

$$H_0 = -\frac{\hbar^2}{2m} \nabla^2$$

$$H_0 |\phi_k\rangle = E |\phi_k\rangle$$

$$E = \frac{\hbar^2 k^2}{2m}$$

$$\langle r | \phi_k \rangle = e^{i\vec{k} \cdot \vec{r}}$$

$$H = -\frac{\hbar^2}{2m} \nabla^2 + V(\vec{r})$$

$$H |\psi_{ka}^+\rangle = E |\psi_{ka}^+\rangle$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + \tilde{U} \right] \tilde{\psi}_b^{(-)} = E \tilde{\psi}_b^{(-)}$$

$$\psi_b^{(-)} \xrightarrow{r \rightarrow \infty} e^{i\vec{k}_b \cdot \vec{r}} + f_b^{(-)}(\Omega) \frac{e^{-ikr}}{r}$$

$$\left[-\frac{\hbar^2}{2m} \nabla^2 + U \right] \psi_a^{(+)} = E \psi_a^{(+)}$$

$$\psi_a^{(+)} \xrightarrow{r \rightarrow \infty} e^{i\vec{k}_a \cdot \vec{r}} + f_a^{(+)}(\Omega) \frac{e^{ikr}}{r}$$

Moltiplicando a sinistra la prima eq. per $\psi_a^{(+)*}$ e $\psi_b^{(-)}$ il complesso coniugato della seconda si può dimostrare che:

$$\langle \tilde{\psi}_b^{(-)} | U - \tilde{U} | \psi_a^{(+)} \rangle = -\frac{2\pi\hbar^2}{m} \left[f_a^{(+)}(\Omega_b) - f_b^{(-)*}(-\Omega_a) \right]$$

A) $U=V \quad \tilde{U}=0$

$$\langle \phi_b | V | \psi_a^+ \rangle = -\frac{2\pi\hbar^2}{m} f_a^{(+)}(\Omega_b)$$

B) $U=\tilde{U}=V$

$$f_a^{(+)}(\Omega_b) = f_b^{(-)*}(-\Omega_a)$$

C) $U=0 \quad \tilde{U}=V$

$$\langle \psi_b^{(-)} | V | \phi_a \rangle = -\frac{2\pi\hbar^2}{m} f_b^{(-)*}(-\Omega_a)$$