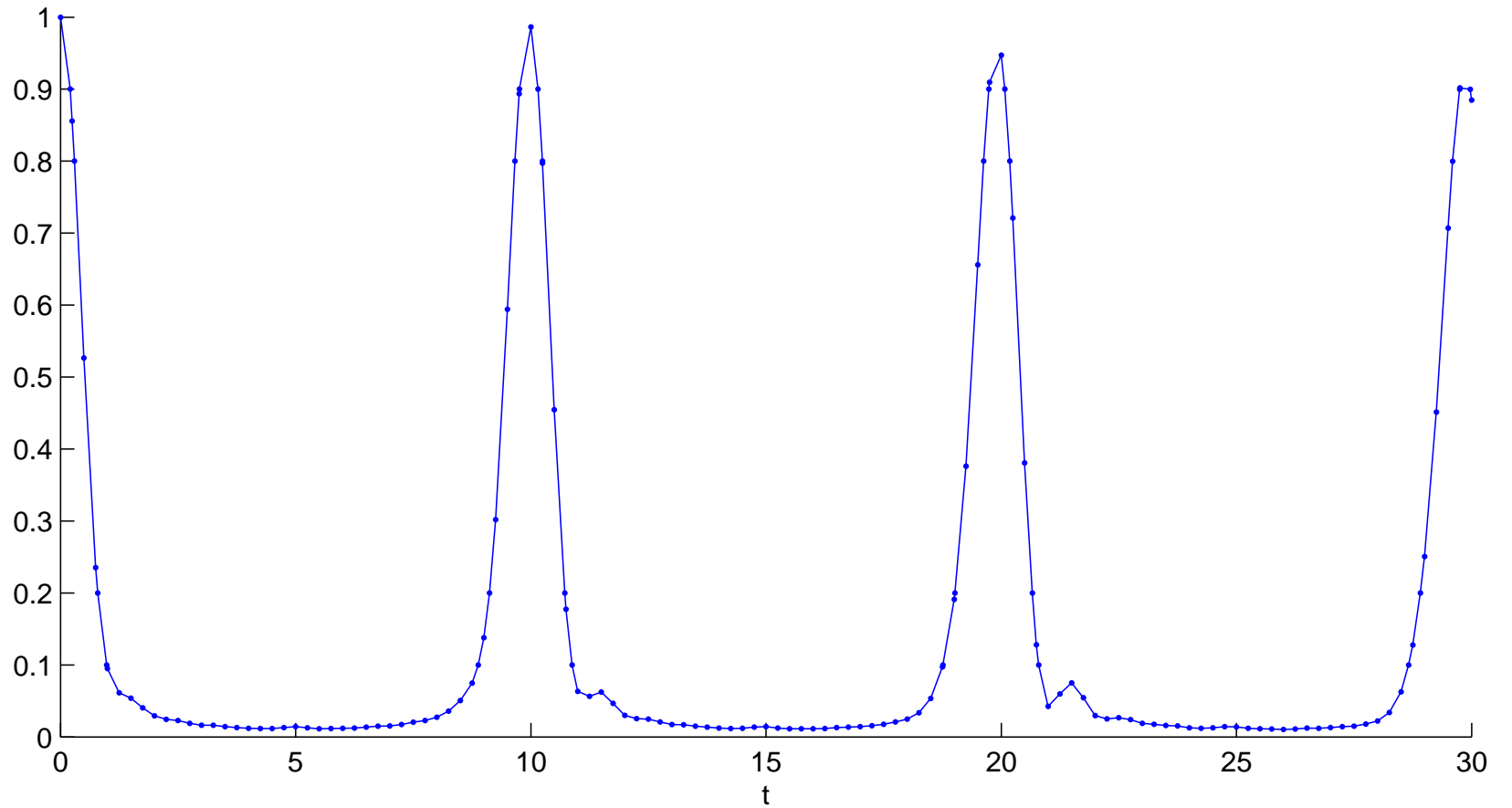
$$\frac{d}{dt}\mathbf{u} = -(\mathbf{A}_3\mathbf{u})(\mathbf{D}_c\mathbf{u}) - \delta^2\mathbf{D}_{3c}\mathbf{u}$$

Recurrence



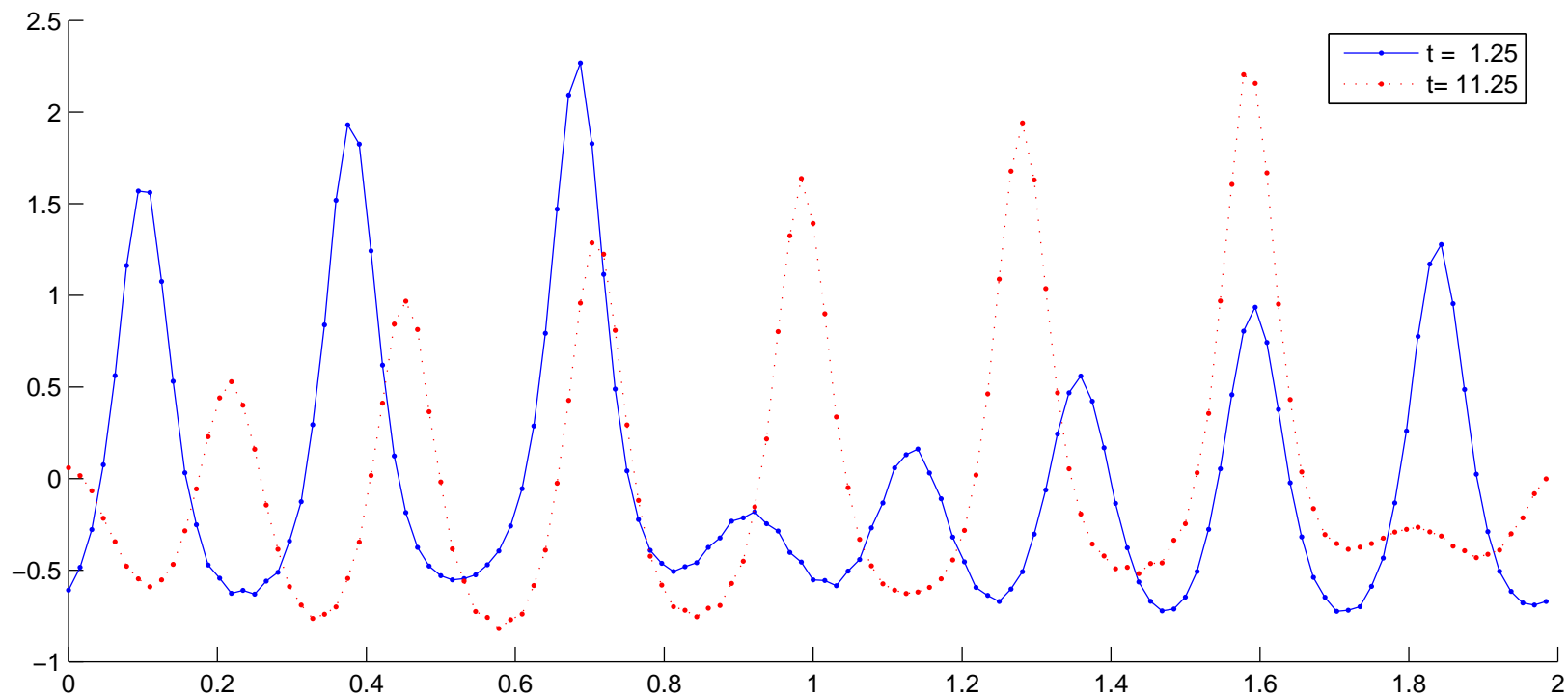
Simulation with $N = 128$

Recurrence



Simulation with $N = 128$

Emergence of Solitons



Simulation with $N = 128$

Emergence of Solitons

KdV solitary waves

$$u(x, t) = 12\delta^2 k^2 \operatorname{sech}^2 [k(x - 4\delta^2 k^2 t)]$$

are the limit of periodic, cnoidal-wave solutions

$$u(x, t) = 12\delta^2 k^2 m \operatorname{cn}^2 [k(x - 4\delta^2 k^2 (2m - 1)t, m)]$$

The solitary waves are *solitons*:

they pass through one another and regain their original shape & velocity.

Recurrence

Recurrence has been explained in terms of the solitons.

[Zabusky & Kruskal, 1965; Osborne & Bergamasco, 1986]

Q: How does recurrence depend on the *spatial* discretization?

Hypotheses:

H1: Higher-accuracy discretization gives better recurrence.

H2: Integrable discretization gives better recurrence.

High-accuracy Discretizations

Pseudo-spectral:

$$\frac{d}{dt}\mathbf{u} = -\mathbf{u}(\mathbf{D}_f\mathbf{u}) - \delta^2\mathbf{D}_f^3\mathbf{u}$$

Pseudo-spectral (Conservation Form):

$$\frac{d}{dt}\mathbf{u} = -\frac{1}{2}\mathbf{D}_f\mathbf{u}^2 - \delta^2\mathbf{D}_f^3\mathbf{u}$$

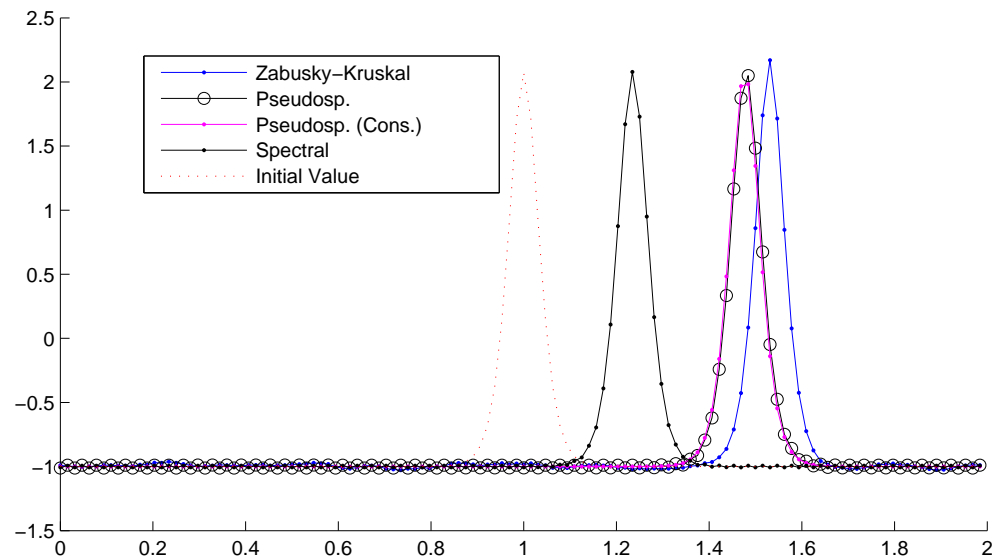
Spectral:

$$\frac{d}{dt}\mathbf{u} = -\mathbf{D}_f\mathbf{F}^{-1}[(\mathbf{F}\mathbf{u}) * (\mathbf{F}\mathbf{u})] - \delta^2\mathbf{D}_f^3\mathbf{u}$$

$\mathbf{D}_f = \mathbf{F}^{-1}\mathbf{\Omega}\mathbf{F}$: Fourier differentiation matrix

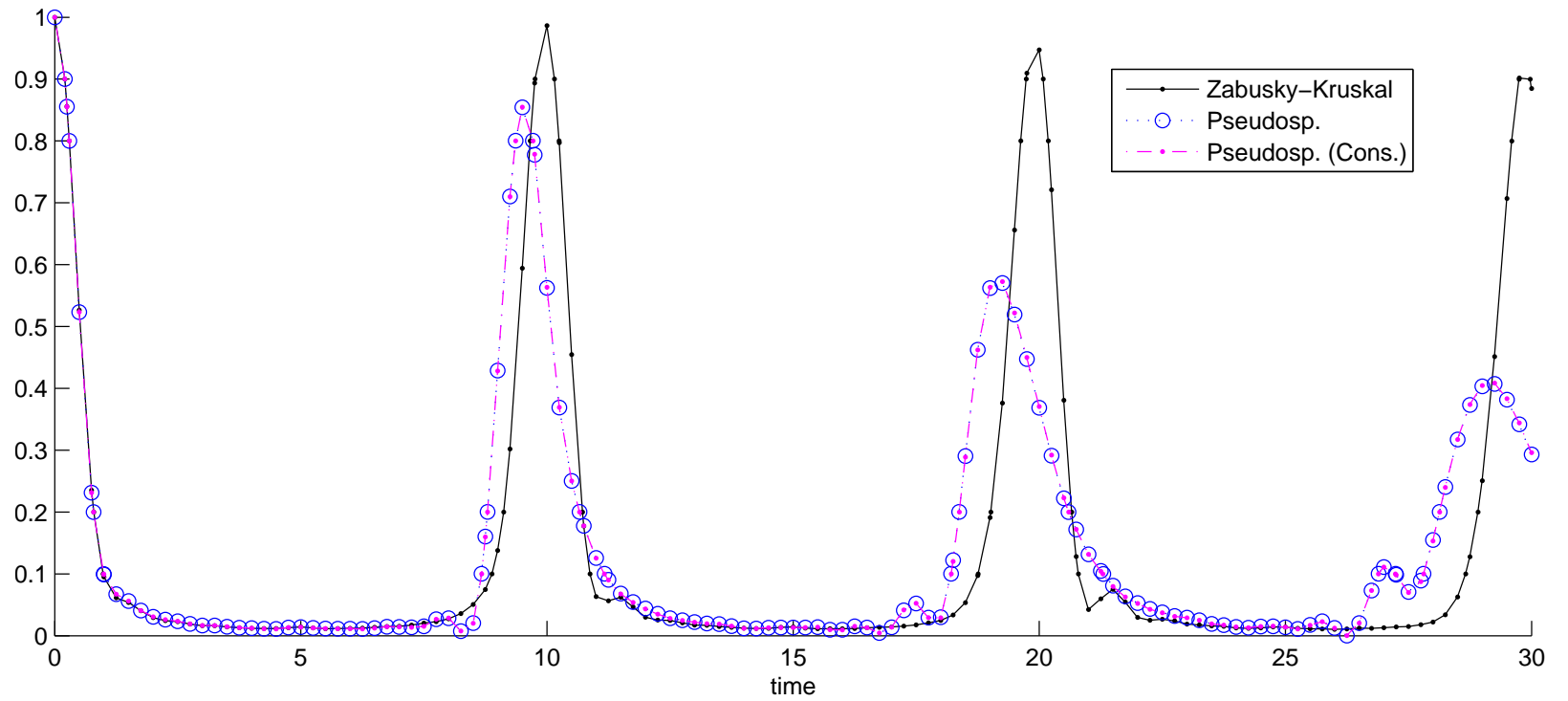
\mathbf{F} : discrete Fourier transform matrix

Solitons



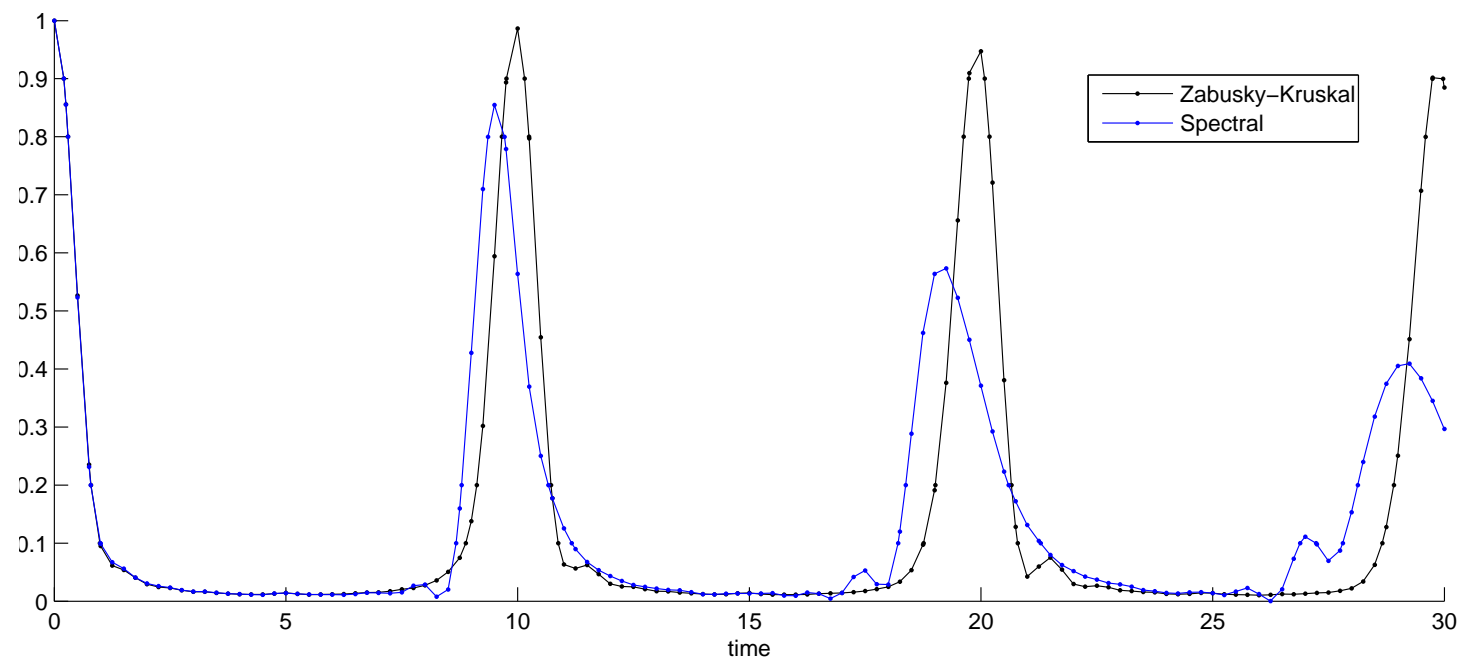
Simulation with $N = 128$ plotted at $t = 9.75$

Recurrence



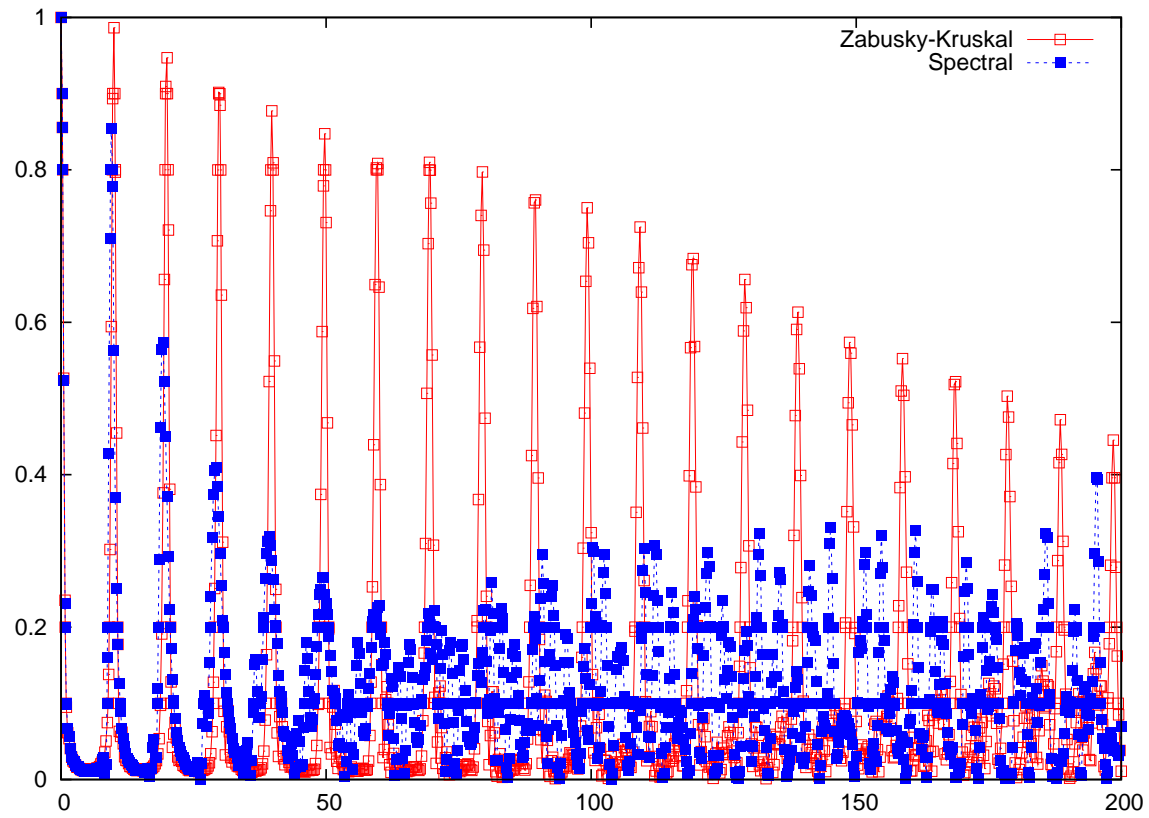
Simulation with $N = 128$

Recurrence



Simulation with $N = 128$

Recurrence



Simulation with $N = 128$

Recurrence

Surprise(?): Recurrence is *worse* in pseudospectral and spectral schemes.

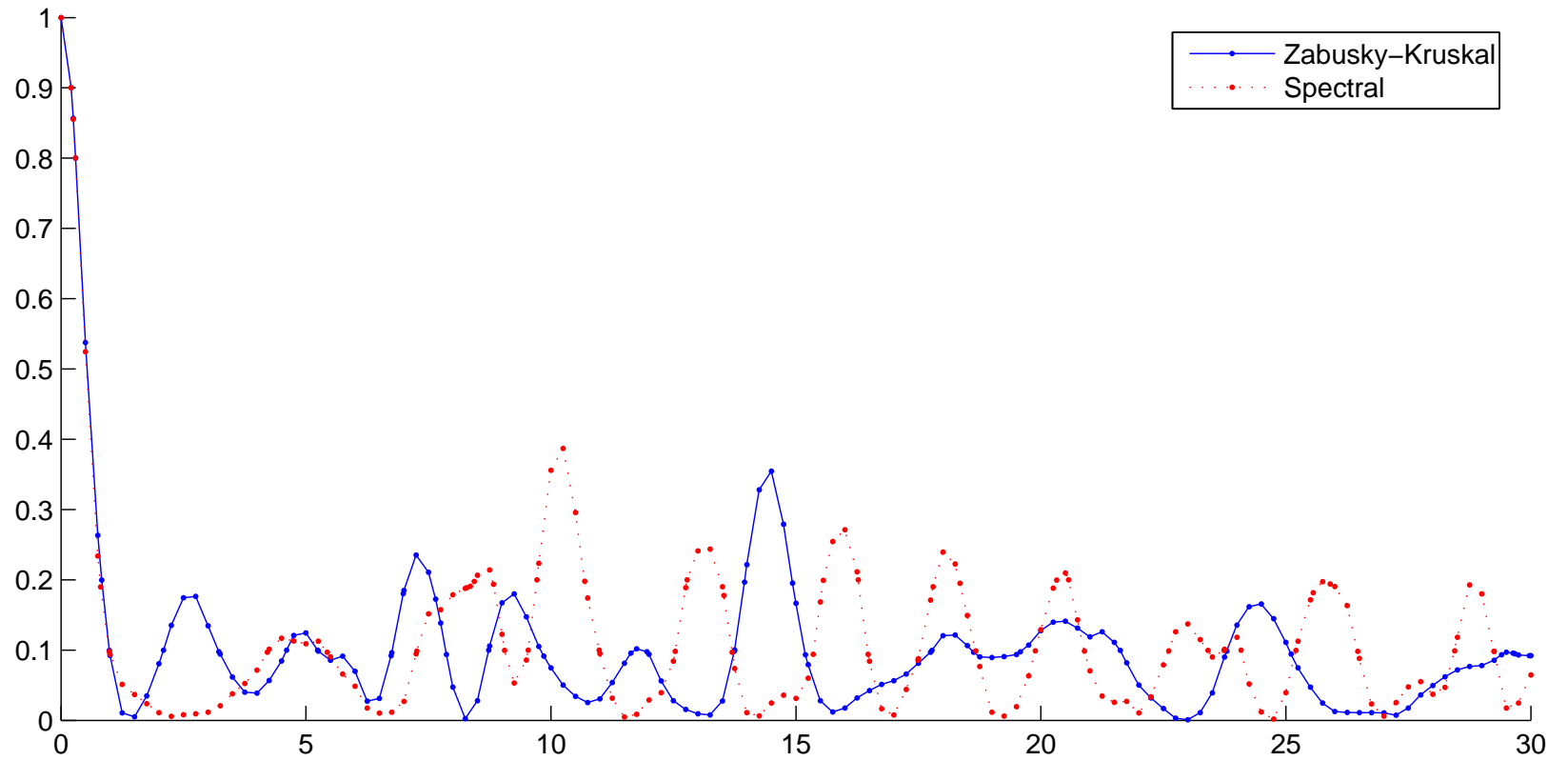
Higher accuracy does *not* yield better recurrence.

Q: What about a “rougher” grid (e.g., $N = 64$)?

A1: Zabusky-Kruskal and Spectral discretizations don't show recurrence.

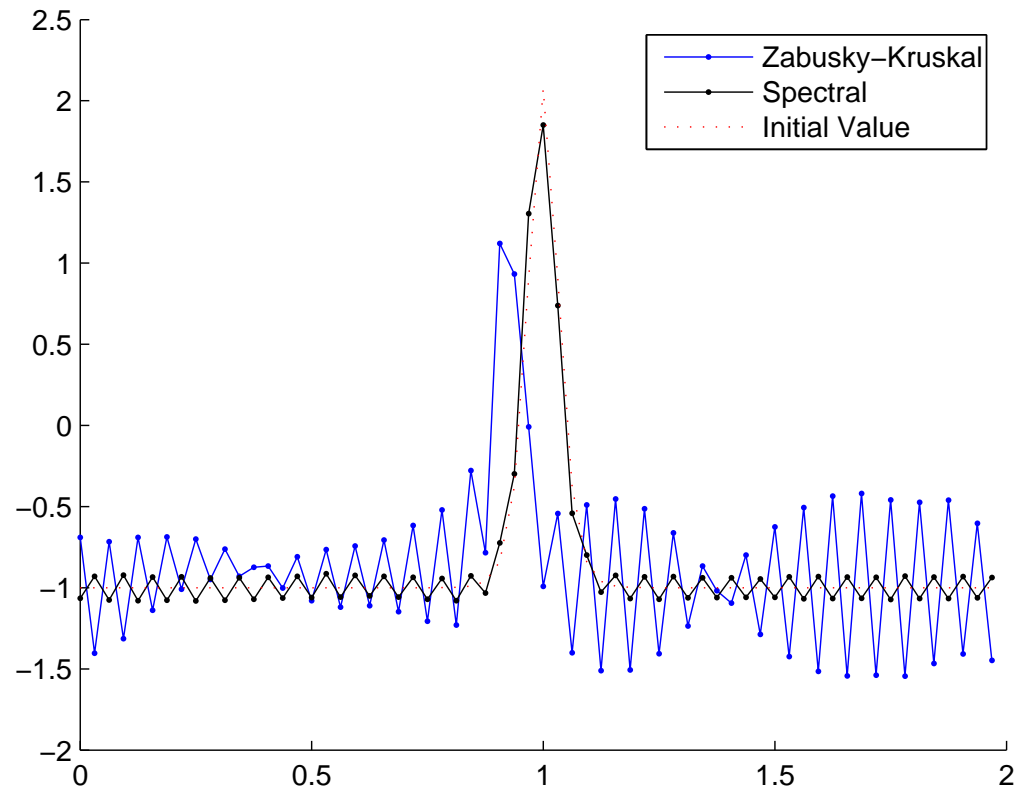
A2: Pseudospectral discretizations manifest a nonlinear instability.

No Recurrence



Simulation with $N = 64$

Solitons(?)



Simulation with $N = 64$ plotted at $t = .75$

Nonlinear Instability

In pseudospectral discretizations there is rapid uncontrolled growth of the solution for “rough” grids.

The nonlinear terms induce aliasing.

Preservation of $\sum_n u_n^2$ in other discretizations precludes the instability.

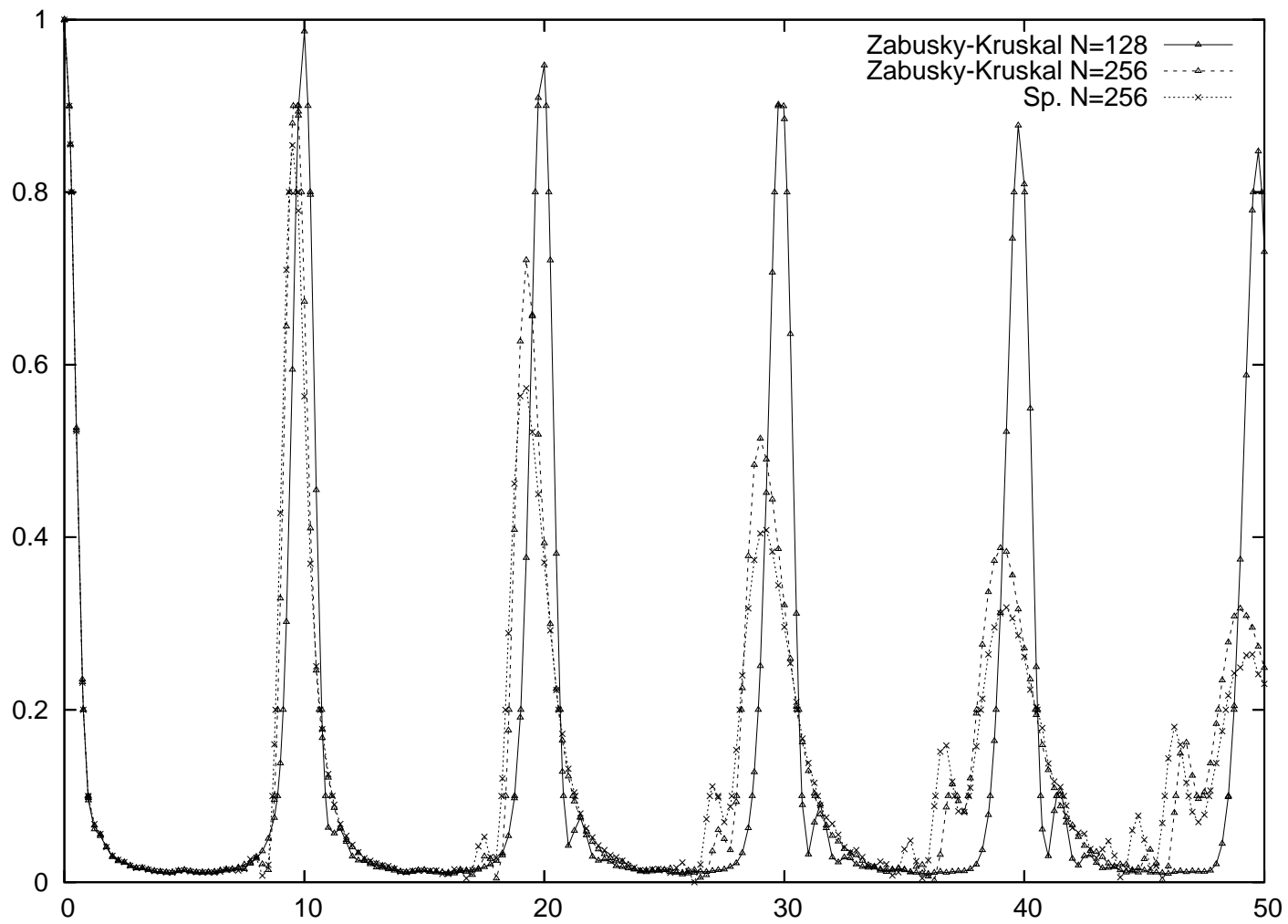
A similar instability exists in:

simple finite-difference discretizations of KdV,

discretizations of viscous Burger’s equation [Maritz & Schoombie],

discretizations inviscid Burger’s equation. [Majda & Timofeyev, 2002]

Finer Grid



Integrable Discretization of KdV

KdV can be associated with the Zakharov-Shabat Scattering problem in the form

$$\psi_x = \begin{pmatrix} ik & u \\ -1 & -ik \end{pmatrix} \psi$$

Forward difference \Rightarrow discrete (Ablowitz-Ladik) scattering problem:

$$\psi_{n+1} = \begin{pmatrix} z & U_n \\ \alpha & z^{-1} \end{pmatrix} \psi_n = \mathbf{S}_n \psi_n$$

where

$$U_n = hu_n, \quad \alpha = -h, \quad z = e^{ikh}$$

Integrable Discretization of KdV

Discrete Compatibility Condition:

$$\frac{d}{d\tau} \mathbf{S}_n = \mathbf{T}_{n+1} \mathbf{S}_n - \mathbf{S}_n \mathbf{T}_n$$

where

$$\psi_{n+1} = \mathbf{S}_n \psi_n \quad \frac{d}{d\tau} \psi_n = \mathbf{T}_n \psi$$

Compatibility condition is equivalent to:

$$\begin{aligned} \frac{d}{d\tau} U_n = & (1 - \alpha U_n) [-\alpha U_{n-1} (U_{n-2} - U_n) - \alpha U_{n+1} (U_n - U_{n+2}) \\ & - \alpha (U_{n-1} + 2U_n + U_{n+1}) (U_{n-1} - U_{n+1}) \\ & + U_{n-2} - 2U_{n-1} + 2U_{n+1} - U_{n+2}] \end{aligned}$$

Integrable Discretization of KdV

Rescale:

$$U_n \rightarrow \frac{h}{6} u_n, \quad \alpha \rightarrow -\frac{h}{\delta^2} \quad \tau \rightarrow \frac{3\delta^2}{h^4}$$

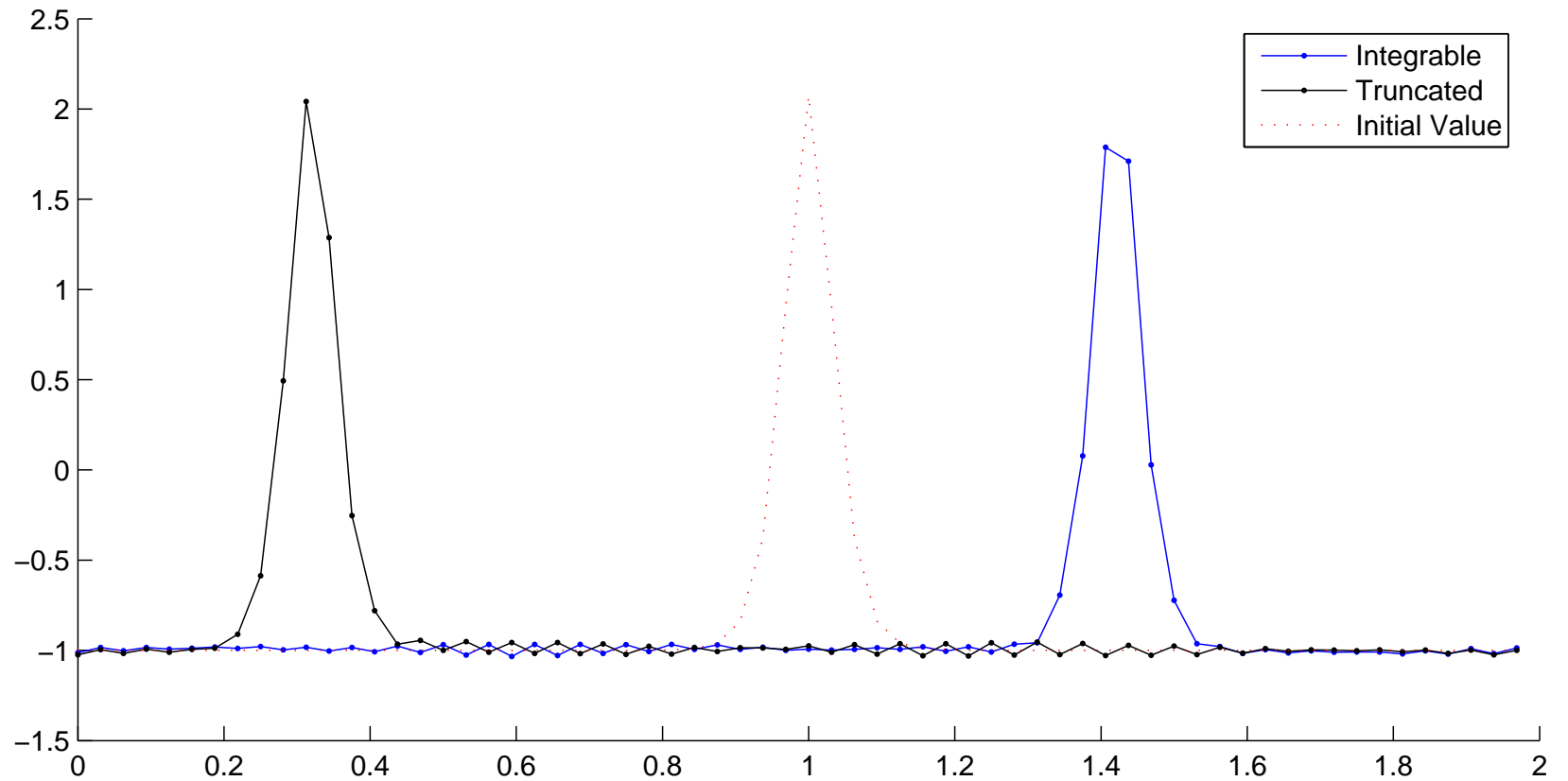
IDKdV:

$$\begin{aligned} \frac{d}{dt} u_n = & \left(1 + \frac{h^2 u_n}{6\delta^2} \right) \left[\frac{u_{n-1}(u_{n-2} - u_n)}{12h} + \frac{u_{n+1}(u_n - u_{n+2})}{12h} \right. \\ & \left. + \frac{(u_{n-1} + 2u_n + u_{n+1})(u_{n-1} - u_{n+1})}{12h} + \delta^2 \frac{u_{n-2} - 2u_{n-1} + 2u_{n+1} - u_{n+2}}{2h^3} \right] \end{aligned}$$

Truncated (non-integrable):

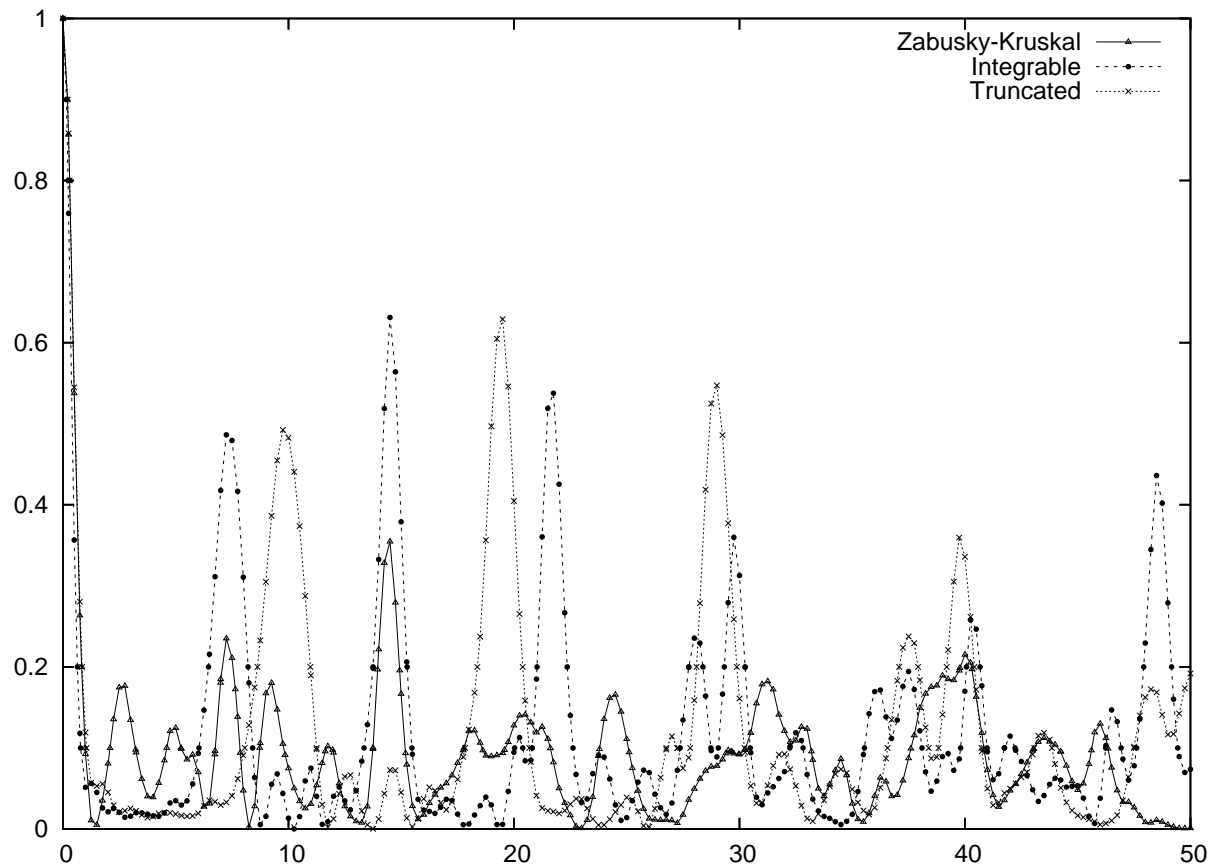
$$\begin{aligned} \frac{d}{dt} u_n = & \frac{u_{n-1}(u_{n-2} - u_n)}{12h} + \frac{u_{n+1}(u_n - u_{n+2})}{12h} \\ & + \frac{(u_{n-1} + 2u_n + u_{n+1})(u_{n-1} - u_{n+1})}{12h} + \delta^2 \frac{u_{n-2} - 2u_{n-1} + 2u_{n+1} - u_{n+2}}{2h^3} \end{aligned}$$

Solitons



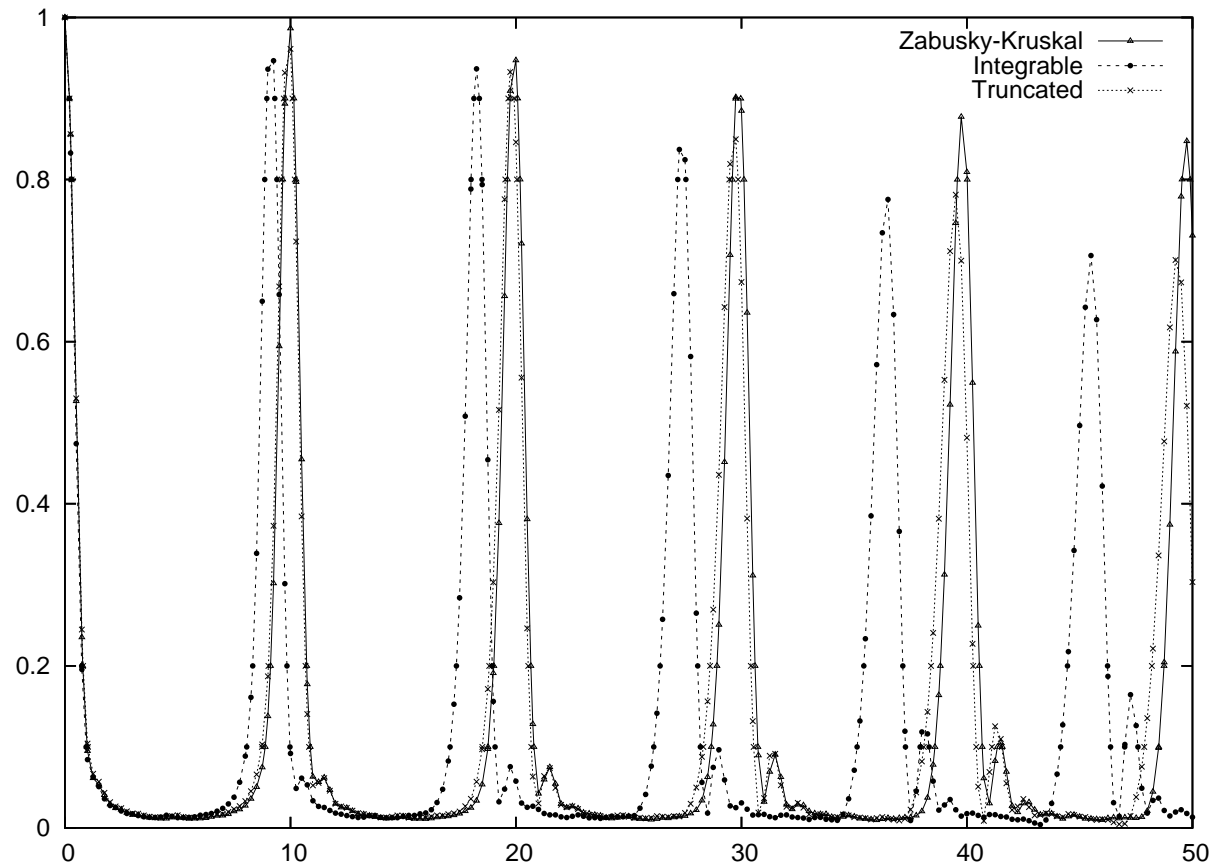
Simulation with $N = 64$ plotted at $t = 39.75$

Recurrence



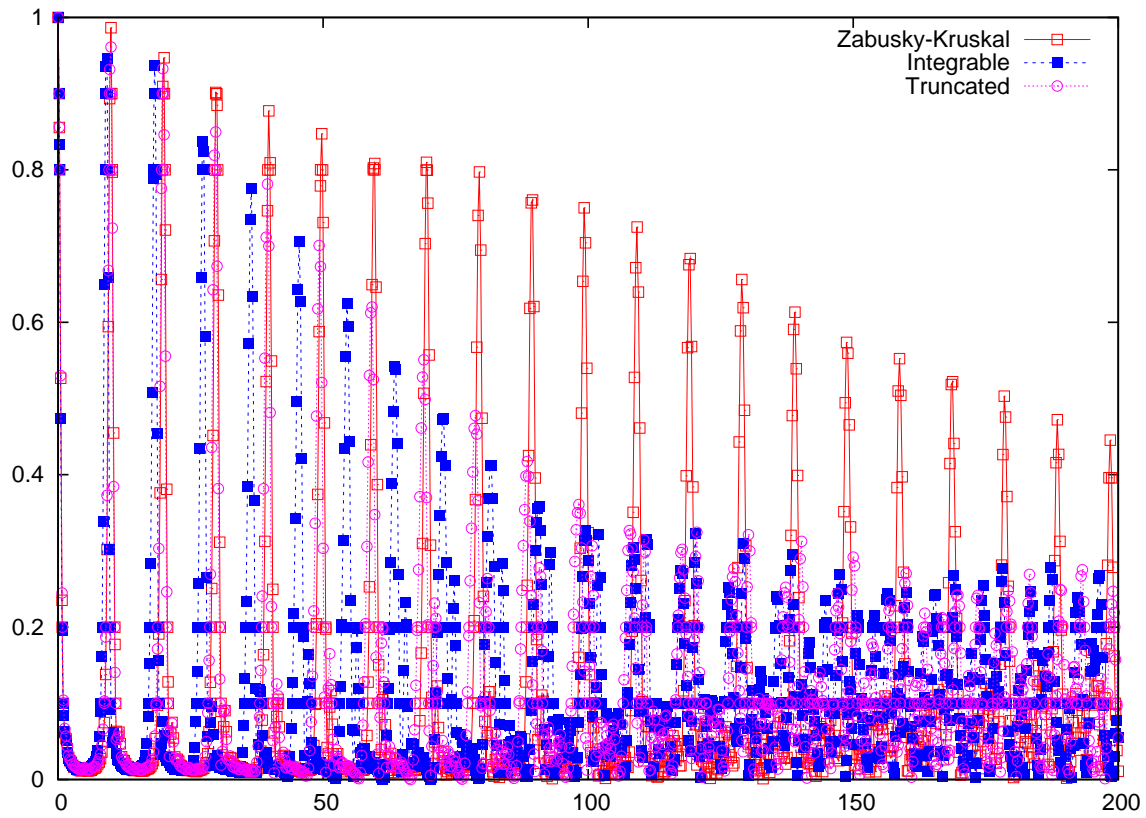
Simulation with $N = 64$ ($\frac{h^2}{6\delta^2} = .336$)

Recurrence



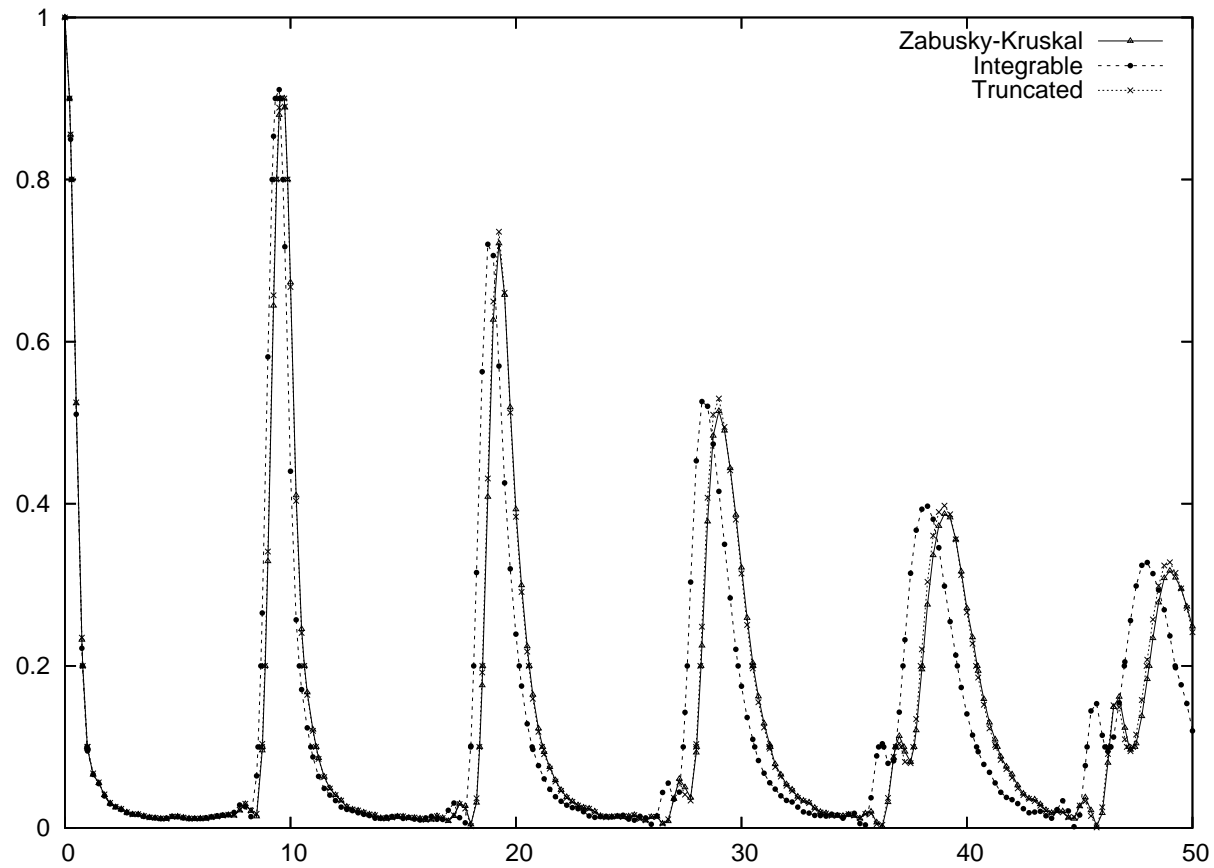
Simulation with $N = 128$ ($\frac{h^2}{6\delta^2} = .0841$)

Recurrence



Simulation with $N = 128$

Recurrence



Simulation with $N = 256$ ($\frac{h^2}{6\delta^2} = .0210$)

Observations

Presence of solitons in discretization not sufficient to capture recurrence.

Recurrence is *worse* in pseudospectral and spectral discretizations than carefully-chosen (Zabusky-Kruskal) finite-differences.

Recurrence not always strengthened by decrease the grid size.

Integrable discretization does not capture recurrence better than Zabusky-Kruskal discretization.

(Related?) Phenomena:

- uncontrolled nonlinear instability
- (modulated) grid-scale oscillations