

Transport properties of β -stable nuclear matter

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Outline

- 1 Introduction
 - Motivations & Objectives

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- 2 Theory and tools for calculation
 - Nuclear Matter & Neutron stars structure
 - Landau theory for normal Fermi liquids
 - Effective potential, effective mass, proton fraction

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 - Effective potential, effective mass, proton fraction
- 3 Results
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 - Conclusions & Outlooks

Motivations & Objectives

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- Oscillating modes of rotating stars are driven unstable through the emission of gravitational waves (CFS instability)
- Dissipation effects interior to the stars, as viscosity, can damp, or even suppress, these unstable oscillating modes
- **The knowledge of the viscosity is an essential element to predict the stability of a rotating star**

Nuclear Matter

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Nuclear Matter

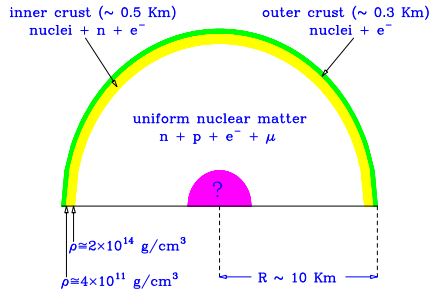
- **Equilibrium properties:** obtained through many body theory with realistic dynamical models
- **Transport properties:** usually obtained with simpler dynamical models

Our effort has been put in calculating both equilibrium and transport properties with the same dynamical model:

Landau-AK formalism using the effective potential derived in the Correlated Basis Function (CBF) theory

Neutron stars structure

External core: homogeneous fluid of neutrons, protons and electrons
 stable with respect to β -decay and electronic capture



Landau theory for normal Fermi liquids

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Quasiparticle distribution

$$n_{\mathbf{p}} = \frac{1}{e^{\beta(\epsilon_{\mathbf{p}}[n_{\mathbf{p}}] - \mu)} + 1}$$

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- The Boltzmann-Landau transport equation

$$\frac{\partial n_{\mathbf{p}}(\mathbf{r}, t)}{\partial t} + \frac{d\mathbf{r}}{dt} \frac{\partial n_{\mathbf{p}}(\mathbf{r}, t)}{\partial \mathbf{r}} + \frac{d\mathbf{p}}{dt} \frac{\partial n_{\mathbf{p}}(\mathbf{r}, t)}{\partial \mathbf{p}} = I[n_{\mathbf{p}}(\mathbf{r}, t)]$$

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- Considering a small deviation from equilibrium of the state of the system, $n_{\mathbf{p}} = n_{\mathbf{p}}^0 + \delta n_{\mathbf{p}}$, the transport equation reads

$$\mathbf{v}_{\mathbf{p}} \cdot \nabla_{\mathbf{r}} \delta n_{\mathbf{p}} = I[\delta n_{\mathbf{p}}]$$

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- The collisions integral

$$\begin{aligned} I[n_1] &= \frac{1}{V^2} \sum_2 \sum_{34} \delta_{\mathbf{p}_1 + \mathbf{p}_2, \mathbf{p}_3 + \mathbf{p}_4} \delta_{\sigma_1 + \sigma_2, \sigma_3 + \sigma_4} \delta(\epsilon_1 + \epsilon_2 - \epsilon_3 - \epsilon_4) \\ &\times W(12; 34) [n_3 n_4 (1 - n_1) (1 - n_2) - n_1 n_2 (1 - n_3) (1 - n_4)] \end{aligned}$$

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- Rewrite the transport equation in terms of a regular variable Φ

$$\delta n_i \equiv -\frac{\partial n_i^0}{\partial \epsilon_i} \Phi_i$$

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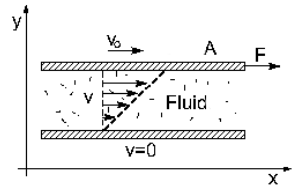
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- Write the shear viscosity tensor in terms of a current of the system

$$\sigma'_{xy} = \eta \frac{\partial u_x}{\partial y} = -\sum_{\sigma} \int \frac{d^3 p}{(2\pi\hbar)^3} p_x (v_{\mathbf{p}})_y \delta n_{\mathbf{p}}$$



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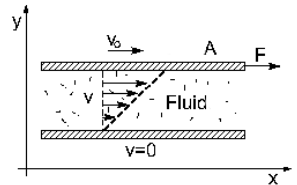
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- Solve the transport equation

Landau theory for normal Fermi liquids

Abrikosov-Khalatnikov solution (1959):

$$\eta_{AK} = \frac{1}{5} \rho m^* v_F^2 \tau \frac{2}{\pi^2 (1 - \lambda_\eta)}$$

$$\tau \equiv \frac{8\pi^4}{m^{*3} \langle W \rangle T^2}, \quad \langle W \rangle \equiv \int \frac{d\Omega}{2\pi} \frac{W(\theta, \phi)}{\cos(\theta/2)}, \quad \lambda_\eta = \frac{\langle W(\theta, \phi) [1 - 3\sin^4(\theta/2) \sin^2 \phi] \rangle}{\langle W(\theta, \phi) \rangle}$$

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Brooker-Sykes solution (1968):

$$\eta = \frac{1}{5} \rho m^* v_F^2 \tau \frac{2}{\pi^2 (1 - \lambda_\eta)} C(\lambda_\eta)$$

$$C(\lambda_\eta) = \frac{1 - \lambda_\eta}{4} \sum_{k=0}^{\infty} \frac{4k+3}{(k+1)(2k+1)[(k+1)(2k+1) - \lambda_\eta]}$$

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Comparison between solutions:

$$0.75 < \frac{\eta}{\eta_{AK}} < 0.92$$

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External core of neutron stars as a multicomponent Fermi liquid:

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External core of neutron stars as a multicomponent Fermi liquid:

- Three components Boltzmann-Landau transport equation

$$\frac{\partial n_\alpha}{\partial t} + \frac{\partial n_\alpha}{\partial \mathbf{r}} \cdot \frac{\partial \epsilon_{\mathbf{p}\alpha}}{\partial \mathbf{p}} - \frac{\partial n_\alpha}{\partial \mathbf{p}} \cdot \frac{\partial \epsilon_{\mathbf{p}\alpha}}{\partial \mathbf{r}} = I_\alpha$$

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- Three components shear viscosity

$$\eta = \eta_n + \eta_p + \eta_e$$

$$\eta_\alpha = \frac{1}{5} \rho_\alpha m_\alpha^* v_{F\alpha}^2 \tau_\alpha \frac{2}{\pi^2 (1 - \lambda_{\alpha\alpha})} C(\lambda_{\alpha\alpha})$$

$$\tau_\alpha = \frac{8\pi^4}{\sum_\beta m_\alpha^* (m_\beta^{*2} \langle W_{\alpha\beta} l_{\alpha\beta} \rangle) T^2}, \quad \lambda_{\alpha\alpha} = \frac{\sum_\beta \langle W_{\alpha\beta}(\theta, \phi) l_{\alpha\beta}^\alpha \rangle}{\sum_\beta \langle W_{\alpha\beta}(\theta, \phi) l_{\alpha\beta} \rangle}$$

The effective potential

- Two body cluster contribution in the CBF theory:

$$v_{\text{eff}}(ij) = f_{ij} \left(-\frac{1}{m} \nabla^2 + v(ij) \right) f_{ij} = \sum_{n=1}^6 v_{\text{eff}}^n(r_{ij}) O_{ij}^n$$

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- Three-nucleon contributions are taken into account considering the approach proposed by Lagaris-Pandharipande

$$v(ij) \rightarrow \tilde{v}(ij) = \sum [v_{\pi}^n(r_{ij}) + v_I^n(r_{ij})e^{-\gamma_1 \rho} + v_S^n(r_{ij})] O_{ij}^n$$

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- Calculate the scattering cross section and obtain the scattering probability to evaluate the mean lifetime of the quasiparticle state

$$W = W(\theta, \phi) = \left(\frac{2\pi}{\mu} \right)^2 \left(\frac{d\sigma}{d\Omega} \right)_{cm} [E_{cm} = f(\theta), \theta_{cm} = \phi]$$

The effective neutron mass

- Calculate the single particle energy using the effective potential

$$e_{\lambda}(p) = \frac{p^2}{2m} + \frac{\rho}{2} \sum_{\mu} \sum_n x_{\mu} \int d^3r v_{\text{eff}}^n(r) [A_{\lambda\mu}^n - B_{\lambda\mu}^n j_0(pr) \ell(p_F r)]$$

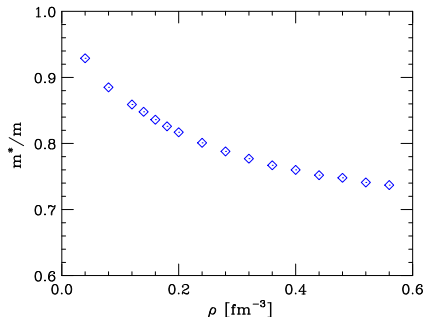
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- Obtain the effective neutron mass

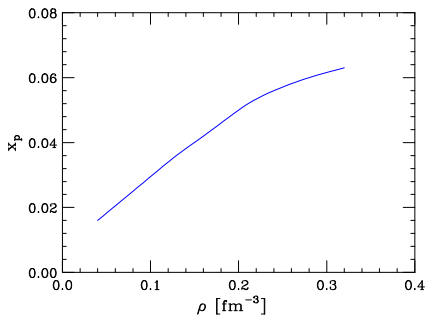
$$\frac{1}{m^*} = \frac{1}{p} \frac{de}{dp}$$



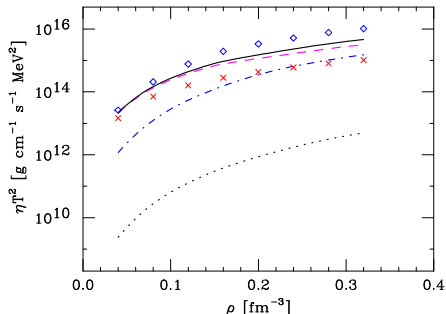
Proton fraction

Solve the coupled equations which define the equilibrium of the system:

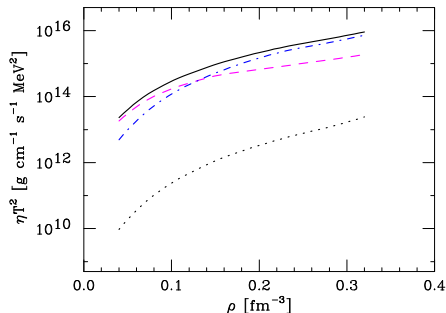
- β -stability: $\mu_n = \mu_p + \mu_e$
- charge neutrality: $\rho_p = x_p \rho = \rho_e$



Results I:



Results II: double proton fraction



- solid line= β -stable matter; dashed=n; dot-dashed=e; dotted=p; diamonds=PNM; cross=Argonne v8'

Conclusions & Outlooks

- We calculated the viscosity of the external core of a neutron star, using a realistic model of nuclear matter and of the dynamic to its interior
- The percentage of electrons play a fundamental role in determining the weight of the viscosity of each particle family
- This work can be improved including the presence of muons and successively hyperons (if that's what really neutron stars are made of!)
- Calculate consistently the shear viscosity of the superfluid / superconducting phase

Thank you! ¹

¹Arianna Carbone, PhD student with Dr. Artur Polls & Dr. Arnau Rios on
“Three body forces with self-consistent Green’s functions”.

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