Reciprocal transformations for a Class of Weakly Nonlinear Hydrodynamic-Type Systems

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$$\boldsymbol{u}_t = K(\boldsymbol{u})\boldsymbol{u}_x,\tag{1}$$

where $\boldsymbol{u} = (u_1, \dots, u_n)^T$, K is an $n \times n$ matrix, called hydrodynamic-type or dispersionless systems.

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The construction of WNSH systems in a coordinate-free fashion ?

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Starting object: (1,1)-tensor L on $M \in \mathbb{R}^n$, of vanishing Nijenhuis torsion and simple eigenvalues. Target: an infinite family of d-Killing systems together with their general solutions.

The construction consists of three steps:

1. From tensor *L* we construct seed d-Killing systems (*d-Killing systems associated with the L-tensor*) together with infinitely many conservation laws numerated by integer $\beta > 0$.

2. Using any *r* conservation laws β_r , $1 \le r \le n$ we define *r*-time reciprocal transformations (RT)and a set of related d-Killing systems.

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3. We show that all constructed systems have common Riemann invariants λ^j , j = 1, ..., n, which are eigenvalues of the tensor *L*. In Riemann invariants we derive the general solution for seed systems and the general solution of related d-Killing systems



Consider: $M \in \mathbb{R}^n$, $u = (u^1, \dots, u^n)^T$, (2,0)-metric G, (1,1)-Killing tensors K_i , $i = 1, \dots, n$ which have a common orthogonal eigenframe of closed one-forms.

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The flows (3) with i = 1, ..., n, commute for any fixed j.



Construction of seed systems from *L*-tensor:

$$K_1 = \mathbb{I}, \quad K_r = \sum_{k=0}^{r-1} \rho_k L^{r-1-k}, \quad r = 2, \dots, n,$$
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where

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(1,1)-tensors K_j are Killing tensors of contravariant metric G (covariant metric $g = G^{-1}$) such that

 $\nabla_m L_j^i = \delta_m^i \partial_j (\operatorname{tr} L) + \sum_{k=1}^n g_{jm} G^{ik} \partial_k (\operatorname{tr} L), \quad m, i, j = 1, \dots, n.$

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where m, β^i and $\gamma^{ij} = \gamma^{ji}$ are arbitrary constants.

For the seed systems we have an infinite set of conservation laws (Blaszak 1999, 2003):

 $D_{t_i}(V_j^{(k)}) = D_{t_j}(V_i^{(k)}), \quad i, j = 1, \dots, n, \quad i \neq j, \quad k \in \mathbb{Z}.$ (5)

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where

$$V_r^{(k)} = V_{r+1}^{(k-1)} - \rho_r V_1^{(k-1)}, k \in \mathbb{Z},$$

with the initial condition

$$V_r^{(0)} = -\delta_r^n, \qquad r = 1, \dots, n.$$



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The numbers s_i , i = 1, ..., k, are a k-tuple of distinct integers from the set $\{1, ..., n\}$, and β_j are positive integers:

$$\beta_1 > \beta_2 > \cdots > \beta_k > n-1.$$

The inverse transformation:

$$dt_{s_i} = -\sum_{j=1}^n \tilde{V}_j^{(n-s_i)} d\tilde{t}_j, \quad i = 1, \dots, k,$$
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$$0 = V_j^{(m)} + \sum_{p=1}^k \tilde{V}_{s_p}^{(m)} V_j^{(\beta_p)}, \ j = s_1, \dots, s_k,$$
$$\tilde{V}_j^{(m)} = V_j^{(m)} + \sum_{p=1}^k \tilde{V}_{s_p}^{(m)} V_j^{(\beta_p)}, \ j \neq s_1, \dots, s_k.$$

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The RT sends the set of the seed systems into:

$$\tilde{K}_{1}^{-1}\boldsymbol{u}_{\tilde{t}_{1}} = \tilde{K}_{2}^{-1}\boldsymbol{u}_{\tilde{t}_{2}} = \dots = \tilde{K}_{n}^{-1}\boldsymbol{u}_{\tilde{t}_{n}},$$

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for any fixed $j \in \{1, \ldots, n\}$, where

$$\tilde{K}_{s_i} = -\sum_{j=1}^k \tilde{V}_{s_i}^{(n-s_j)} K_{s_j}, \quad i = 1, \dots, k,$$

$$\tilde{K}_m = K_m - \sum_{j=1}^k \tilde{V}_m^{(n-s_j)} K_{s_j}, \quad m = 1, 2, \dots, n, \quad m \neq s_i$$

It is readily verified that new d-Killing systems possesses infinite set of conservation laws:

$$D_{\tilde{t}_i}(\tilde{V}_j^{(m)}) = D_{\tilde{t}_j}(\tilde{V}_i^{(m)}), \quad i, j = 1, \dots, n, \quad i \neq j.$$



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$$\sum_{j=1}^{n} \Phi_{i}^{j}(\lambda^{i}) H_{j} = f_{i}(\lambda^{i}) \mu_{i}^{2} + \gamma_{i}(\lambda^{i}), \quad i = 1, \dots, n, \quad (8)$$

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 H_j are separable Hamiltonians and $f_i(\xi)$, $\gamma_i(\xi)$ are arbitrary functions of a single variable.

Solving (8) with respect to Hamiltonians $H_i = H_i(\lambda, \mu)$, i = 1, ..., n, where $\mu = (\mu_1, ..., \mu_n)^T$, we get

 $H_i = \mu^T K_i G_f \mu + V_i^{\gamma}(\lambda), \quad i = 1, \dots, n, \quad K_1 = \mathbb{I},$

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From the first half of Hamiltonian equations we have

$$\lambda_{t_k} = \frac{\partial H_i}{\partial \mu} = 2K_k G_f \mu, \qquad k = 1, ..., n.$$

We can rewrite it in the form

$$K_{1}^{-1}\boldsymbol{\lambda}_{t_{1}} = \dots = K_{n}^{-1}\boldsymbol{\lambda}_{t_{n}} \iff \frac{\lambda_{t_{1}}^{i}}{v_{1}^{i}} = \dots = \frac{\lambda_{t_{n}}^{i}}{v_{n}^{i}}, \quad i = 1, \dots, n,$$

$$(9)$$
as $K_{i}(\lambda) = \operatorname{diag}(\nu_{i}^{1}(\lambda), \dots, \nu_{i}^{n}(\lambda)), i = 1, \dots, n$ and

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$$v_r^i = (-1)^{r+1} \frac{\det \Phi^{ir}}{\det \Phi^{i1}}$$

where $\Phi = (\Phi_j^i(\lambda^i))$ is so called *Stäckel matrix* and Φ^{ik} is the $(n-1) \times (n-1)$ matrix obtained from Φ by removing its *i*th row and *k*th column.

On the other hand, such constructed functions v^i for arbitrary r are general solution (Ferapontov 1991) for the system of equations:

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$$\frac{\partial}{\partial \lambda_j} \left(\frac{\partial v^i / \partial \lambda^k}{v^k - v^i} \right) = \frac{\partial}{\partial \lambda_k} \left(\frac{\partial v^i / \partial \lambda^j}{v^j - v^i} \right), \qquad i \neq j \neq k \neq i = 1, \dots, n.$$
$$\frac{\partial v^i / \partial \lambda^i}{\partial \lambda^i} = 0, \qquad i = 1, \dots, n.$$

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The first equation means *semi-Hamiltonian* property while the second one *weakely nonlinear* property.

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where $\varphi_i(\xi)$ are arbitrary functions of a single variable

$$\varphi_i(\xi) = [f_i(\xi)(\sum_{j=1}^n \Phi_i^j(\xi)a_j - \gamma_i(\xi))]^{1/2}, \quad a_j = const.$$

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Separation relations for systems (7) after RT (Blaszak, Sergyeyev 2008): for i = 1, ..., n

$$\sum_{j=1}^{k} (\lambda^{i})^{\beta_{j}} \tilde{H}_{s_{j}} + \sum_{p=1, p \neq s_{1}, \dots, s_{k}}^{n} (\lambda^{i})^{n-p} \tilde{H}_{p} = f_{i}(\lambda^{i})\mu_{i}^{2} + \gamma_{i}(\lambda^{i}).$$

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$$\sum_{j=1}^{n} \int^{\lambda^{j}} \frac{\xi^{\beta_{q}}}{\varphi_{j}(\xi)} d\xi = \tilde{t}_{s_{q}}, \quad q = 1, \dots, k,$$

$$\sum_{j=1}^{n} \int^{\lambda^{j}} \frac{\xi^{n-i}}{\varphi_{j}(\xi)} d\xi = \tilde{t}_{i}, \quad i = 1, \dots, n, \quad i \neq s_{q}, \quad q = 1, \dots, k.$$

THE END