

In questo limite

$$\frac{j_{\bar{\sigma}+1}(kr)}{j_{\bar{\sigma}-1}(kr)} \xrightarrow{kr \rightarrow 0} \frac{(kr)^{\bar{\sigma}+1}}{(2\bar{\sigma}+3)!!} \frac{(2\bar{\sigma}-1)!!}{(kr)^{\bar{\sigma}-1}} = \frac{(kr)^2}{(2\bar{\sigma}+3)(2\bar{\sigma}+1)}$$

questo per affermare che il termine $\bar{\sigma}-1$ è dominante in questo limite.

La parte longitudinale ha una struttura analoga

$$\vec{A}_{\bar{\sigma}M}(\vec{r}, b) = \left(\frac{\bar{\sigma}}{2\bar{\sigma}+1}\right)^{1/2} j_{\bar{\sigma}-1}(kr) \frac{\vec{Y}_{\bar{\sigma}-1}^M(\hat{r})}{\bar{\sigma}, \bar{\sigma}-1} + \left(\frac{\bar{\sigma}+1}{2\bar{\sigma}+1}\right)^{1/2} j_{\bar{\sigma}+1}(kr) \frac{\vec{Y}_{\bar{\sigma}+1}^M(\hat{r})}{\bar{\sigma}, \bar{\sigma}+1} = \frac{1}{k} \vec{\nabla} \cdot j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M}$$

escludendo i termini $\bar{\sigma}+1$ si ha che

$$\begin{aligned} \vec{A}_{\bar{\sigma}M}(\vec{r}, e) &\xrightarrow{kr \ll 1} \left(\frac{\bar{\sigma}+1}{2\bar{\sigma}+1}\right)^{1/2} j_{\bar{\sigma}+1}(kr) \frac{\vec{Y}_{\bar{\sigma}+1}^M(\hat{r})}{\bar{\sigma}, \bar{\sigma}+1} \rightarrow \left(\frac{\bar{\sigma}+1}{\bar{\sigma}}\right)^{1/2} \vec{A}_{\bar{\sigma}M}(\vec{r}, b) \\ &\rightarrow \left(\frac{\bar{\sigma}+1}{\bar{\sigma}}\right)^{1/2} \frac{1}{k} \vec{\nabla} \cdot j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M} \end{aligned}$$

$$\begin{aligned} M_{\beta\alpha}(\vec{E}\vec{\sigma}) &= \sqrt{2\pi} [\bar{\sigma}] (i)^{\bar{\sigma}+1} \langle \beta | \vec{J} | \alpha \rangle \int d^3r \vec{A}(\vec{r}, e) \Rightarrow kr \ll 1 \\ &= \sqrt{2\pi} [\bar{\sigma}] (i)^{\bar{\sigma}+1} \langle \beta | \vec{J} | \alpha \rangle \left(\frac{\bar{\sigma}+1}{\bar{\sigma}}\right)^{1/2} \frac{1}{k} \vec{\nabla} \cdot j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M} d^3r = \text{per parti} \end{aligned}$$

$$= \sqrt{2\pi} [\bar{\sigma}] (i)^{\bar{\sigma}+1} \left(\frac{\bar{\sigma}+1}{\bar{\sigma}}\right)^{1/2} \frac{1}{k} \left[\langle \beta | \vec{J} | \alpha \rangle \cdot j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M} \right]_{-\infty}^{\infty} - \int \vec{\nabla} \cdot \langle \beta | \vec{J} | \alpha \rangle j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M} d^3r$$

= il primo termine è nullo e nel secondo uso l'eq. di continuità (65)

$$\begin{aligned} &= \sqrt{2\pi} [\bar{\sigma}] (i)^{\bar{\sigma}+1} \left(\frac{\bar{\sigma}+1}{\bar{\sigma}}\right)^{1/2} \frac{1}{k} \cdot \frac{\vec{E}_{\beta} - \vec{E}_{\alpha}}{\hbar} \int \langle \beta | \vec{J} | \alpha \rangle j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M} d^3r \\ &= \sqrt{2\pi} (i)^{\bar{\sigma}} \left[\frac{(2\bar{\sigma}+1)(\bar{\sigma}+1)}{\bar{\sigma}} \right]^{1/2} \frac{1}{k} \cdot \frac{1}{\hbar} \cdot \frac{1}{kc} \int \langle \beta | \vec{J} | \alpha \rangle j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M} d^3r \\ &= -c \sqrt{2\pi} (i)^{\bar{\sigma}} \left[\frac{(2\bar{\sigma}+1)(\bar{\sigma}+1)}{\bar{\sigma}} \right]^{1/2} \int j_{\bar{\sigma}}(kr) \frac{Y_{\bar{\sigma}}^M(\hat{r})}{\bar{\sigma}M} \langle \beta | \vec{J} | \alpha \rangle d^3r \end{aligned}$$

eq. 73) è nota come teorema di Siegert.