

# $\gamma$ DECAY OF GIANT RESONANCES WITHIN THE SKYRME FRAMEWORK

Marco Brenna

&

Gianluca Colò



UNIVERSITÀ  
DEGLI STUDI  
DI MILANO



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# Outline

- ✿ Introduction
- ✿ Formalism for the calculation of the  $\gamma$  decay width
  - ✿ Decay to the ground state
  - ✿ Decay to the low-lying states
- ✿ Results

# Giant Resonances

- ✿ Nuclear collective modes, in which almost all the nucleons are involved
- ✿ Parameters of the Equation of State of nuclear matter

Monopole	⇒	Compressibility
Dipole	⇒	Simmetry Energy

- ✿ Over than 60 year of studies (since 1947)
- ✿ More exclusive experiments feasible ⇒  $\gamma$  decay (LNL – INFN, June 2010)
  - ✿ Resonances in exotic nuclei (n – rich)
- ✿ Microscopically: coherent superposition of particle – hole excitation

✿ RPA (linear response) →

Fully self – consistent calculations with microscopic interactions (Skyrme, Gogny, RMF)

# Giant Resonances

## Main Properties

- ✱ Energy: 10 – 30 MeV
- ✱ Width: 2 – 5 MeV
- ✱ High percentage of EWSR

## Decay

- Particle emission (neutron)
- Coupling with doorway states → compound nucleus
- $\gamma$  decay

# Giant Resonances

## Decay

- Particle emission (neutron)
- Coupling with doorway states → compound nucleus
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- Suppressed with respect to particle emission ( $\sim 10^{-3}$ )
- Sensitive to resonance multipolarity
- Direct decay: complementary to inelastic scattering data  
(based on not completely under control ingredients –  
reaction model, optical potential...)
- Confront with experiment:  $\gamma$  decay from compound nucleus

# Decay width

## DECAY WIDTH

$$\Gamma_\gamma(E\lambda; i \rightarrow f) \propto E^{2\lambda+1} B(E\lambda; i \rightarrow f)$$

## REDUCED TRANSITION PROBABILITY

$$B(E\lambda; i \rightarrow f) = \frac{1}{2J_i + 1} |\langle J_f \| Q_\lambda^{(E)} \| J_i \rangle|^2$$

## ELECTROMAGNETIC OPERATOR (LONG-WAVELENGTH LIMIT)

$$Q_{\lambda\mu}^{(E)} = \sum_{i=1}^A e_i^\lambda i^\lambda r_i^\lambda Y_{\lambda\mu}^*(\hat{\mathbf{r}}_i)$$

## EFFECTIVE CHARGE DUE TO NUCLEAR RECOIL IN $E\lambda$ TRANSITIONS:

$$e_p^\lambda = e \left[ \left(1 - \frac{1}{A}\right)^\lambda + (-)^\lambda \frac{Z-1}{A^\lambda} \right] \quad e_n^\lambda = eZ \left(-\frac{1}{A}\right)^\lambda$$

# Nuclear Field Theory

Perturbative theory that describes the interweaving between single particle (fermionic) and phonon (bosonic) degrees of freedom

P. F. Bortignon et al., Phys. Rep.**30**(1977)305



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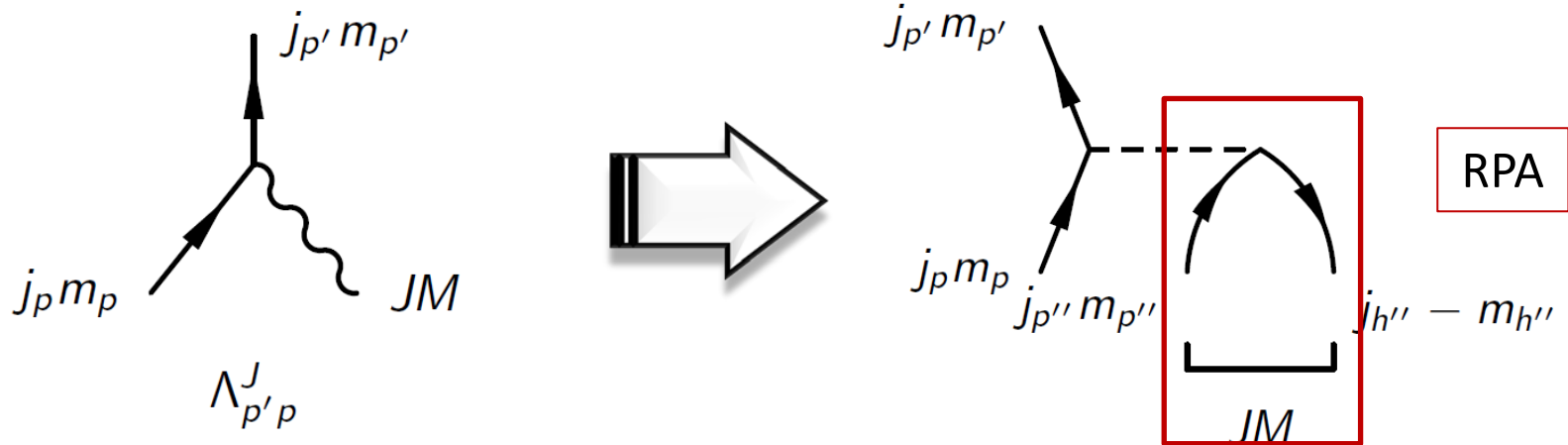
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- Single particle moments
- Multiplets of even/odd nuclei
- Giant resonances width  $\Gamma$
- **Appropriate framework to describe the decay to low-lying states**



# Particle-Vibration Coupling vertex (PVC)



$$\langle i || V || j, nJ \rangle = \sqrt{2J+1} \sum_{ph} X_{ph}^{nJ} V_J(ihjp) + (-)^{j_h - j_p + J} Y_{ph}^{nJ} V_J(ipjh)$$

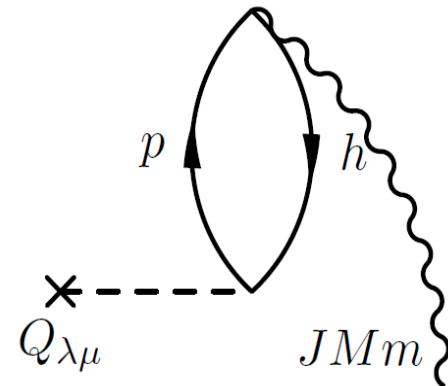
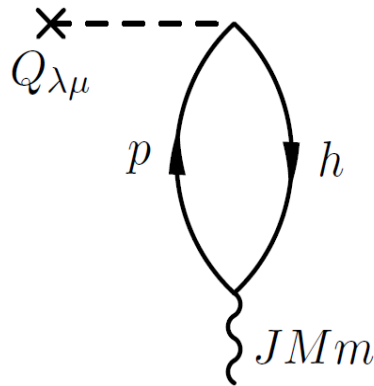
$$V_J(ihjp) = \sum_{\{m\}} (-)^{j_j - m_j + j_h - m_h} \langle j_i m_i j_j - m_j | JM \rangle \langle j_p m_p j_h - m_h | JM \rangle v_{ihjp}$$

Consistent treatment of the coupling vertex in the Skyrme framework:  
single particle states, RPA phonons, microscopic interaction

G. Colò, H. Sagawa, P. F. Bortignon, PRC82(2010)064307

# Decay to the ground state

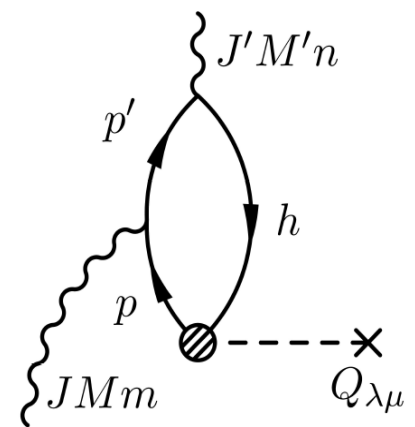
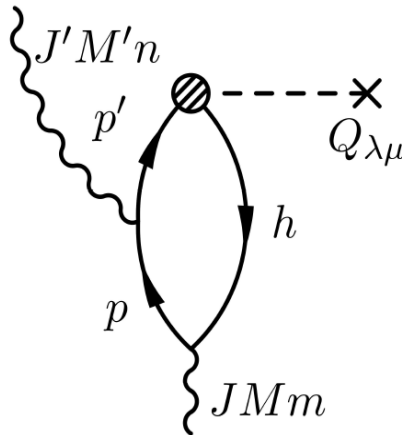
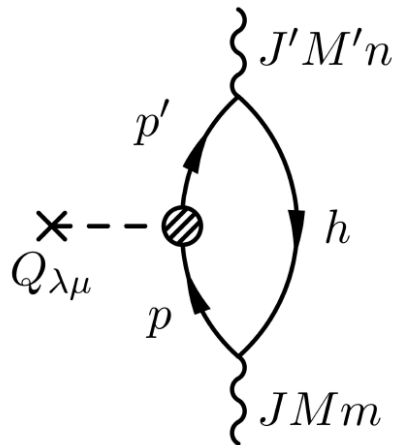
$$\langle 0 | Q_{\lambda\mu} | GR, JM \rangle$$



$$\langle 0 || Q_{\lambda} || GR, J \rangle = \sum_{ph} \langle p || Q_{\lambda} || h \rangle \left( \frac{\Lambda_{ph}^J}{E_J - \epsilon_{ph} + i\eta} - \frac{\Lambda_{ph}^J}{E_J + \epsilon_{ph} + i\eta'} \right)$$

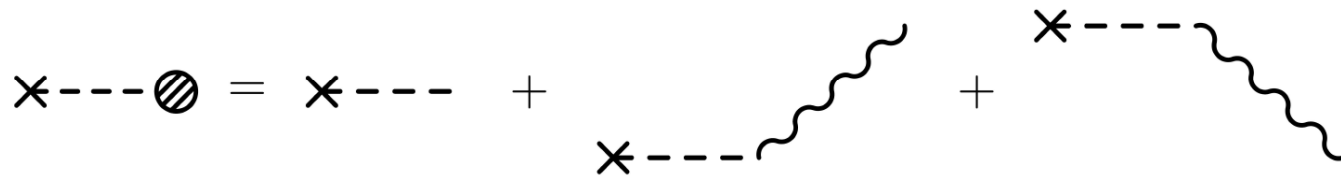
# Decay to low-lying states

NFT: 12 diagrams contribute to the matrix element



$$\begin{aligned}
 &= \sum_{pp'h} (-)^{J+\lambda+J'+1} \begin{Bmatrix} J & \lambda & J' \\ j_{p'} & j_h & j_p \end{Bmatrix} \frac{\langle p \| V \| h, nJ \rangle \langle h, mJ' \| V \| p' \rangle Q_{p'p}^{\lambda pol}}{(E_J - \epsilon_{ph} + i\eta) (\hbar\omega_{J'} - \epsilon_{p'h})}
 \end{aligned}$$

# Decay to low-lying states



Polarization charge:

External field partially screened through the interaction with intermediate states

$$Q_{ij}^{\lambda pol} = \langle i || Q_{\lambda} || j \rangle + \sum_{n'} \frac{1}{\sqrt{2\lambda + 1}} \left[ \frac{\langle 0 || Q_{\lambda} || n' \lambda \rangle \langle i, n' \lambda || V || j \rangle}{(E_J - \hbar\omega_{J'}) - \hbar\omega_{\lambda} + i\eta} - \frac{\langle i || V || j, n' \lambda \rangle \langle n' \lambda || Q_{\lambda} || 0 \rangle}{(E_J - \hbar\omega_{J'}) + \hbar\omega_{\lambda} + i\eta} \right]$$

# Results – $^{208}\text{Pb}$

- ✿ Consistent approach to the coupling vertex:
  - ✿ single particle states – HF
  - ✿ phonons – self consistent RPA with Skyrme functional
  - ✿ microscopic Skyrme interaction
- ✿ No phenomenological ingredient
- ✿ 4 Skyrme parametrization: SLy5, SGII, SkP, LNS
- ✿  $\gamma$  decay to the ground state and to first  $J^\pi = 3^-$  state of the Isoscalar Giant Quadrupole Resonance (ISGQR) in  $^{208}\text{Pb}$

# Energy and collectivity of the states

$J^\pi$	$2^+$		$3^-$	
	E [MeV]	EWSR [%]	E [MeV]	$B(E3) \uparrow [10^5 e^2 \text{fm}^6]$
Experimental	$10.9 \pm 0.3$	100	$2.6145 \pm 0.0003$	$6.11 \pm 0.09$
SLy5	12.28	69.27	3.62	6.54
SGII	11.62	69.66	2.92	6.83
SkP	10.28	81.79	3.29	5.11
LNS	12.19	64.90	3.37	5.46

Experimental data from NDS**108**(2007)*1583*

# Decay to the ground state

	$E_{GQR}(\text{MeV})$	$\Gamma_\gamma(\text{eV})$	
		RPA	RPA'
SLy5	12.28	231.54	160
SGII	11.62	170.09	154
SkP	10.28	119.18	169
LNS	12.19	176.19	135
Beene et al., PRC <b>39</b> (1989)1307	10.60	146±36 – <i>exp.</i>	
Speth et al., PRC <b>85</b> (1985)2310	10.60	112 – <i>theor.</i>	
Beene et al., PLB <b>164</b> (1985)19	11.20	175 – <i>theor.</i>	

Consistent with the experimental value through  
an energy and EWSR scaling of the GQR  
( $\Delta E = 1 \text{ MeV} \Rightarrow$  increase in  $\Gamma_\gamma$  of 40 %)

# Decay to the $3^-$ state

	$E_{tran}$ (MeV)	$\Gamma_\gamma$ (eV)
SLy5	8.66	5.07
SGII	8.70	33.25
SkP	6.99	10.99
LNS	8.82	54.60
Beene et al., PRC <b>39</b> (1989)1307	7.99	$5 \pm 5 - exp.$
Speth et al., PRC <b>85</b> (1985)2310	7.99	$4.00 - theor.$
Bortignon et al., PLB <b>148</b> (1984)20	8.60	$3.50 - theor.$



# Decay to the 3<sup>-</sup> state

SLy5

<b><math>\Gamma_\gamma</math> for a typical ph (eV) at 8.5 MeV</b>		<b><math>1.2 \cdot 10^3</math></b>
Quenching factors	<b>Recoupling</b>	<b>3</b>
	<b><math>\pi - \nu</math> cancellation</b>	<b>5</b>
	<b>p - h cancellation</b>	<b>3 - 4</b>
	<b>Polarization factor</b>	<b>4</b>
<b><math>\Gamma_\gamma</math>(eV)</b>		<b>5.07</b>

# Decay to the 3<sup>-</sup> state

SLy5

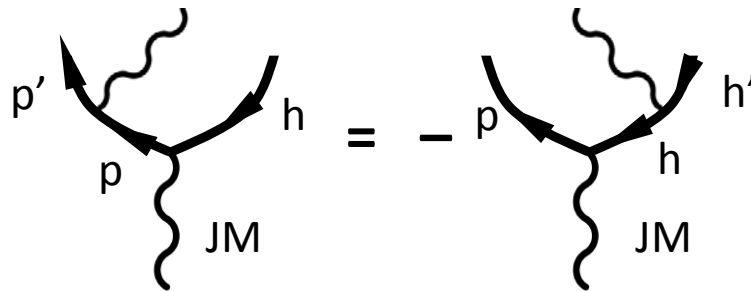
<b><math>\Gamma_\gamma</math> for a typical ph (eV) at 8.5 MeV</b>		<b><math>1.2 \cdot 10^3</math></b>
Quenching factors	<b>Recoupling</b>	3
	<b><math>\pi - \nu</math> cancellation</b>	5
	<b>p - h cancellation</b>	3 - 4
	<b>Polarization factor</b>	4
<b><math>\Gamma_\gamma</math>(eV)</b>		<b>5.07</b>

$$Q_{ij} = \left( \tau_z - \frac{N - Z}{A} \right)_j \langle i || r^\lambda Y_\lambda || j \rangle$$

# Decay to the 3<sup>-</sup> state

SLy5

$\Gamma_\gamma$ for a typical ph (eV) at 8.5 MeV	$1.2 \cdot 10^3$	
Quenching factors	<b>Recoupling</b>	3
	$\pi - \nu$ <b>cancellation</b>	5
	<b>p - h</b> cancellation	3 - 4
	<b>Polarization factor</b>	4
$\Gamma_\gamma$ (eV)	5.07	





# Conclusions

- Microscopic and consistent treatment  $\gamma$  decay
- $\gamma$  decay to the GS: not able to discriminate between models
- $\gamma$  decay to the  $3^-$  state: sensitive to the interaction used
  - Dipole spectrum
- Comparison with the experiment at LNL – INFN (June 2010)
- Other closed – shell nuclei:  $^{90}\text{Zr}$  (LNL - 2010) ...



# Decay of the compound nucleus

$$\langle \Gamma_{\gamma 0}^{CN} \rangle = \frac{X(\lambda) b_{E\lambda}(E) \left(\frac{E}{\hbar c}\right)^{2\lambda+1}}{\rho_I(E)}$$

$$X(\lambda) = \frac{8\pi(\lambda + 1)}{\lambda[(2\lambda + 1)!!]^2}$$

- $\rho_I(E)$ : density of state of the compound nucleus with spin  $I$  and energy  $E$

- $b_{E\lambda}(E)$  reduced transition probability per unit of energy

