

Superscaling in neutrino/antineutrino CCQE scattering from MiniBooNE to NOMAD energies

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- SuSA and RFG
- Form Factors' Parametrizations
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3 Conclusions

- Summary and Main Conclusions
- Further Work

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Quasielastic Regime

CCQE scattering

$$\nu_\mu(\bar{\nu}_\mu) + A \rightarrow \mu^-(\mu^+) + p(n) + (A-1)$$

Impulse Approximation (IA)

The neutrino only interacts with a single bound nucleon.

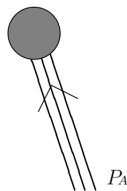
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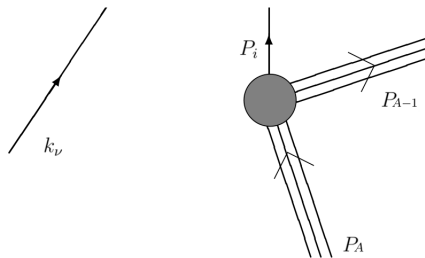
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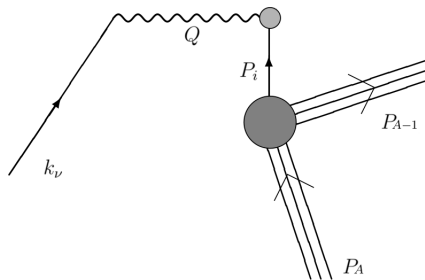
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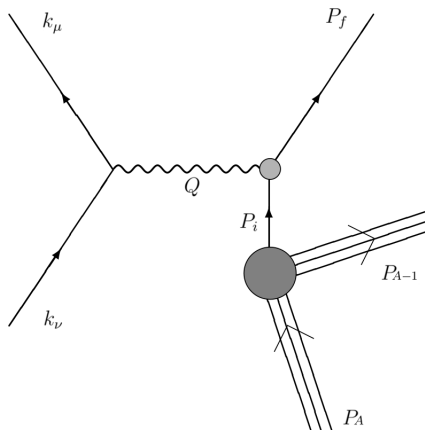
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Motivation

Two main related objectives:

- 1 Complete theoretical description of the CCQE neutrino-nucleus interaction and the weak structure of the nucleon.
- 2 Full analysis of the experimental data in all range of energies from intermediate $E_\nu \sim 1 \text{ GeV}$ (MiniBooNE) to high values $E_\nu \sim 10 - 100 \text{ GeV}$ (NOMAD).

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Relevance of this investigation:

- To know better the hadronic structure of the nucleon and other nuclear properties such as correlations or 2p-2h MEC contributions.
- To analyze better neutrino oscillations experiments.

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MiniBooNE

MiniBooNE (Fermilab)

- Measurement of CCQE ν_μ ($\bar{\nu}_\mu$) cross sections on a ^{12}C target in the 1 GeV region.
- Discrepancy between the data and traditional nuclear models.

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Two options to solve this puzzle:

- 1 MiniBooNE proposed a higher nucleon axial mass value: $M_A = 1,35 \text{ GeV}/c^2 \implies$ It does not work at high energy data (NOMAD).
 - 2 Microscopic explanations based on multi-nucleon excitations, such as the evaluation of the Meson Exchange Currents (MEC) within the 2p-2h approach \implies Do not give a full account for the discrepancy.
- ➡ A consistent evaluation of MEC is hard to achieve.

NOMAD

NOMAD (CERN)

- ⇒ CCQE $\nu_\mu(\bar{\nu}_\mu)-^{12}\text{C}$ cross sections measurements go from 3 to 100 GeV.
- ⇒ Results do not call for a large axial-vector mass (M_A) and do not seem to match with the MiniBooNE results.

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Main goal

- ➡ To perform a consistent theoretical analysis of the CCQE neutrino-nucleus interaction in the entire 0-100 GeV region.

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Theoretical descriptions: Nuclear model dependence

- ➡ For this purpose we need to employ a nuclear model which can be applied up to very high energies.
- ➡ Two basic requirements: it has to be relativistic and it must describe QE electron scattering data from intermediate up to high energies.

SuperScaling Approach (SuSA)

- Based on the superscaling function extracted from QE electron scattering data.
- This model does not account for the $\sim 10\%$ scaling violations of the transverse channel, which are associated with 2p-2h MEC. This should therefore be added.
- An evaluation of MEC is very hard to achieve. Our present MEC parametrization does not work well for $E_\nu \gtrsim 2\text{ GeV}$ (work in progress).

Theoretical descriptions: Nuclear model dependence

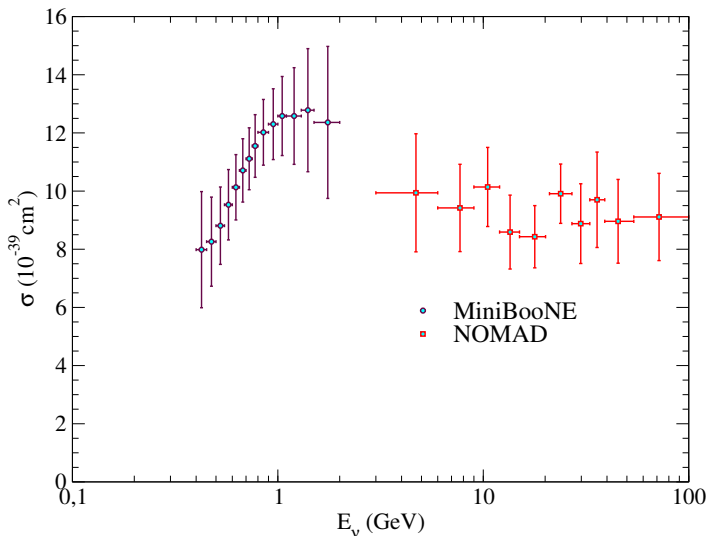
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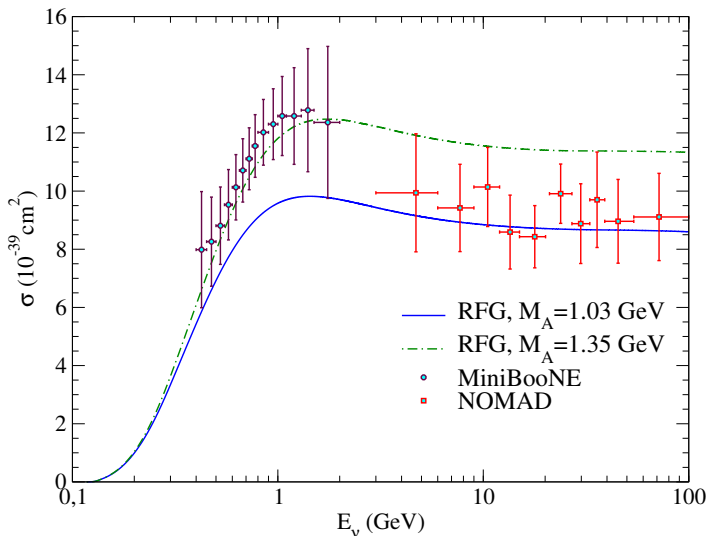
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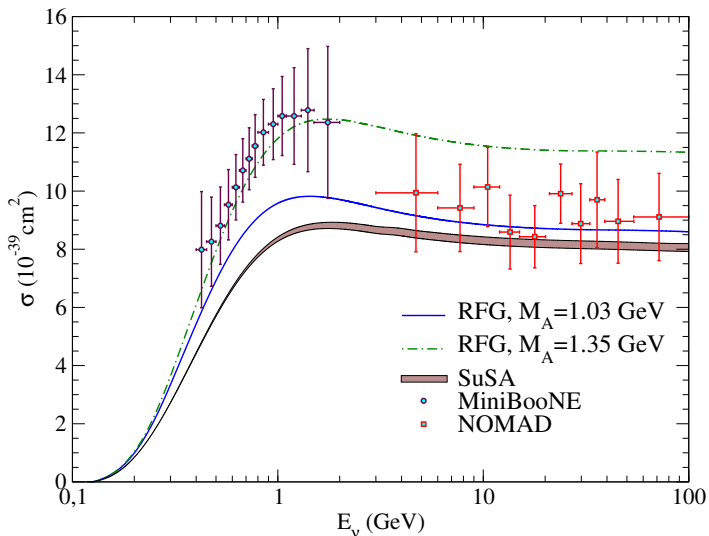
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ν_μ - ^{12}C CCQE scattering

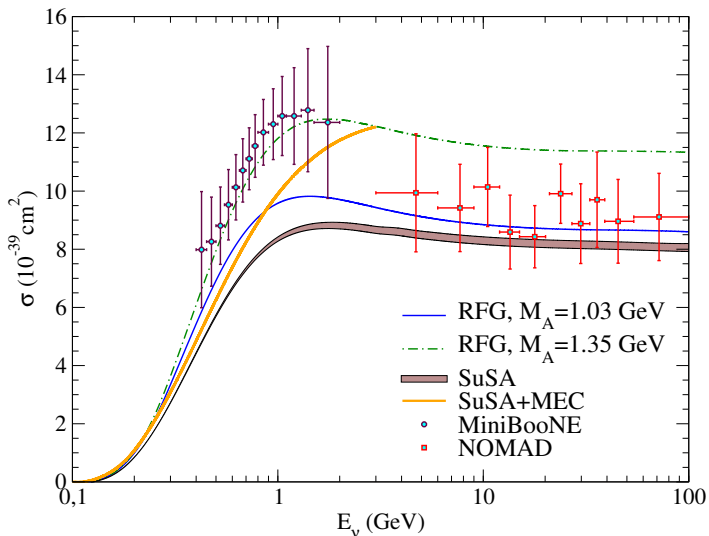
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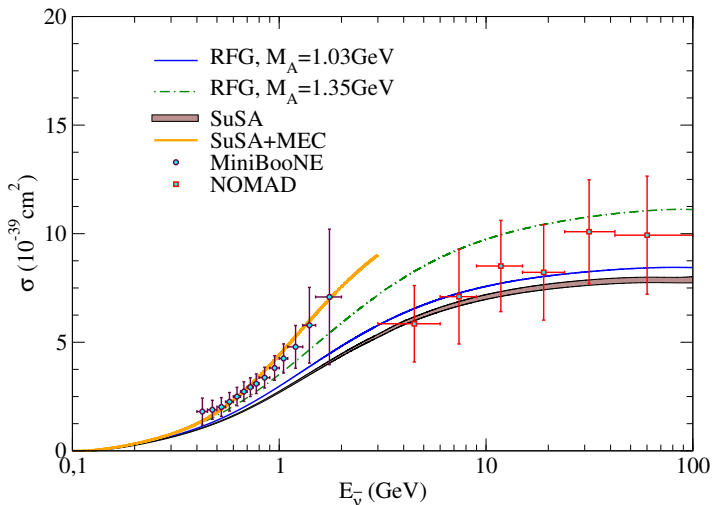


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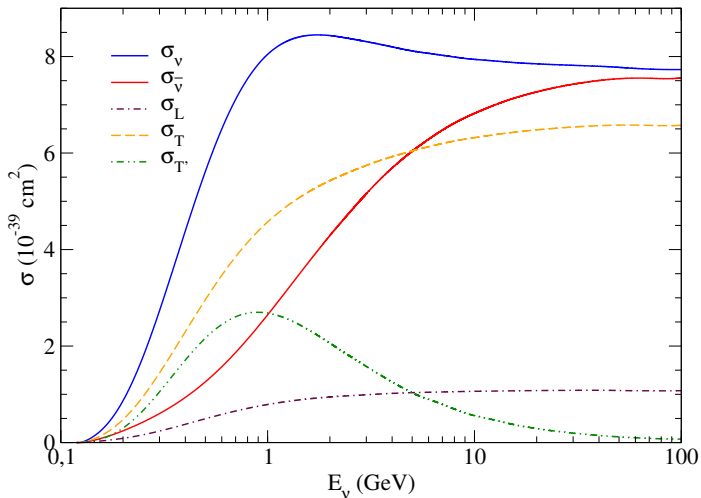


ν_μ - ^{12}C CCQE scattering



$\bar{\nu}_\mu$ - ^{12}C CCQE scattering

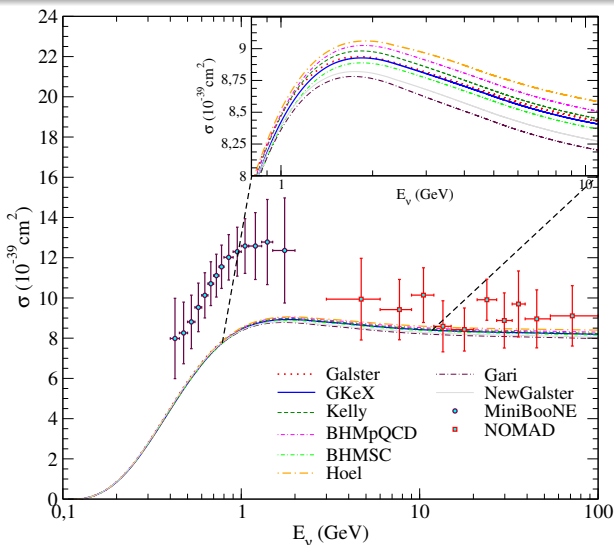
Separated Contributions in the SuSA Model



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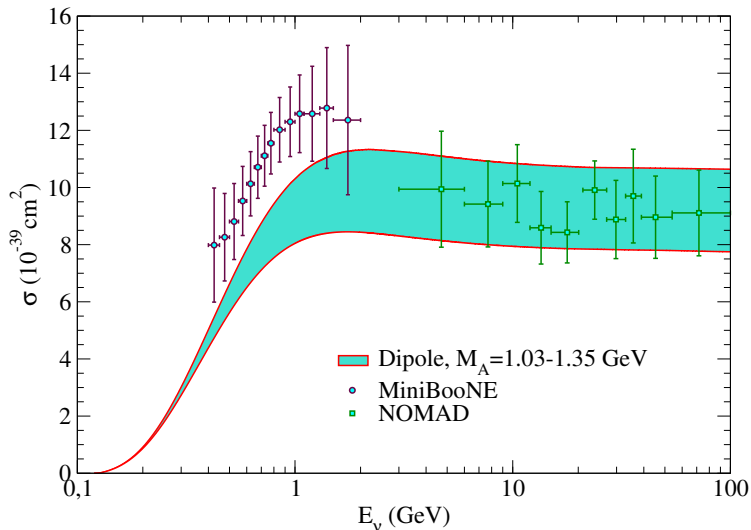
Parametrization of the nucleon EM form factors



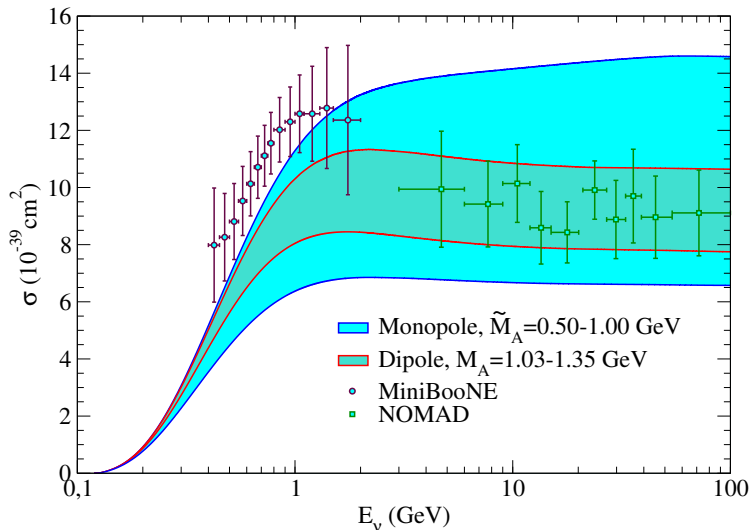
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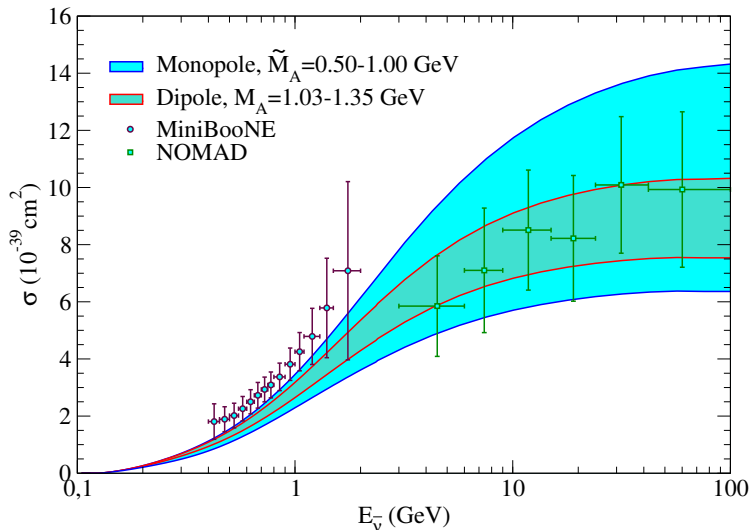
Dipolar axial form factor



Monopolar axial form factor



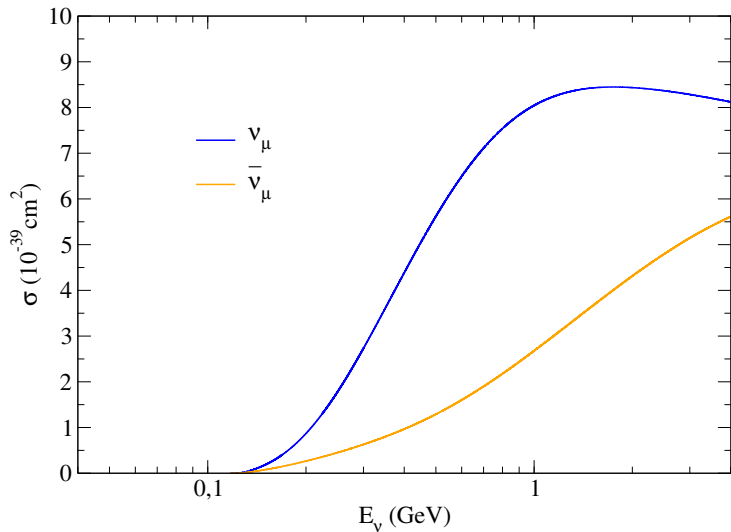
Monopolar axial form factor

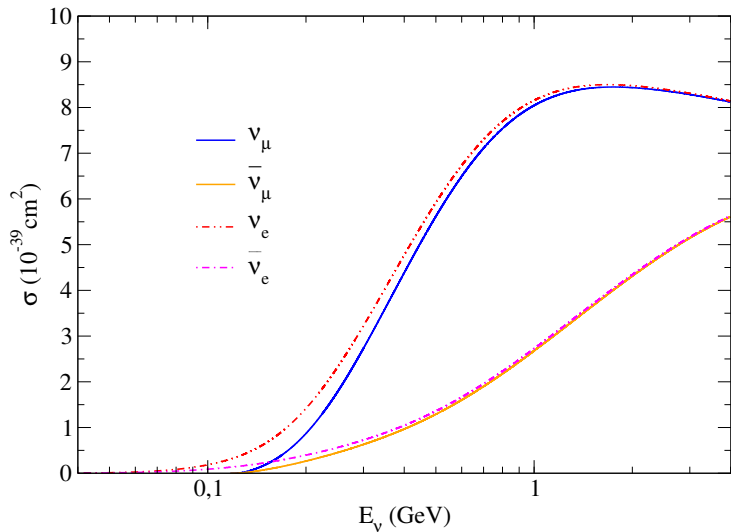


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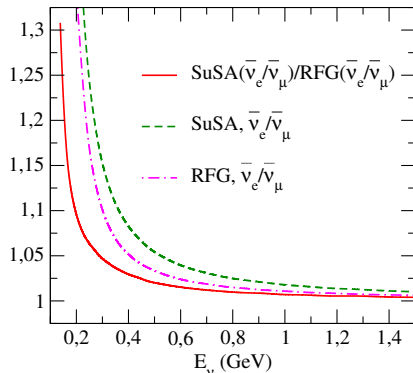
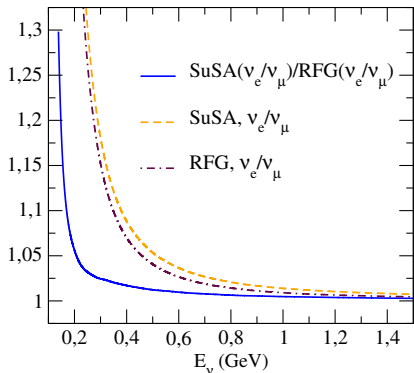
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ν_μ vs. ν_e CCQE Cross Section



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Summary and Main Conclusions

arXiv:1305.6884 [nucl-th]

- First SuSA results from intermediate to high energies. In spite of similar results between RFG and SuSA, RFG fails to reproduce (e, e') data whereas SuSA agrees with them by construction.
- SuSA model has to be completed with effects that go beyond the Impulse Approximation (correlations, MEC) \Rightarrow increase of the cross sections at low energies (10-15 %) \Rightarrow better data explanation without increasing M_A .
- Moreover, our predictions corresponding to ν STORM kinematics can be useful to get information about the electroweak nuclear matrix elements and the dipole or monopole nature of the axial-vector form factor.



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Further Work

ν_μ →

n →

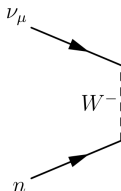
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➤ Recent CCQE $\nu_\mu(\bar{\nu}_\mu)$ - ^{12}C data from the Minerva Collaboration reject $M_A = 1,35 \text{ GeV}$ in favor of using $M_A \approx 1 \text{ GeV}$ and a TEM enhancement.



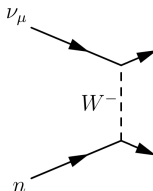
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- Recent CCQE $\nu_\mu(\bar{\nu}_\mu)^{-12}\text{C}$ data from the Minerva Collaboration reject $M_A = 1,35$ GeV in favor of using $M_A \approx 1$ GeV and a TEM enhancement.
- SuSAv2: two different scaling functions for the longitudinal and the transverse channel, taking into account the isovector structure of the CC.



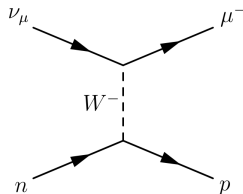
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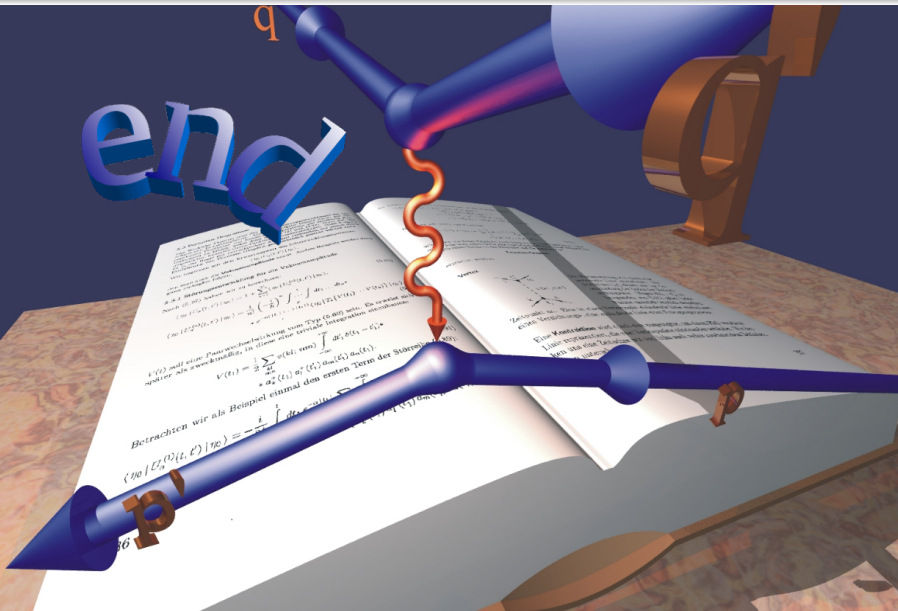


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- Evaluation of MEC at higher energies and extension to the VA interference channel (T').
- Study of future νSTORM , Minerva and T2K results.



end



$$V(t) = \frac{1}{2} \sum_{k=1}^{\infty} \omega(k, \tau_m) \int_{-\infty}^{\infty} dx_1^2 \delta(x - t_1) \cdot$$

$$+ a_2^2(t_1) a_1^2(t_1) \delta(t_1) a_m(t_1) a_m(t_1)$$

Betrachten wir als Beispiel einmal den ersten Term der Störreihe

$$\langle \tau_0 | U_0^{(1)}(t, t') | \tau_0 \rangle = -\frac{i}{2} \int_{t'}^t dt_1 \int_{-\infty}^{\infty} dx_1^2 \delta(x - t_1) \cdot$$