

Neutrino cross sections and nuclear structure

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NOW08 Conca Specchiulla, Sept. 11, 2008

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Aim of the talk

How the uncertainties on the nuclear structure affects the
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Specific case

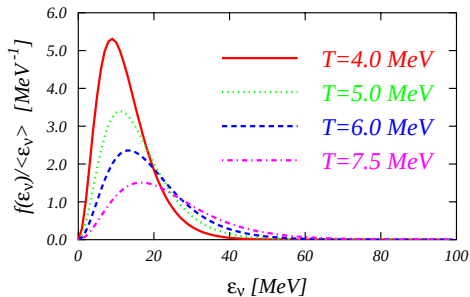
Precision ν astronomy

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How the uncertainties on the nuclear structure affects the
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Precision ν astronomy
SUPERNOVAE



Disentangling

$$\nu_x, \bar{\nu}_x \rightarrow T = 6.0 \text{ MeV} \quad T = 7.5 \text{ MeV}$$

Excitation energy $\omega \leq 100 \text{ MeV}$

Momentum transfer $q \leq 1.0 \text{ fm}^{-1}$

Ab initio calculations

Realistic nucleon-nucleon interaction

Nuclear Structure

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Effective theories

Nucleon-nucleon interactions linked to the theory

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How the uncertainties related to the choice of the effective interaction affect the ν -nucleus cross section?

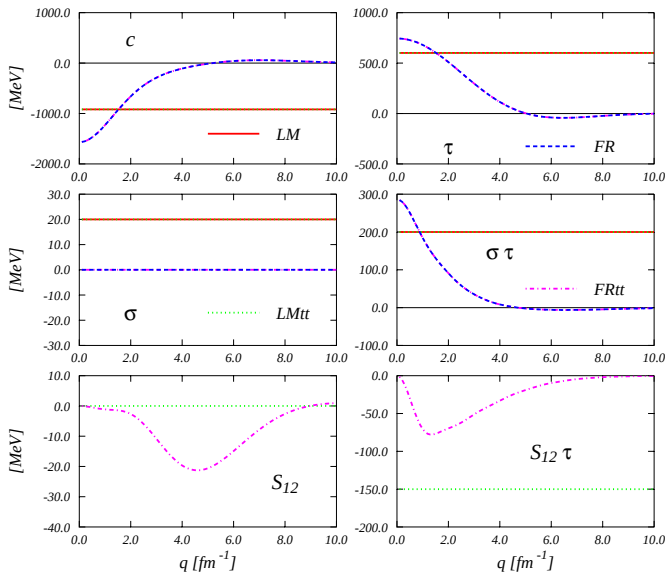
$$\begin{aligned} V_{\text{eff}}(r) &= V_1(r, \rho) + V_2(r) \vec{\tau}(1) \cdot \vec{\tau}(2) \\ &+ V_3(r) \vec{\sigma}(1) \cdot \vec{\sigma}(2) + V_4(r) \vec{\sigma}(1) \cdot \vec{\sigma}(2) \vec{\tau}(1) \cdot \vec{\tau}(2) \\ &+ V_5(r) S_{12}(r) + V_6(r) S_{12}(r) \vec{\tau}(1) \cdot \vec{\tau}(2) \end{aligned}$$

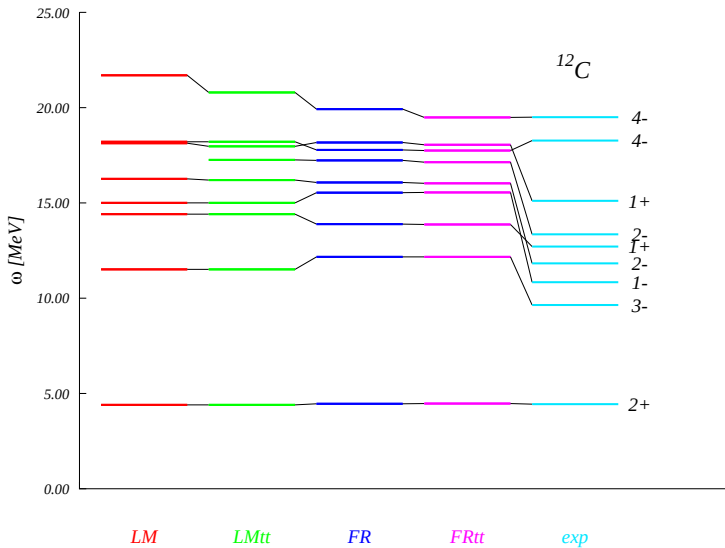
$$S_{12}(r) = 3 \frac{\vec{\sigma}(1) \cdot \vec{r} \vec{\sigma}(2) \cdot \vec{r}}{r^2} - \vec{\sigma}(1) \cdot \vec{\sigma}(2)$$

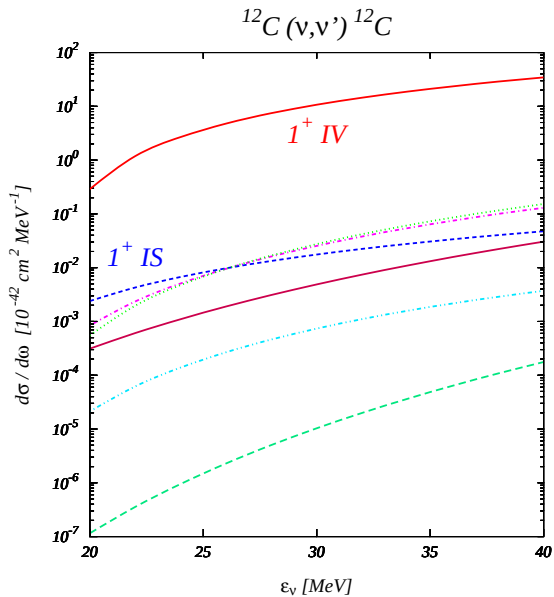
$$V_1(r, \rho) = V_c(r) + V_\rho(r) \rho(r)^\alpha$$

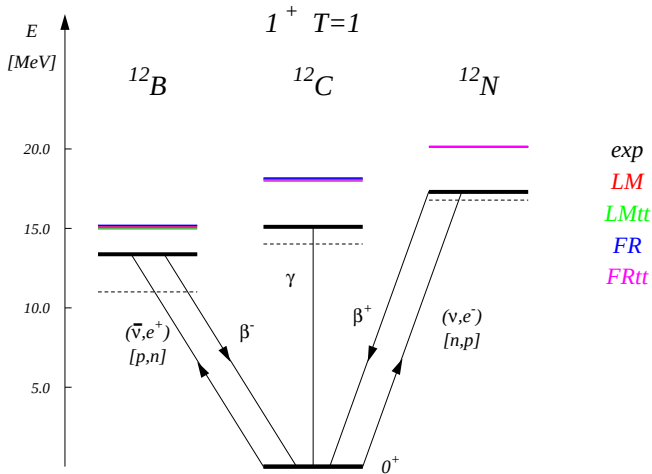
$$\rho(r_{12}) = [\rho(r_1) \rho(r_2)]^{1/2}$$

$$\alpha = 1$$

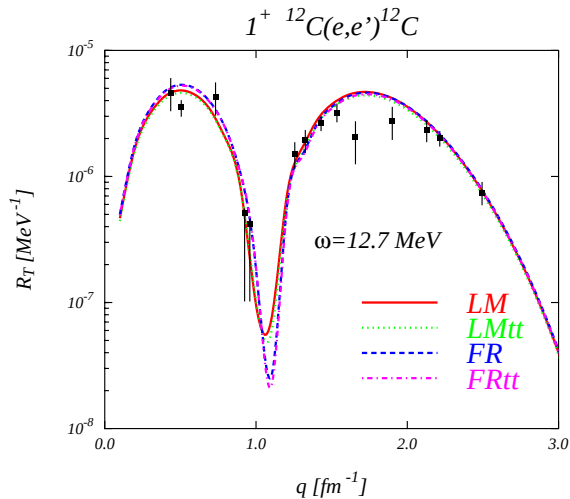


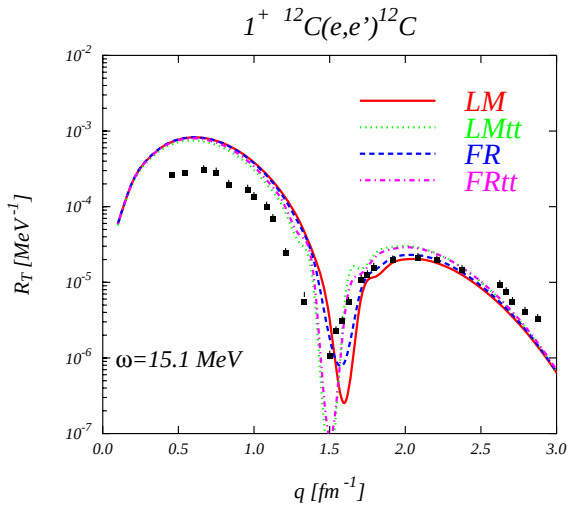






$^{12}\text{C} : 1^+$ Isoscalar

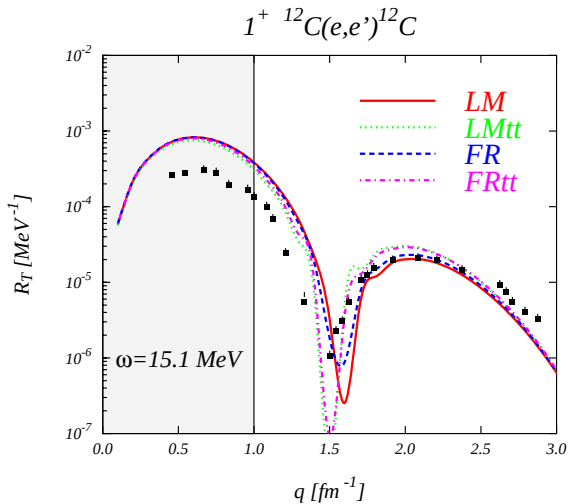


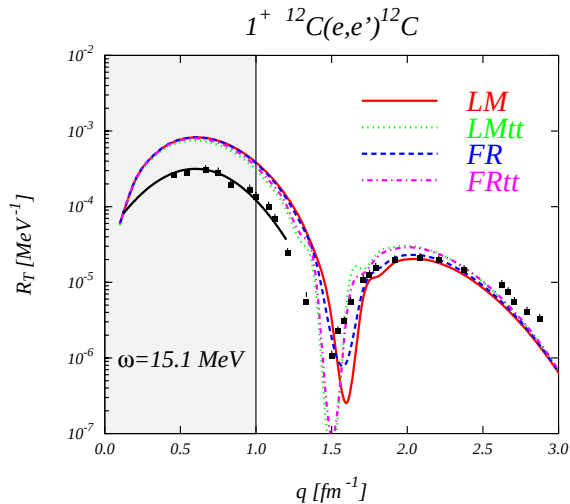


$$\left(\frac{d^2\sigma}{d\Omega d\epsilon} \right)_{e,e'} = \sigma_M \frac{1}{2J_i + 1} \left[|v_L| < J_f || C_J^Y || J_i >|^2 + |v_T| < J_f || T_J^Y || J_i >|^2 \right]$$

$$\begin{aligned} \left(\frac{d^2\sigma}{d\Omega d\epsilon} \right)_{\nu,\nu'} &= G^2 \frac{\epsilon'_\nu}{\epsilon_\nu} \frac{1}{2J_i + 1} \left[|w_C^A| < J_f || C_J^A || J_i >|^2 + \right. \\ &+ |w_L^A| < J_f || L_J^A || J_i >|^2 + |w_L^Y| < J_f || C_J^Y || J_i >|^2 + \\ &+ |w_{LC}^A| \operatorname{Re} (< J_f || C_J^A || J_i > < J_f || C_J^A || J_i >^*) \\ &+ |w_T^A| < J_f || T_J^A || J_i >|^2 + |w_T^Y| < J_f || T_J^Y || J_i >|^2 \\ &+ |w_T^{AV}| \operatorname{Re} (< J_f || T_J^A || J_i > < J_f || T_J^Y || J_i >^*) \left. \right] \end{aligned}$$

$^{12}\text{C} : 1^+$ Isovector

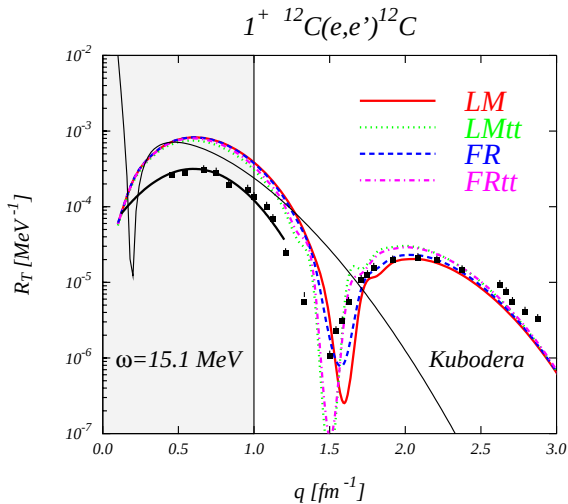




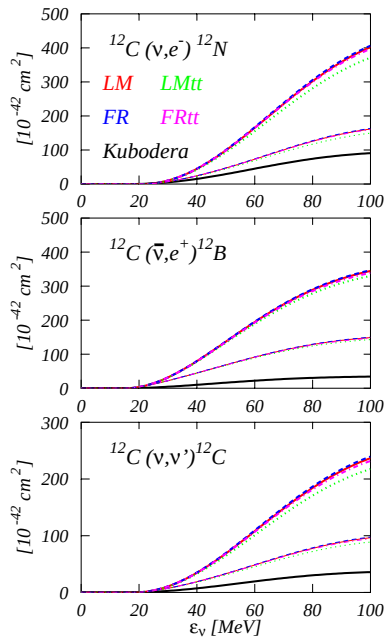
$$Q(q) = \frac{|\langle J_f || T_J^{\text{exp}}(q) || J_i \rangle|^2_{e,e'}}{|\langle J_f || T_J^V(q) || J_i \rangle|^2_{RPA}}$$

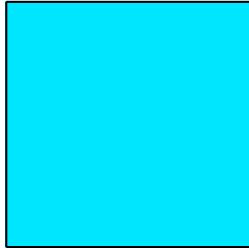
$$\left(\frac{d^2\sigma}{d\Omega d\epsilon} \right)_{\nu,\nu'}^{\text{scal}} = Q(q) \left(\frac{d^2\sigma}{d\Omega d\epsilon} \right)_{\nu,\nu'}^{\text{RPA}}$$

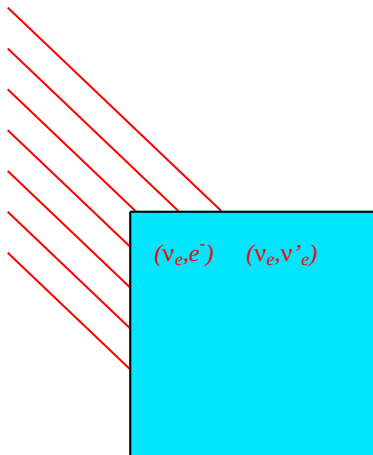
$^{12}\text{C} : 1^+$ Isovector



M. Fukugita, Y. Kohyama and K. Kubodera, Phys. Lett. B 212 (1988) 139.

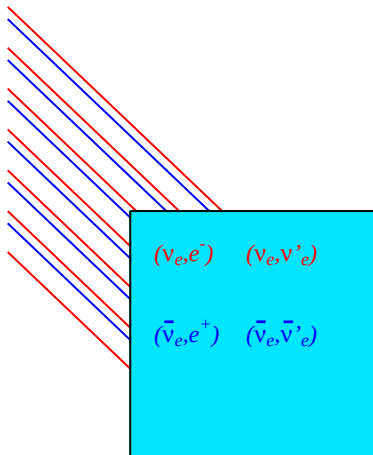






$$N_{nc} = (\nu_e, \nu'_e)$$

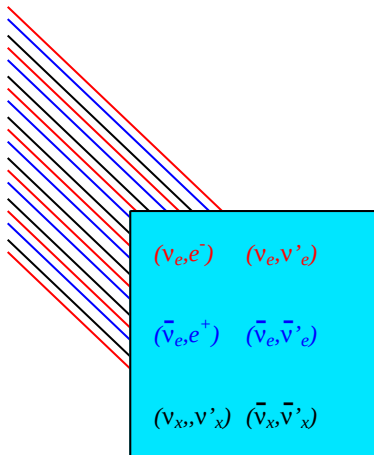
$$N_{np} = (\nu_e, e^-)$$



$$N_{nc} = (\nu_e, \nu'_e) + (\bar{\nu}_e, \bar{\nu}'_e)$$

$$N_{np} = (\nu_e, e^-)$$

$$N_{pn} = (\bar{\nu}_e, e^+)$$



$$N_{nc} = (\nu_e, \nu'_e) + (\bar{\nu}_e, \bar{\nu}'_e) + (\nu_x, \nu'_x) + (\bar{\nu}_x, \bar{\nu}'_x)$$

$$N_{np} = (\nu_e, e^-)$$

$$N_{pn} = (\bar{\nu}_e, e^+)$$

$$N_i = \frac{1}{4\pi} \frac{E_B f_i}{D^2} \frac{n_t \eta}{\langle \epsilon_\nu \rangle} \int_{e_{th}}^{\infty} f(\epsilon_\nu) \sigma(\epsilon_\nu) d\epsilon_\nu$$

Supernova parameters

$E_B \sim 5.0 \cdot 10^{52}$ erg = $8.01 \cdot 10^{60}$ MeV (Energy)

$f_i = 1/6$ (Fraction of the energy carried by ν_i)

$D = 10$ kpc (Distance from the earth)

LVD parameters [1]

$n_t = 4.23 \cdot 10^{31}$ (Number of ^{12}C target nuclei)

$\eta_{NC} = 0.55$, $\eta_{np}(\nu, e^-) = 0.75$, $\eta_{pn}(\bar{\nu}, e^+) = 0.90$ (efficiencies)

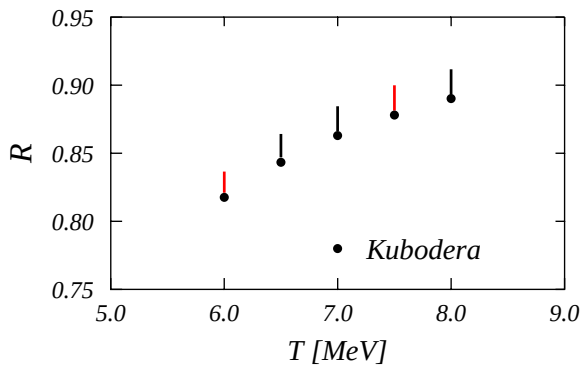
Number of events	$\nu + e^-$	$\bar{\nu} + e^+$	NC T=6.0 MeV	NC T=7.5 MeV
Kubodera	7	9	25	32
N_{min}	8	43	64	84
N_{max}	14	63	112	151

[1] N. Yu. Agafonova *et al.*, *Astrop. Phys.* 27 (2007) 254

$$N_{NC} = N_{NC}^T + N(\nu_e, \nu'_e)_4 + N(\bar{\nu}_e, \bar{\nu}'_e)_5$$

$$N_{NC}^T = N(\nu_x, \nu'_x)_T + N(\bar{\nu}_x, \bar{\nu}'_x)_T$$

$$R = \frac{N_{NC}^T}{N_{NC}} = \frac{N_{NC} - [N(\nu_e, \nu'_e)_4 + N(\bar{\nu}_e, \bar{\nu}'_e)_5]}{N_{NC}}$$

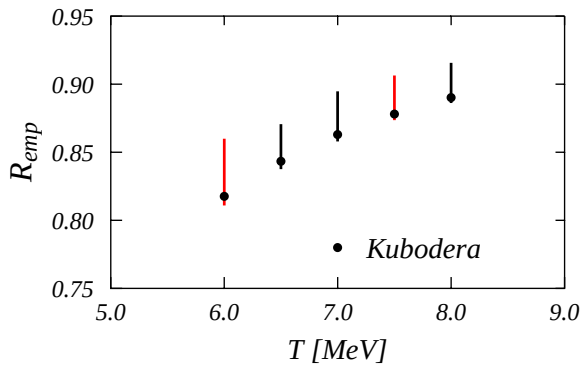


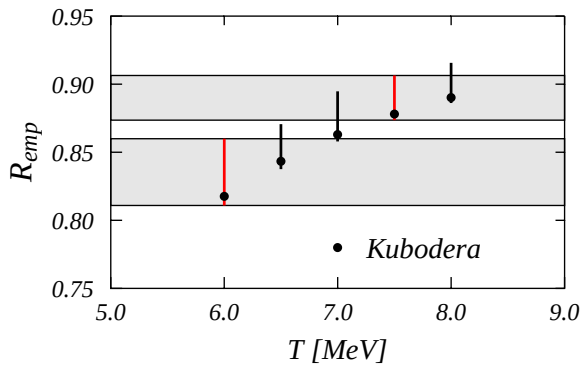
$$N_{np} = N(\nu_e, e^-)_4$$

$$N_{pn} = N(\bar{\nu}_e, e^+)_5$$

$$R_{emp} = \frac{1}{N_{NC}^{emp}} \left[N_{NC}^{emp} - N_{np}^{emp} \frac{\int_{e_{th}}^{\infty} f(\epsilon) \sigma_{\nu_e, \nu'_e}(\epsilon) d\epsilon}{\eta_{np} \int_{e_{th}}^{\infty} f(\epsilon) \sigma_{\nu_e, e^-}(\epsilon) d\epsilon} - N_{pn}^{emp} \frac{\int_{e_{th}}^{\infty} f(\epsilon) \sigma_{\bar{\nu}_e, \bar{\nu}'_e}(\epsilon) d\epsilon}{\eta_{pn} \int_{e_{th}}^{\infty} f(\epsilon) \sigma_{\bar{\nu}_e, e^+}(\epsilon) d\epsilon} \right]$$

In our simulation $N^{emp} = N(\text{Kubodera})$





Conclusions

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Discrete excitations better described.
Uncerainties on the global magnitude of the cross sections

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Nuclear structure uncertainties are relevant for precision ν astrometry

Uncertainties are large above the nucleon emission threshold (giant resonance region)

Discrete excitations better described
Uncertainties on the magnitude of the cross sections

Ratios of observables stable against uncertainties

μ capture rate

10^3 s^{-1}	LM	LMtt	FR	FRtt	exp [1]
no scal	34.98	33.38	35.14	34.30	
scal	14.10	13.35	14.14	13.76	6.2 ± 0.3

$\langle \sigma_e \rangle$ from μ decay at rest

10^{-42} cm^2	LM	LMtt	FR	FRtt	Kub.	exp [2]
no scal	26.55	25.22	26.85	26.64		
scal	14.24	13.55	14.40	14.30	9.19	8.9 ± 0.12

[1] D.F. Measday, Phys. Rep. 354 (1001) 243

[2] L.B. Auerbach *et al.* (LSND), Phys. Rev. C 64 (2001) 065501

$^{208}\text{Pb} (\nu, \nu') ^{208}\text{Pb}$

