# Radiation-matter interaction

- 1. General details for charged particles
- 2. Heavy charged particles
- 3. Electrons and positrons
- 4. Photons
- 5. Neutrons

6. Simulating particle transport in matter

1. General details for charged particles

#### • problem:

- -assume a particle that collides with an atomic electron
- -the particle moves rapidly with respect to the electron
- -the energy transferred is large compared to the binding energy

\*the electron is free and at rest

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 $\frac{\text{relativistic expression}}{Q_{\text{max}}} = \frac{2\gamma^2 m V_{\text{i}}^2}{1 + 2\gamma m/M + m^2/M^2}$  $\gamma = (1 - \beta^2)^{-1/2}; \quad \beta = V_{\text{i}}/c$ 

1. General details for charged particles

 -electrons may loss a large fraction of energy in a collision: tortuous paths in matter
 -heavy charged particles loss small fractions of energy:

straight paths in matter

energy \*<u>the electron is free and at rest</u>

-to calculate the maximum energy transfer
-head-on collision

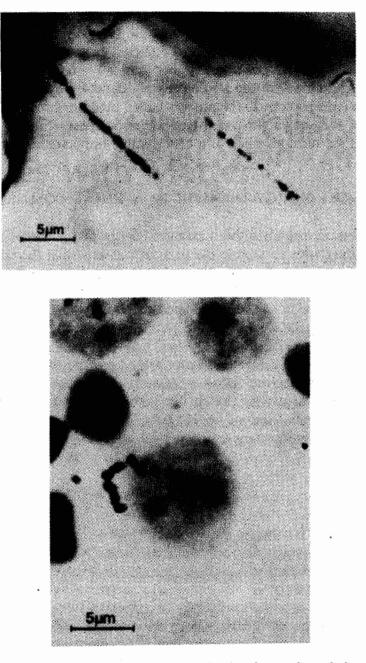
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General details for charged particles
 -electrons:

tortuous paths in matter -heavy charged particles: straight paths in matter General details for charged particles
 -electrons:

tortuous paths in matter -heavy charged particles: straight paths in matter



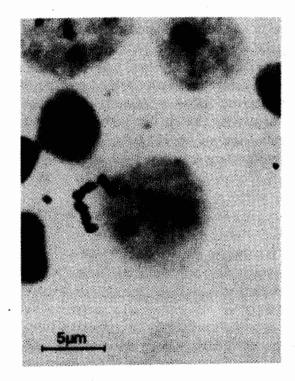
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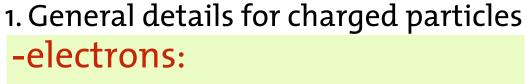
# tortuous paths in matter -heavy charged particles: straight paths in matter

Proton Kinetic Energy E (MeV)	Q <sub>max</sub> (MeV)	Maximum Percentage Energy Transfer 100Q <sub>max</sub> /E
0.1	0.00022	0.22
1	0.0022	0.22
10	0.0219	0.22
100	0.229	0.23
$10^{3}$	3.33	0.33
$10^{4}$	136	1.4
10 <sup>5</sup>	$1.06 \ge 10^4$	10.6
10 <sup>6</sup>	$5.38 \times 10^5$	53.8
10 <sup>7</sup>	9.21 x 10 <sup>6</sup>	92.1

Maximum Possible Energy Transfer, Q<sub>max</sub>, in Proton Collision with Electron







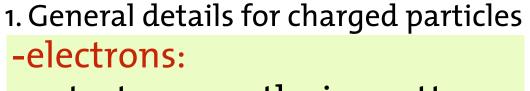
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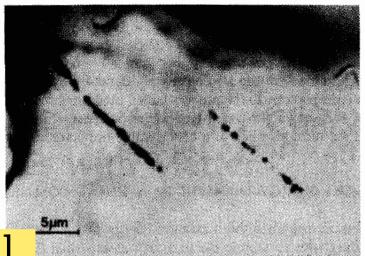


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# tortuous paths in matter -heavy charged particles: straight paths in matter



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(MeV)	energy loss small		
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**Proton Kinetic** Energy E (MeV)

0.1

10

100  $10^{3}$ 

 $10^{4}$ 

 $10^{5}$ 

10<sup>6</sup>

 $10^{7}$ 

Maximum Possible Energy Transfer, Omer. in Proton Collision with Electron

0.00022

0.0022 0.0219

0.229

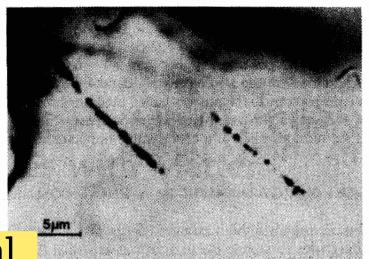
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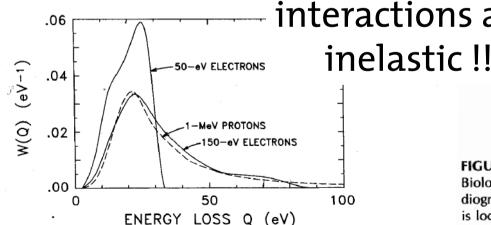
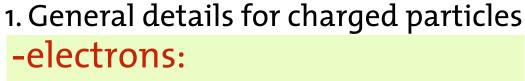
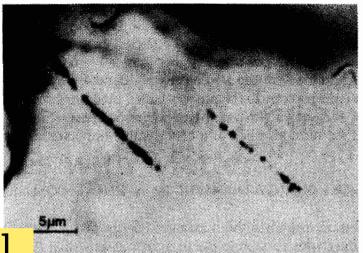


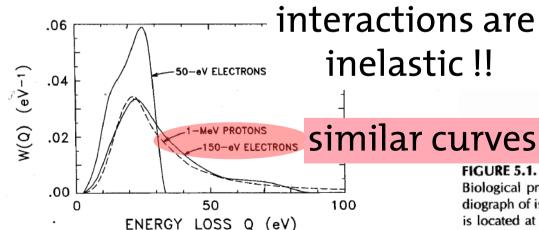
FIGURE 5.3. Single-collision energy-loss spectra for 50-eV and 150-eV electrons and 1-MeV protons in liquid water. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)

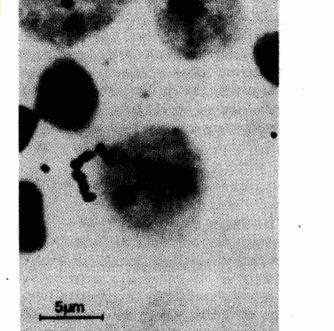


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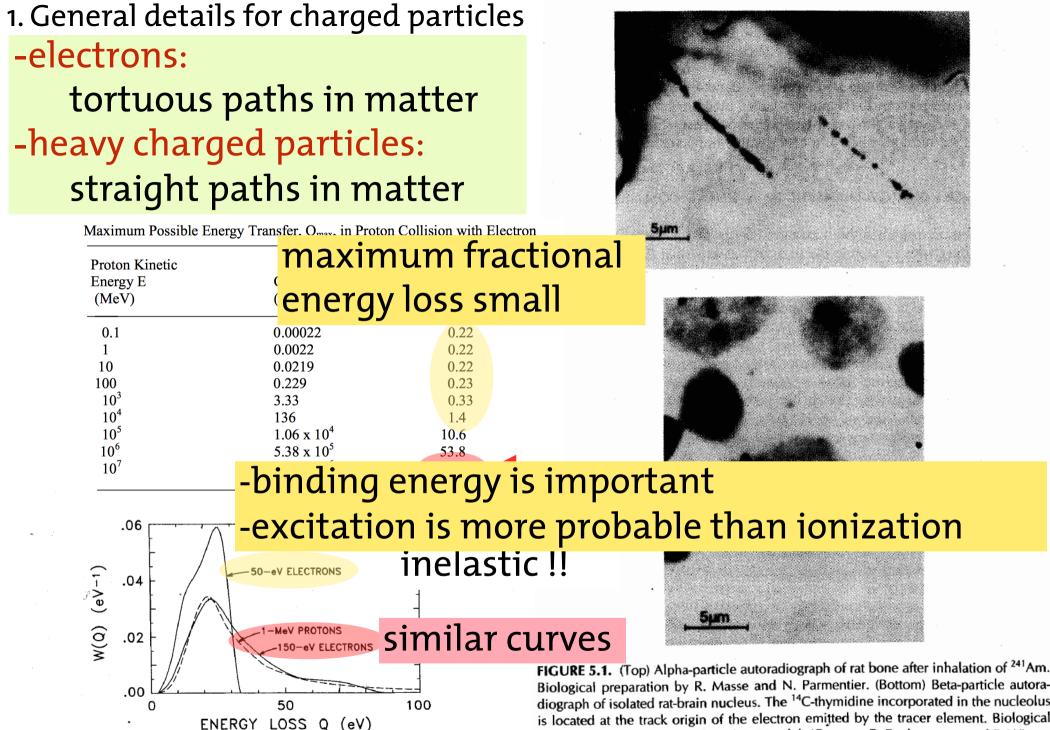


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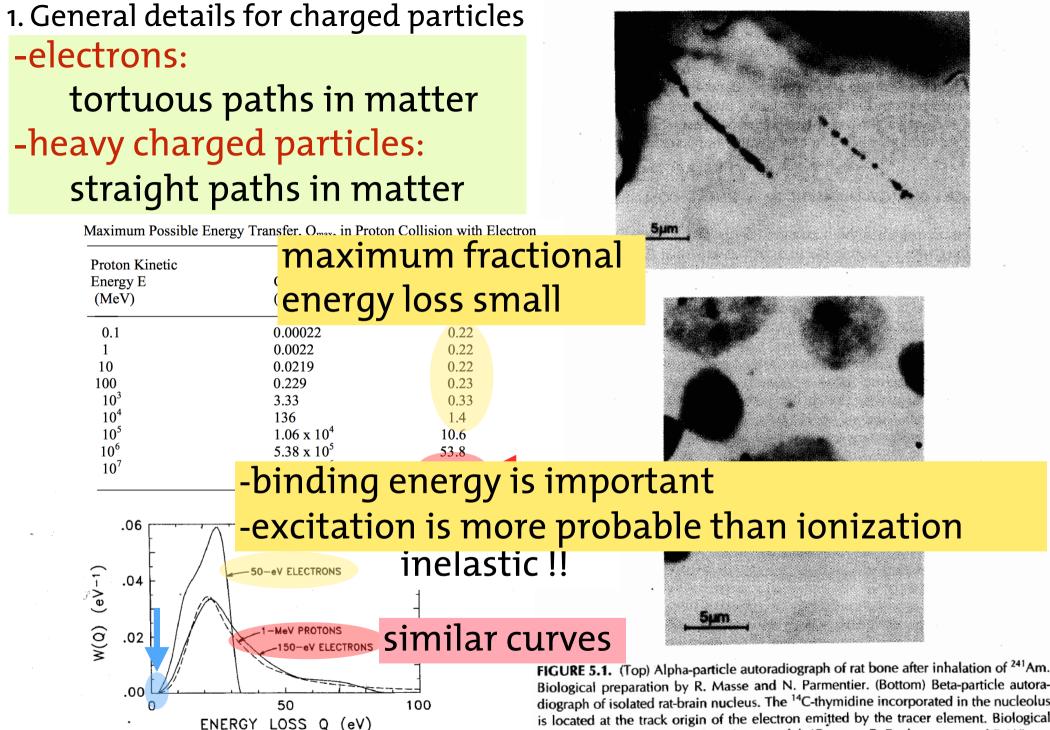


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#### stopping power

- average energy loss of a charged particle per unit path length

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# • Bethe's formula:

$$-\frac{\mathrm{d}E}{\mathrm{d}x} = \frac{4\pi k_0^2 z^2 e^4 n}{mc^2 \beta^2} \left[ \ln \frac{2mc^2 \beta^2}{I(1-\beta^2)} - \beta^2 \right]$$

- $k_{o} = 8.99 \text{ x } 10^9 \text{ N } \text{m}^2 \text{ C}^{-2}$ , (the Boltzman constant)
- z = atomic number of the heavy particle,
- e = magnitude of the electron charge,
- n = number of electrons per unit volume in the medium,
- m = electron rest mass,
- c = speed of light in vacuum,
- $\beta = V/c$  = speed of the particle relative to c,
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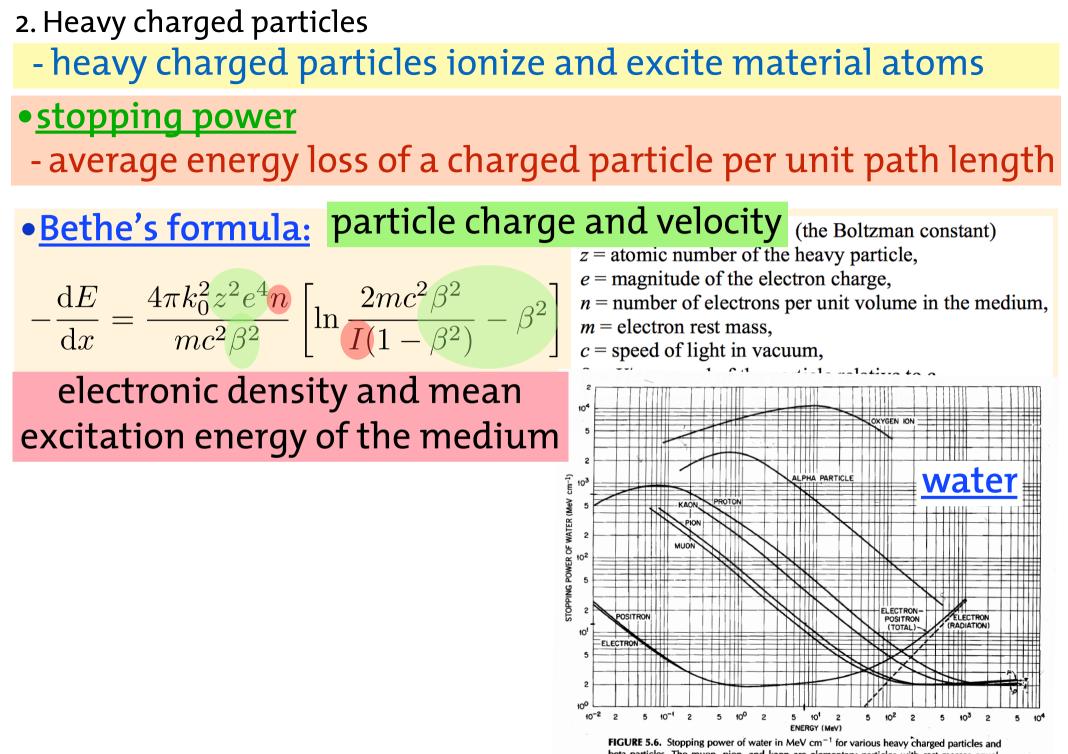
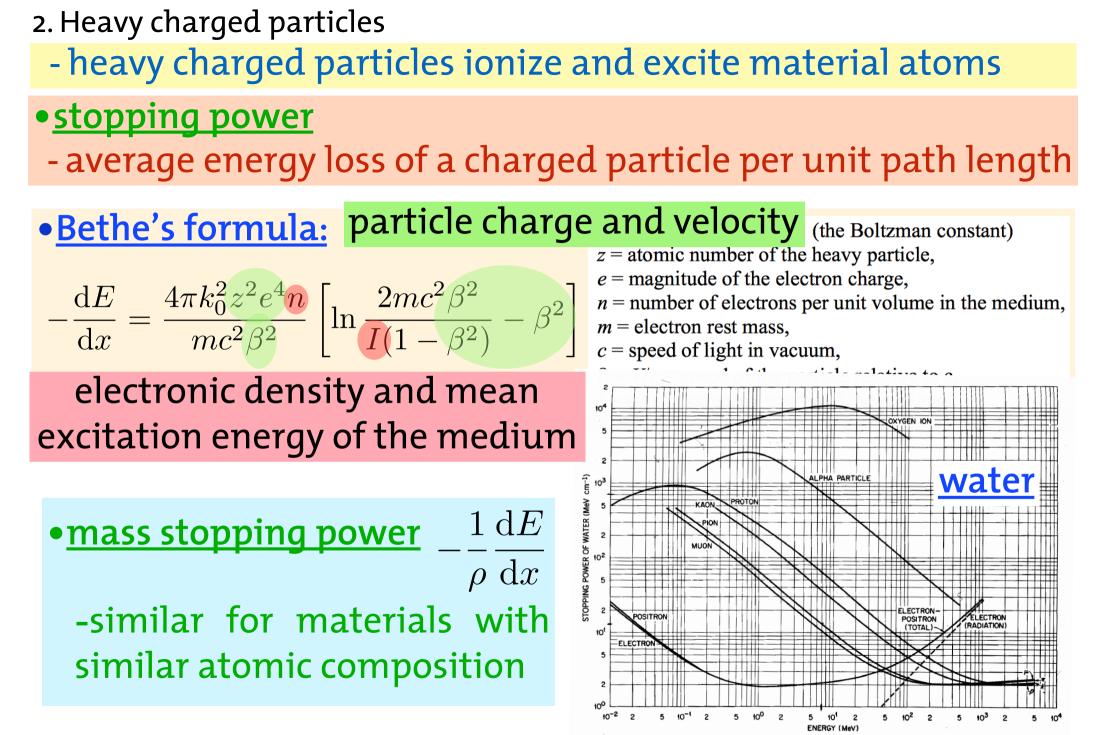


FIGURE 5.6. Stopping power of water in MeV cm<sup>-1</sup> for various heavy charged particles and beta particles. The muon, pion, and kaon are elementary particles with rest masses equal, respectively, to about 207, 270, and 967 electron rest masses. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)



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TABLE 5.3.	Mass Stopping Powe	$r = dE/\rho dx$ and Range $R_P$ for Protection	ons in Water
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	stopping rower – uzip	ux and kange Rp for Froton	s in water
Kinetic Energy		$-dE/\rho dx$	R <sub>p</sub>
(MeV)	$\beta^2$	$(MeV cm^2 g^{-1})$	(g cm <sup>-2</sup> )
0.01	.000021	500.	$3 \times 10^{-5}$
0.04	.000085	860.	$6 \times 10^{-5}$
0.05	.000107	910.	$7 \times 10^{-5}$
0.08	.000171	920.	$9 \times 10^{-5}$
0.10	.000213	910.	$1 \times 10^{-4}$
0.50	.001065	428.	$8 \times 10^{-4}$
1.00	.002129	270.	0.002
2.00	.004252	162.	0.007
4.00	.008476	95.4	0.023
6.00	.01267	69.3	0.047
8.00	.01685	55.0	0.079
10.0	.02099	45.9	0.118
12.0	.02511	39.5	0.168
14.0	.02920	34.9	0.217
16.0	.03327	31.3	0.280
18.0	.03731	28.5	0.342
20.0	.04133	26.1	0.418
25.0	.05126	21.8	0.623
30.0	.06104	18.7	0.864
35.0	.07066	16.5	1.14
40.0	.08014	14.9	1.46
45.0	.08948	13.5	1.80
50.0	.09867	12.4	2.18
60.0	.1166	10.8	3.03
70.0	.1341	9.55	4.00
80.0	.1510	8.62	5.08
90.0	.1675	7.88	6.27
100.	.1834	7.28	7.57
150.	.2568	5.44	15.5
200.	.3207	4.49	25.5
300.	.4260	3.52	50.6
400.	.5086	3.02	80.9
500.	.5746	2.74	115.
600.	.6281	2.55	152.
700.	.6721	2.42	192.
800.	.7088	2.33	234.
900.	.7396	2.26	277.
1000.	.7658	2.21	321.
2000.	.8981	2.05	795.
4000.	.9639	2.09	1780.

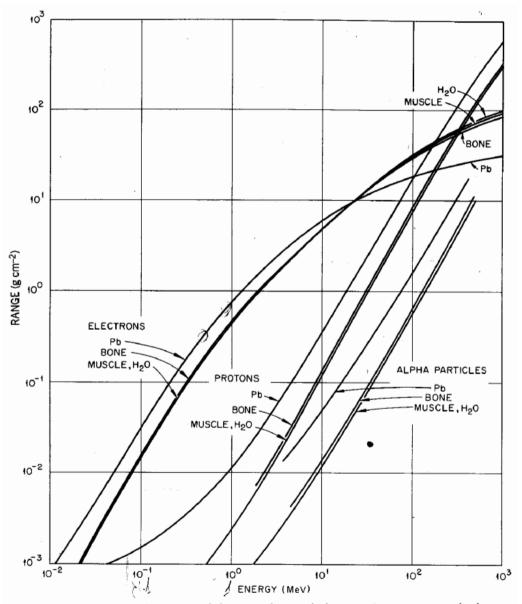
#### • range

- distance traveled by the particle before coming to rest  $R(K) = \int_0^K dE \left(-\frac{dE}{dx}\right)^{-1}$ 

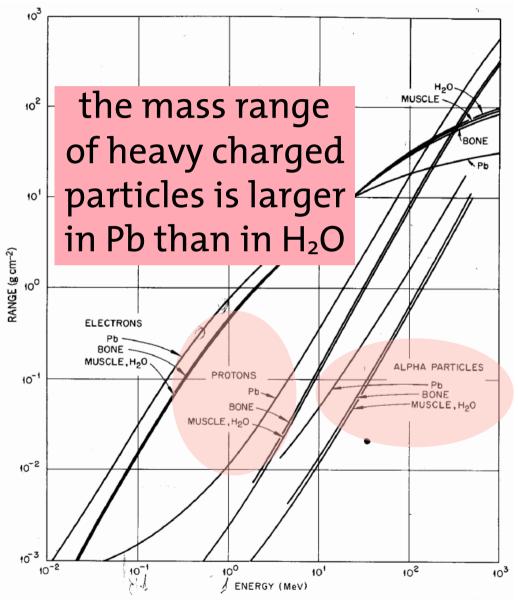
-ranges for two heavy particles with the same initial velocity  $R_1(\beta) = \frac{z_2^2 M_1}{z_1^2 M_2} R_2(\beta)$ 

TABLE 5.3. Mass Stopping Power  $- dE/\rho dx$  and Range  $R_P$  for Protons in Water

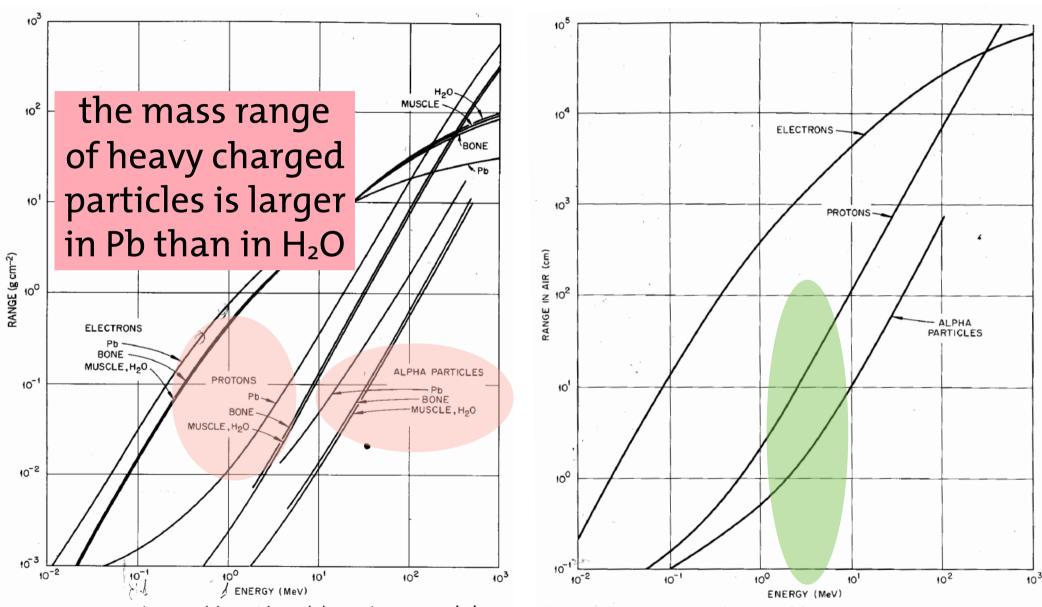
Kinetic Energy (MeV) $-dE/\rho dx$ (MeV cm² g <sup>-1</sup> )0.01.000021500.0.04.000085860.0.05.000107910.0.08.000171920.0.10.000213910.0.50.001065428.1.00.002129270.2.00.004252162.4.00.00847695.46.00.0126769.38.00.0168555.010.0.0299945.912.0.0251139.514.0.0292034.916.0.0332731.318.0.0373128.520.0.0413326.125.0.0512621.830.0.0610418.735.0.0706616.540.0.0801414.945.0.0894813.550.0.0986712.460.0.116610.870.0.13419.5580.0.15108.6290.0.16757.8810018347.2815025685.44	$\frac{R_{\rm p}}{({\rm g \ cm^{-2}})}$ $3 \times 10^{-5}$
	(g cm <sup>-2</sup> )
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$3 \times 10^{-5}$
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$6 \times 10^{-5}$
	$7 \times 10^{-5}$
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$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$1 \times 10^{-4}$
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$\begin{array}{cccccccccccccccccccccccccccccccccccc$	0.002
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50.0.0986712.460.0.116610.870.0.13419.5580.0.15108.6290.0.16757.8810018347.28	1.46
60.0.116610.870.0.13419.5580.0.15108.6290.0.16757.8810018347.28	1.80
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80.0.15108.6290.0.16757.8810018347.28	3.03
90.0.16757.8810018347.28	4.00
1001834 7.28	5.08
	6.27
1502568 5.44	7.57
	15.5
2003207 4.49	25.5
3004260 3.52	50.6
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**FIGURE 5.7.** Ranges of protons, alpha particles, and electrons in water, muscle, bone, and lead, expressed in g cm<sup>-2</sup>. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)



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**FIGURE 5.8.** Ranges in cm of protons, alpha particles, and electrons in air at STP. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)

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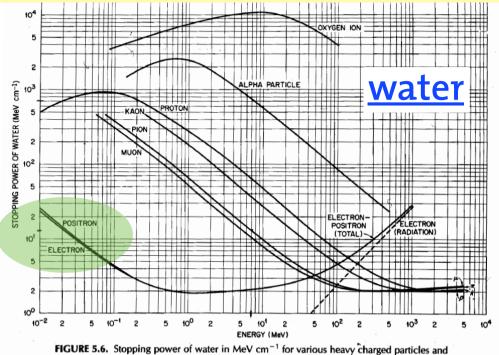
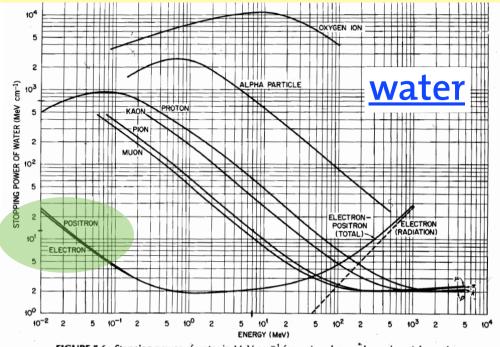


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# • stopping power

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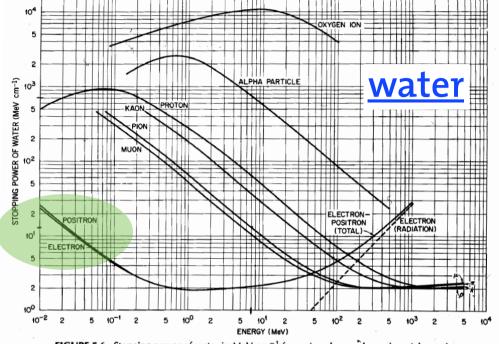


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$$\left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{tot}}^{\pm} = \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{col}}^{\pm} + \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{rad}}^{\pm}$$

# - electrons and positrons ionize and excite material atoms



**FIGURE 5.6.** Stopping power of water in MeV cm<sup>-1</sup> for various heavy charged particles and beta particles. The muon, pion, and kaon are elementary particles with rest masses equal, respectively, to about 207, 270, and 967 electron rest masses. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)

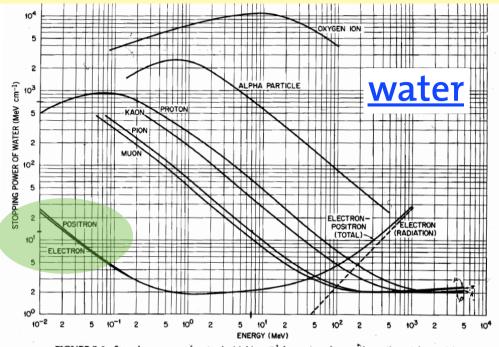
# • <u>collisional stopping power</u> $\left( -\frac{\mathrm{d}E}{\mathrm{d}x} \right)_{\mathrm{col}}^{\pm} = \frac{4\pi k_0^2 z^2 e^4 n}{mc^2 \beta^2} \left[ \ln \frac{mc^2 \tau \sqrt{\tau+2}}{\sqrt{2}I} + F^{\pm}(\beta) \right], \quad \tau = \frac{K}{mc^2}$ $F^{-}(\beta) = \frac{1-\beta^2}{2} \left[ 1 + \frac{\tau^2}{8} - (2\tau+1)\ln 2 \right], \quad F^{+}(\beta) = \ln 2 - \frac{\beta^2}{24} \left[ 23 + \frac{14}{\tau+2} + \frac{10}{(\tau+2)^2} + \frac{4}{(\tau+2)^3} \right]$

## • <u>stopping power</u>

$$\left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{tot}}^{\pm} = \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{col}}^{\pm} + \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)$$

Kinetic Energy	β²	$-\frac{1}{\rho} \left(\frac{dE}{dx}\right)_{col}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	$-\frac{1}{\rho} \left(\frac{dE}{dx}\right)^{-}_{rad}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	$-\frac{1}{\rho} \left( \frac{dE}{dx} \right)_{tot}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	Radiation Yield
10 eV	0.00004	4.0	_	4.0	
30	0.00012	44.	_	44.	
50	0.00020	170.	_	170.	
75	0.00029	272.	_	272.	
100	0.00039	314.	·	314.	
200	0.00078	298.		298.	_
500 eV	0.00195	194.	_	194.	
1 keV	0.00390	126.		126.	_
2	0.00778	77.5	-	77.5	` <del>_</del>
5	0.0193	42.6	—	42.6	_
10	0.0380	23.2	-	23.2	0.0001
25	0.0911	11.4		11.4	0.0002
50	0.170	6.75	_	6.75	0.0004
75	0.239	5.08	_	5.08	0.0006
100	0.301	4.20	—	4.20	0.0007
200	0.483	2.84	0.006	2.85	0.0012
500	0.745	2.06	0.010	2.07	0.0026
700 keV	0.822	1.94	0.013	1.95	0.0036
1 MeV	0.886	1.87	0.017	1.89	0.0049
4	0.987	1.91	0.065	1.98	0.0168
7	0.991	1.93	0.084	2.02	0.0208
10	0.998	2.00	0.183	2.18	0.0416
100	0.999+	2.20	2.40	4.60 🔹	0.317
1000 MeV	0.999+	2.40	26.3	28.7	0.774

# - electrons and positrons ionize and excite material atoms



**FIGURE 5.6.** Stopping power of water in MeV cm<sup>-1</sup> for various heavy charged particles and beta particles. The muon, pion, and kaon are elementary particles with rest masses equal, respectively, to about 207, 270, and 967 electron rest masses. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)

## <u>collisional stopping power</u>

$$\left( -\frac{\mathrm{d}E}{\mathrm{d}x} \right)_{\mathrm{col}}^{\pm} = \frac{4\pi k_0^2 z^2 e^4 n}{mc^2 \beta^2} \left[ \ln \frac{mc^2 \tau \sqrt{\tau+2}}{\sqrt{2}I} + F^{\pm}(\beta) \right], \quad \tau = \frac{K}{mc^2}$$

$$F^{-}(\beta) = \frac{1-\beta^2}{2} \left[ 1 + \frac{\tau^2}{8} - (2\tau+1)\ln 2 \right], \quad F^{+}(\beta) = \ln 2 - \frac{\beta^2}{24} \left[ 23 + \frac{14}{\tau+2} + \frac{10}{(\tau+2)^2} + \frac{4}{(\tau+2)^3} \right]$$

 $\pm$ 

rad

## stopping power

$$\left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{tot}}^{\pm} = \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{col}}^{\pm} + \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)$$

(	$\mathrm{d}E$	$\rangle^{\pm}$
	dx	$I_{\rm rad}$

.

Kinetic		$-\frac{1}{\rho}\left(\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{col}}^{-}$	$-\frac{1}{\rho}\left(\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{rad}}$	$-\frac{1}{\rho}\left(\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{tot}}^{-}$	Radiation
Energy	β²	(MeV cm <sup>2</sup> g <sup>-1</sup> )	(MeV cm <sup>2</sup> g <sup><math>-1</math></sup> )	(MeV cm <sup>2</sup> g <sup>-1</sup> )	Yield
10 eV	0.00004	4.0	_	4.0	
30	0.00012	44.	_	44.	
50	0.00020	170.	_	170.	
75	0.00029	272.	_	272.	
100	0.00039	314.	·	314.	
200	0.00078	298.		298.	_
500 eV	0.00195	194.	_	194.	
1 keV	0.00390	126.		126.	_
2	0.00778	77.5	_	77.5	` —
5	0.0193	42.6		42.6	_
10	0.0380	23.2		23.2	0.0001
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## collisional stopping power

# - electrons and positrons ionize and excite material atoms

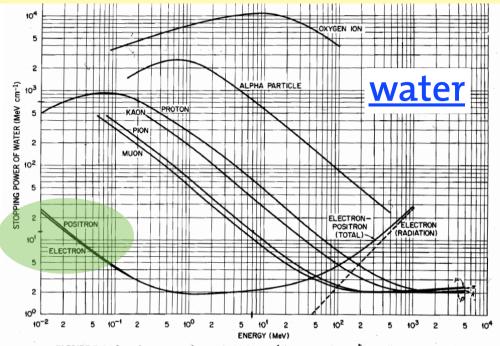
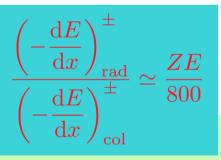


FIGURE 5.6. Stopping power of water in MeV cm<sup>-1</sup> for various heavy charged particles and beta particles. The muon, pion, and kaon are elementary particles with rest masses equal, respectively, to about 207, 270, and 967 electron rest masses. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)



$$\left( -\frac{\mathrm{d}E}{\mathrm{d}x} \right)_{\mathrm{col}}^{\pm} = \frac{4\pi k_0^2 z^2 e^4 n}{mc^2 \beta^2} \left[ \ln \frac{mc^2 \tau \sqrt{\tau+2}}{\sqrt{2}I} + F^{\pm}(\beta) \right], \quad \tau = \frac{K}{mc^2}$$

$$F^{-}(\beta) = \frac{1-\beta^2}{2} \left[ 1 + \frac{\tau^2}{8} - (2\tau+1)\ln 2 \right], \quad F^{+}(\beta) = \ln 2 - \frac{\beta^2}{24} \left[ 23 + \frac{14}{\tau+2} + \frac{10}{(\tau+2)^2} + \frac{4}{(\tau+2)^3} \right]$$

## stopping power

$$\left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{tot}}^{\pm} = \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{col}}^{\pm} + \left(-\frac{\mathrm{d}E}{\mathrm{d}x}\right)$$

Kinetic Energy	β²	$-\frac{1}{\rho} \left(\frac{dE}{dx}\right)_{col}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	$-\frac{1}{\rho} \left(\frac{dE}{dx}\right)_{rad}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	$-\frac{1}{\rho} \left(\frac{dE}{dx}\right)_{tot}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	Radiation Yield
10 eV	0.00004	4.0		4.0	_
30	0.00012	44.	_	44.	
50	0.00020	170.	_	170.	
75	0.00029	272.	_	272.	
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25	0.0911	11.4	-	11.4	0.0002
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75	0.239	5.08	_	5.08	0.0006
100	0.301	4.20	_	4.20	0.0007
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## collisional stopping power

# - electrons and positrons ionize and excite material atoms

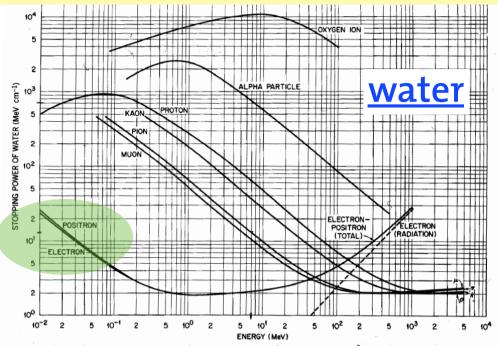
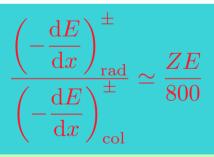


FIGURE 5.6. Stopping power of water in MeV cm<sup>-1</sup> for various heavy charged particles and beta particles. The muon, pion, and kaon are elementary particles with rest masses equal, respectively, to about 207, 270, and 967 electron rest masses. (Courtesy Oak Ridge National Laboratory, operated by Martin Marietta Energy Systems, Inc., for the Department of Energy.)



$$\frac{\text{radiation yield}}{Y \simeq \frac{6 \cdot 10^{-4} Z T}{1 + 6 \cdot 10^{-4} Z T}}$$

$$\left( -\frac{\mathrm{d}E}{\mathrm{d}x} \right)_{\mathrm{col}}^{\pm} = \frac{4\pi k_0^2 z^2 e^4 n}{mc^2 \beta^2} \left[ \ln \frac{mc^2 \tau \sqrt{\tau+2}}{\sqrt{2}I} + F^{\pm}(\beta) \right], \quad \tau = \frac{K}{mc^2}$$

$$F^{-}(\beta) = \frac{1-\beta^2}{2} \left[ 1 + \frac{\tau^2}{8} - (2\tau+1)\ln 2 \right], \quad F^{+}(\beta) = \ln 2 - \frac{\beta^2}{24} \left[ 23 + \frac{14}{\tau+2} + \frac{10}{(\tau+2)^2} + \frac{4}{(\tau+2)^3} \right]$$

 $\pm$ 

rad

#### • range

- distance traveled by the particle before coming to rest  $R(K) = \int_0^K dE \left(-\frac{dE}{dx}\right)^{-1}$ 

#### • range

- distance traveled by the particle before coming to rest  $R(K) = \int_0^K dE \left(-\frac{dE}{dx}\right)^{-1}$ 

TABLE 6.1. Electron Collisonal, Radiative, and Total Mass Stopping Powers; Radiation Yield; and Range in Water

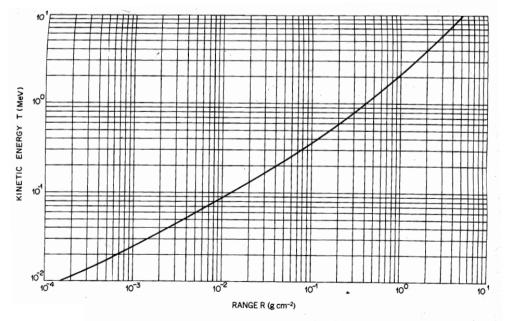
anu kange	III Water					
Kinetic Energy	β²	$-\frac{1}{\rho} \left(\frac{dE}{dx}\right)_{col}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	$-\frac{1}{\rho} \left( \frac{\mathrm{d}E}{\mathrm{d}x} \right)_{\mathrm{rad}}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	$-\frac{1}{\rho} \left(\frac{dE}{dx}\right)_{tot}^{-}$ (MeV cm <sup>2</sup> g <sup>-1</sup> )	Radiation Yield	Range (g cm <sup>-2</sup> )
10 eV	0.00004	4.0		4.0	_	4 × 10 <sup></sup>
30	0.00012	44.	_	44.		$2 \times 10^{-1}$
50	0.00020	170.	_	170.		$3 \times 10^{-1}$
75	0.00029	272.	_	272.		$4 \times 10^{-1}$
100	0.00039	314.		314.		$5 \times 10^{-1}$
200	0.00078	298.		298.	_	$8 \times 10^{-1}$
500 eV	0.00195	194.	_	1 <del>9</del> 4.		$2 \times 10^{-1}$
1 keV	0.00390	126.		126.	_	$5 \times 10^{-1}$
2	0.00778	77.5	_	77.5	· _	$2 \times 10^{-1}$
2 5	0.0193	42.6		42.6	_	$8 \times 10^{-1}$
10	0.0380	23.2	-	23.2	0.0001	0.000
25	0.0911	11.4		11.4	0.0002	0.001
50	0.170	6.75	_	6.75	0.0004	0.004
75	0.239	5.08	_	5.08	0.0006	0.008
100	0.301	4.20	_	4.20	0.0007	0.014
200	0.483	2.84	0.006	2.85	0.0012	0.044
500	0.745	2.06	0.010	2.07	0.0026	0.174
700 keV	0.822	1.94	0.013	1.95	0.0036	0.275
1 MeV	0.886	1.87	0.017	1.89	0.0049	0.430
4	0.987	1.91	0.065	1.98	0.0168	2.00
7	0.991	1.93	0.084	2.02	0.0208	2.50
10	0.998	2.00	0.183	2.18	0.0416	4.88
100	0.999+	2.20	2.40	4.60 🔹	0.317	32.5
1000 MeV	0.999+	2.40	26.3	28.7	0.774	101.

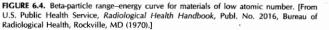
#### • range

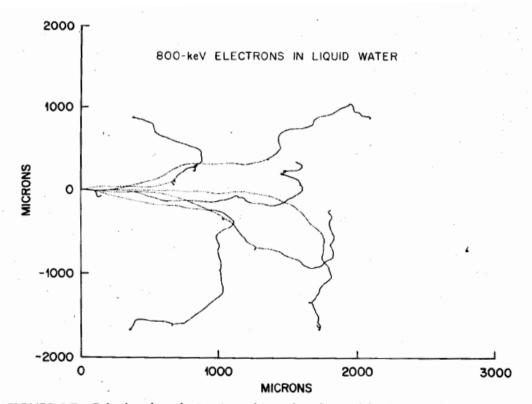
- distance traveled by the particle before coming to rest  $R(K) = \int_0^K dE \left(-\frac{dE}{dx}\right)^{-1}$ 

TABLE 6.1. Electron Collisonal, Radiative, and Total Mass Stopping Powers; Radiation Yield; and Range in Water

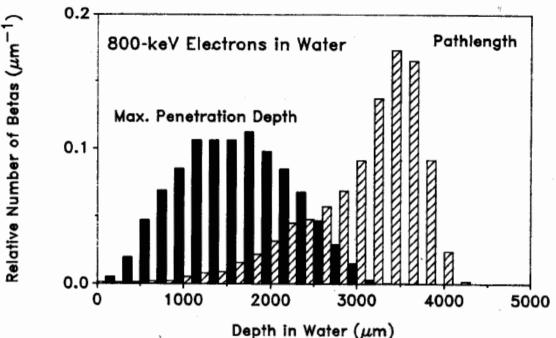
Kinetic		$-\frac{1}{\rho}\left(\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{col}}^{-}$	$-\frac{1}{\rho}\left(\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{rad}}$	$-\frac{1}{\rho}\left(\frac{\mathrm{d}E}{\mathrm{d}x}\right)_{\mathrm{tot}}^{-}$	Radiation	Range
Energy	β²	(MeV cm2 g-1)	$(MeV cm^2 g^{-1})$	(MeV cm <sup>2</sup> g <sup>-1</sup> )	Yield	(g cm <sup>-2</sup> )
10 eV	0.00004	4.0	_	4.0	_	4 × 10 <sup></sup>
30	0.00012	44.	_	44.		$2 \times 10^{-1}$
50	0.00020	170.	_	170.		$3 \times 10^{-1}$
75	0.00029	272.	_	272.		$4 \times 10^{-1}$
100	0.00039	314.		314.		$5 \times 10^{-1}$
200	0.00078	298.		298.	_	$8 \times 10^{-1}$
500 eV	0.00195	194.	-	194.		$2 \times 10^{-1}$
1 keV	0.00390	126.		126.	_	$5 \times 10^{-1}$
2 5	0.00778	77.5	_	77.5	· —	$2 \times 10^{-1}$
	0.0193	42.6		42.6	_	$8 \times 10^{-1}$
10	0.0380	23.2		23.2	0.0001	0.0002
25	0.0911	11.4		11.4	0.0002	0.0012
50	0.170	6.75	_	6.75	0.0004	0.0042
75	0.239	5.08	_	5.08	0.0006	0.008
100	0.301	4.20	_	4.20	0.0007	0.0140
200	0.483	2.84	0.006	2.85	0.0012	0.0440
500	0.745	2.06	0.010	2.07	0.0026	0.174
700 keV	0.822	1.94	0.013	1.95	0.0036	0.275
1 MeV	0.886	1.87	0.017	1.89	0.0049	0.430
4	0.987	1.91	0.065	1.98	0.0168	2.00
7	0.991	1.93	0.084	2.02	0.0208	2.50
10	0.998	2.00	0.183	2.18	0.0416	4.88
100	0.999+	2.20	2.40	4.60 🔹	0.317	32.5
1000 MeV	0.999+	2.40	26.3	28.7	0.774	101.

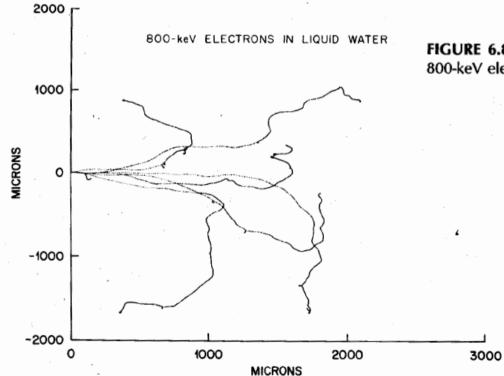






**FIGURE 6.7.** Calculated tracks (projected into the plane of the figure) of 800-keV electrons in water. Each electron starts moving horizontally toward the right from the point 0 on the vertical axis.



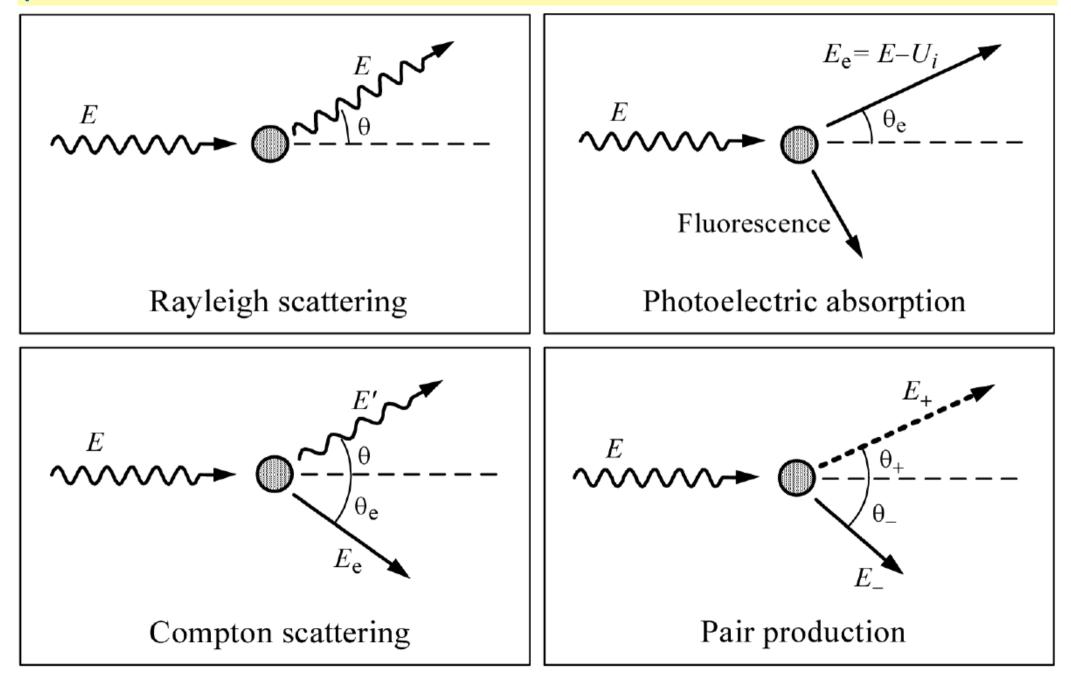


**FIGURE 6.8.** Distributions in pathlength and in maximum depth of penetration for 800-keV electrons in water.

**FIGURE 6.7.** Calculated tracks (projected into the plane of the figure) of 800-keV electrons in water. Each electron starts moving horizontally toward the right from the point 0 on the vertical axis.

# photon interactions

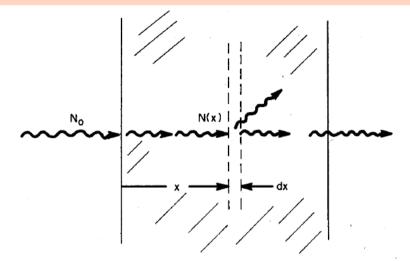
## photon interactions



- linear attenuation coefficient
  - probability per unit distance travelled that a photon interacts by any physical process

## linear attenuation coefficient

 probability per unit distance travelled that a photon interacts by any physical process

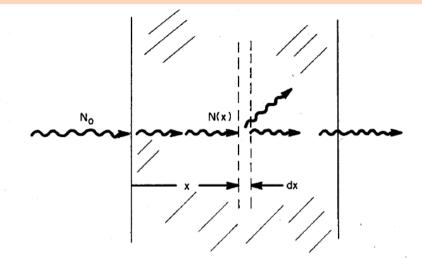


**FIGURE 8.6.** Pencil beam of  $N_0$  monoenergetic photons incident on slab. The number of photons that reach a depth x without having an interaction is given by  $N(x) = N_0 e^{-\mu x}$ , where  $\mu$  is the linear attenuation coefficient.

## linear attenuation coefficient

 probability per unit distance travelled that a photon interacts by any physical process

 $dN(x) = -\mu N(x) dx$  $\Downarrow$  $N(x) = N_0 \exp(-\mu x)$ 

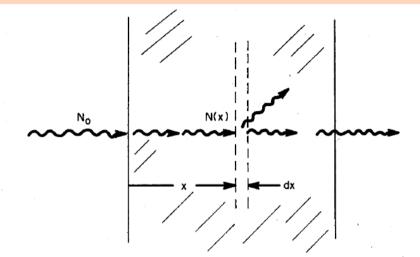


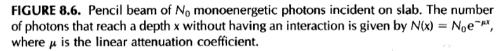
**FIGURE 8.6.** Pencil beam of  $N_0$  monoenergetic photons incident on slab. The number of photons that reach a depth x without having an interaction is given by  $N(x) = N_0 e^{-\mu x}$ , where  $\mu$  is the linear attenuation coefficient.

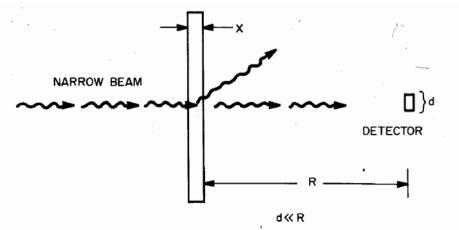
## linear attenuation coefficient

 probability per unit distance travelled that a photon interacts by any physical process

$$dN(x) = -\mu N(x) dx$$
$$\Downarrow$$
$$N(x) = N_0 \exp(-\mu x)$$





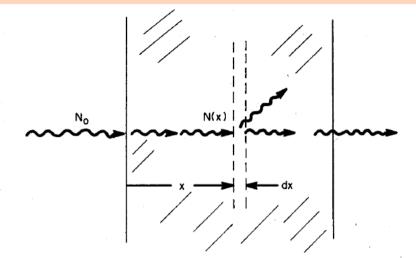


**FIGURE 8.7.** Illustration of "good" scattering geometry for measuring linear attenuation coefficient  $\mu$ . Photons from a *narrow* beam that are absorbed or scattered by the absorber do not reach a small detector placed in beam line some distance away.

## linear attenuation coefficient

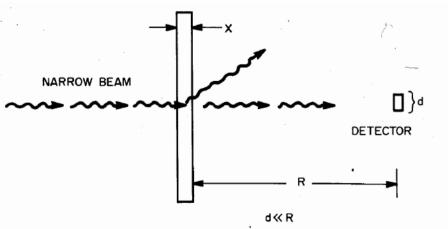
 probability per unit distance travelled that a photon interacts by any physical process

$$dN(x) = -\mu N(x) dx$$
$$\Downarrow$$
$$N(x) = N_0 \exp(-\mu x)$$



 $\exp(-\mu x) = N(x)/N_0$ probability that a normally incident photon will traverse a slab of thickness *x* without interacting

**FIGURE 8.6.** Pencil beam of  $N_0$  monoenergetic photons incident on slab. The number of photons that reach a depth x without having an interaction is given by  $N(x) = N_0 e^{-\mu x}$ , where  $\mu$  is the linear attenuation coefficient.



**FIGURE 8.7.** Illustration of "good" scattering geometry for measuring linear attenuation coefficient  $\mu$ . Photons from a *narrow* beam that are absorbed or scattered by the absorber do not reach a small detector placed in beam line some distance away.

• linear attenuation coefficient  $\mu = \mu_{\rm ph} + \mu_{\rm Co} + \mu_{\rm pp}$ 

• mass attenuation coefficient						
$\mu$ _	$\mu_{\rm ph}$	$\mu_{\rm Co}$	$\mu_{ m pp}$			
$\rho$ –	-	ho	0			

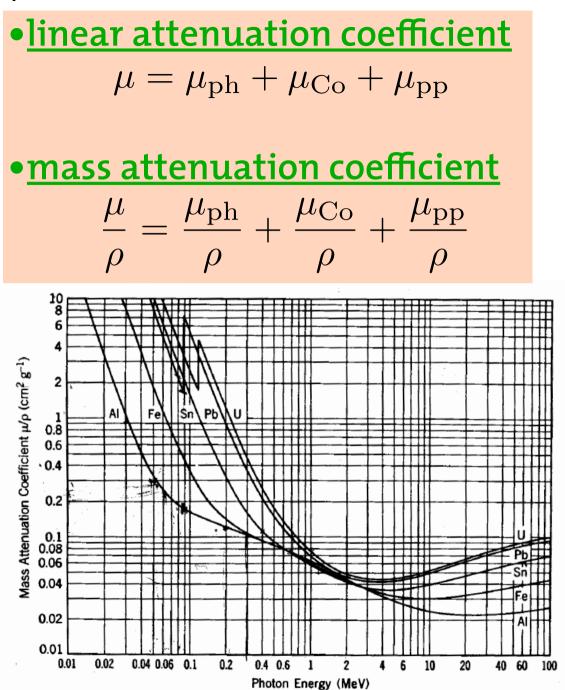
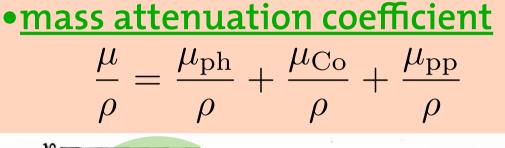


FIGURE 8.8. Mass attenuation coefficients for various elements. [Reprinted with permission from K. Z. Morgan and J. E. Turner, eds., *Principles of Radiation Protection*, Wiley, New York (1967). Copyright 1967 by John Wiley & Sons.]

• linear attenuation coefficient  $\mu = \mu_{\rm ph} + \mu_{\rm Co} + \mu_{\rm pp}$ 



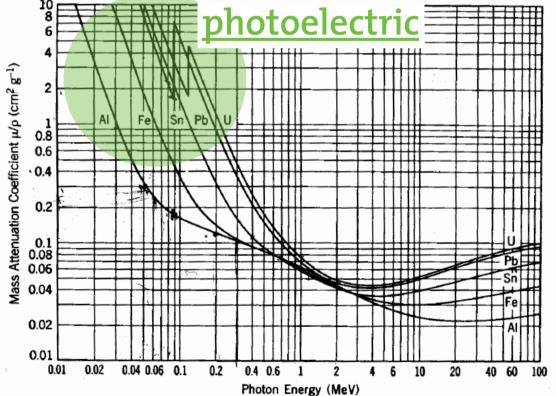


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• linear attenuation coefficient  $\mu = \mu_{\rm ph} + \mu_{\rm Co} + \mu_{\rm pp}$ 

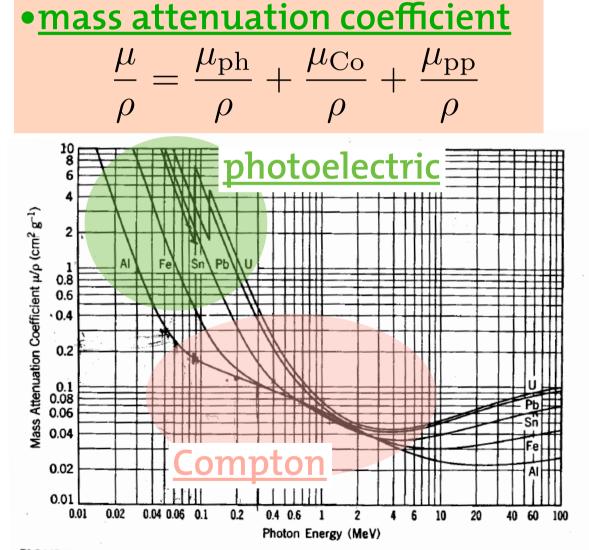
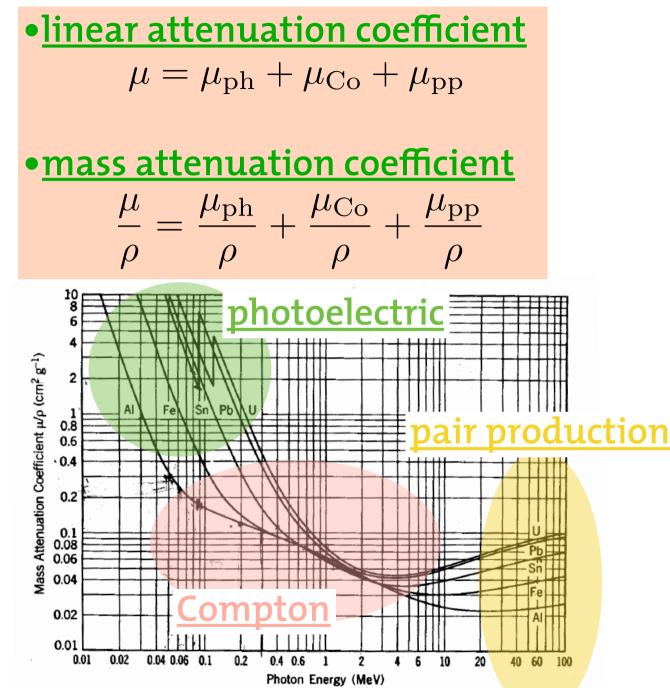
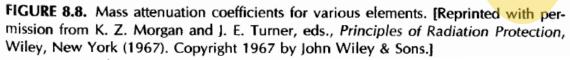
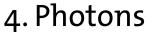


FIGURE 8.8. Mass attenuation coefficients for various elements. [Reprinted with permission from K. Z. Morgan and J. E. Turner, eds., *Principles of Radiation Protection*, Wiley, New York (1967). Copyright 1967 by John Wiley & Sons.]







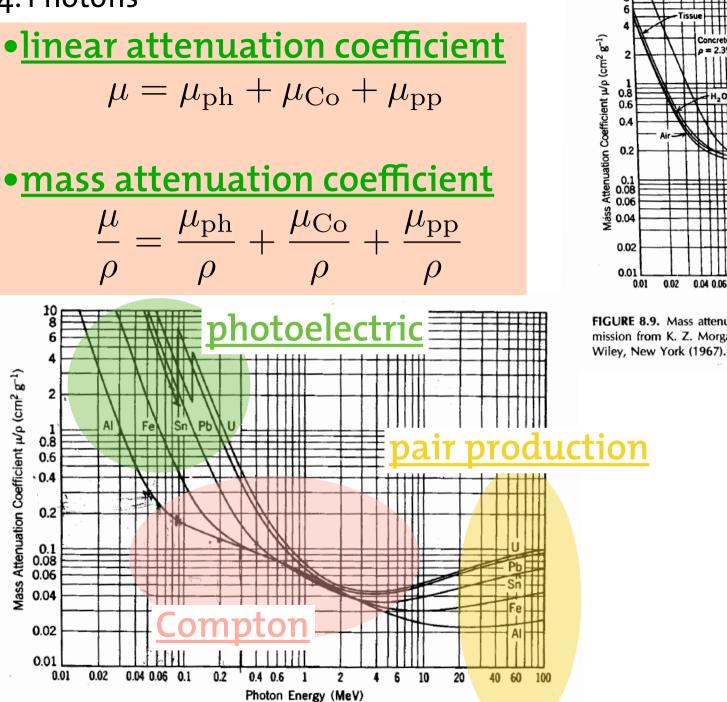


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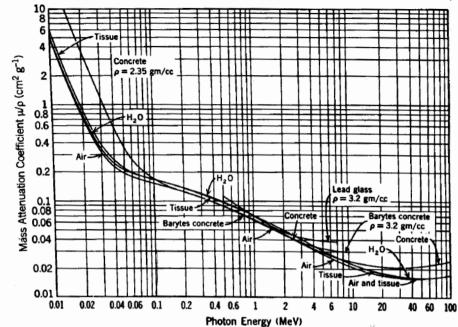


FIGURE 8.9. Mass attenuation coefficients for various materials. [Reprinted with permission from K. Z. Morgan and J. E. Turner, eds., *Principles of Radiation Protection*, Wiley, New York (1967). Copyright 1967 by John Wiley & Sons.]

## • energy transfer coefficients

 not all energy delivered by the incident photons is transferred to electrons (which are those actually depositing the energy in the material)

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-pair production: electron and positron

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$$\frac{\mu^{\rm tr}}{\rho} = \frac{\mu^{\rm tr}_{\rm ph}}{\rho} + \frac{\mu^{\rm tr}_{\rm Co}}{\rho} + \frac{\mu^{\rm tr}_{\rm pp}}{\rho} = \frac{\mu_{\rm ph}}{\rho} \left(1 - \frac{\delta}{h\nu}\right) + \frac{\mu_{\rm Co}}{\rho} \frac{T_{\rm avg}}{h\nu} + \frac{\mu_{\rm pp}}{\rho} \left(1 - \frac{2mc^2}{h\nu}\right)$$

## energy absorption coefficients

- electrons may produce Bremsstrahlung photons

$$\frac{\mu^{\rm en}}{\rho} = \frac{\mu^{\rm tr}}{\rho} (1-g)$$

4. Photons

### energy absorption coefficients

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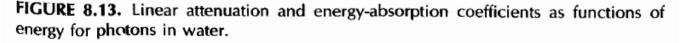
$$\frac{\mu^{\rm en}}{\rho} = \frac{\mu^{\rm tr}}{\rho} (1-g)$$

10 COEFFICIENT (cm<sup>-1</sup>) 0.1 uen 0.01 0.001 0.1 0.01 10 1 100 ENERGY (MeV)

TABLE 8.3. Mass Attenuation, Mass Energy-Transfer, and Mass Energy-Absorption Coefficients ( $\text{cm}^2 \text{ g}^{-1}$ ) for Photons in Water and Lead

Photon Energy (MeV)	Water			Lead		
	μ/ρ	$\mu_{ m tr}/ ho$	$\mu_{\rm en}/ ho$	μίρ	$\mu_{tr}/\rho$	$\mu_{\rm en}/\rho$
0.01	5.33	4.95	4.95	131.	126.	126.
0.10	0.171	0.0255	0.0255	5.55	2.16	2.16
1.0	0.0708	0.0311	0.0310	0.0710	0.0389	0.0379
10.0	0.0222	0.0163	0.0157	0.0497	0.0418	0.0325
100.0	0.0173	0.0167	0.0122	0.0931	0.0918	0.0323

Source: Based on P. D. Higgins, F. H. Attix, J. H. Hubbell, S. M. Seltzer, M. J. Berger, and C. H. Sibata, Mass Energy-Transfer and Mass Energy-Absorption Coefficients, Including In-Flight Positron Annihilation for Photon Energies 1 keV to 100 MeV, NISTIR 4680, National Institute of Standards and Technology, Gaithersburg, MD (1991).



nuclear reactions: production and interaction

- cannot be accelerated !

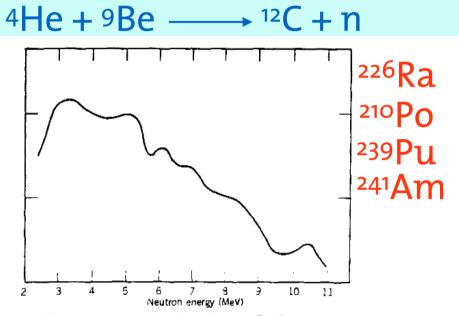
• produced at a high energy and then <u>moderated</u>

5. Neutrons	type	energy	
nuclear reactions: production and interaction	Thermal Epithermal	~0.025 eV ~1 eV	
- cannot be accelerated !	Slow Fast	~1 keV >100 keV	

• produced at a high energy and then <u>moderated</u>



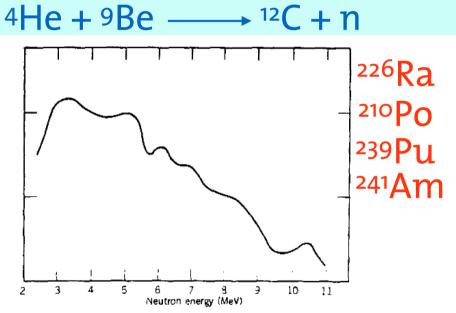
produced at a high energy and then moderated



**Figure 12.1** Neutron energy spectrum from a Ra-Be source, measured with a proton recoil counter. Several neutron groups are present; they result from reactions induced by  $\alpha$ 's with differing energies and in which the <sup>12</sup>C is left either in the ground state or the 4.43- or 7.6-MeV excited states.

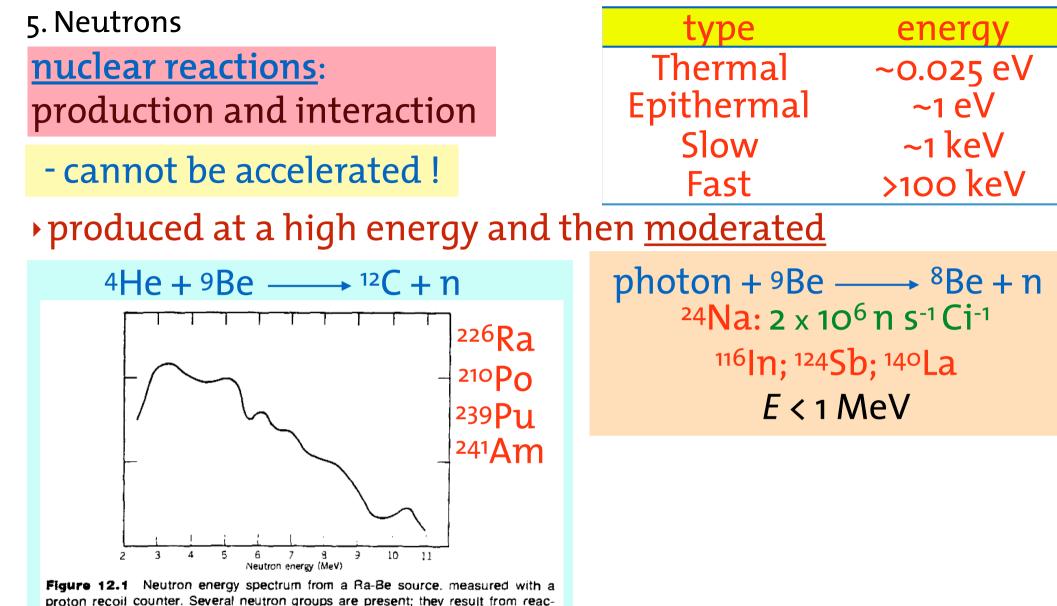
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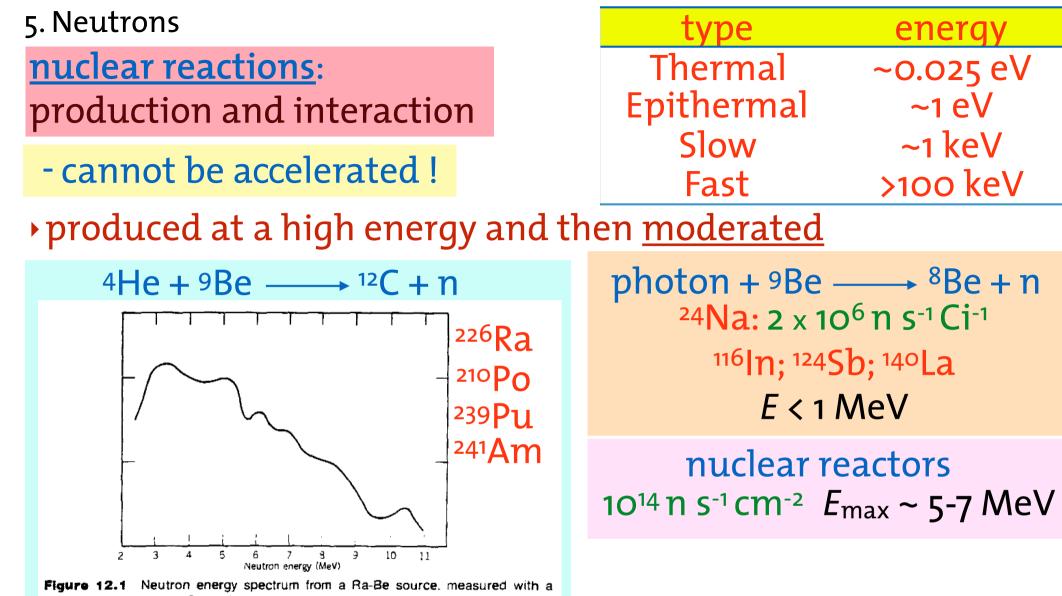
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**spontaneous fission 252Cf;** *E* ~ 1-3 MeV; 4.3 × 10<sup>9</sup> n s<sup>-1</sup> Ci<sup>-1</sup>



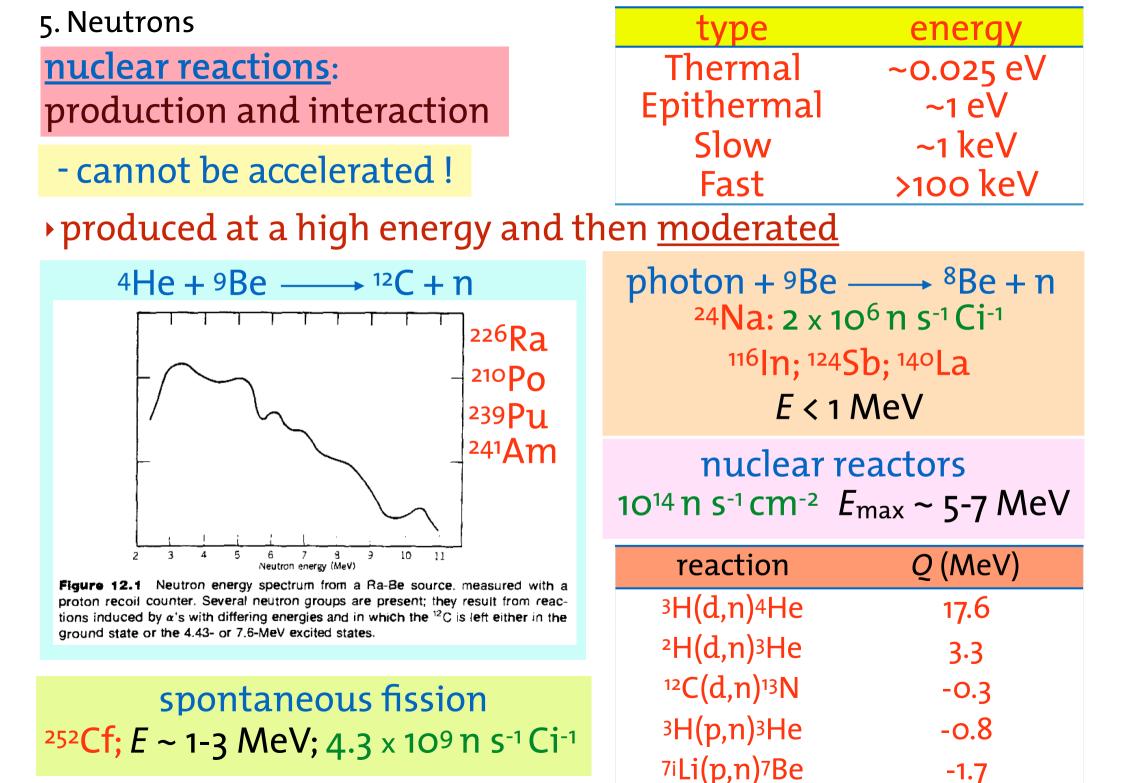
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- minimal interaction with electrons
- travel large distances (even thermal neutrons)

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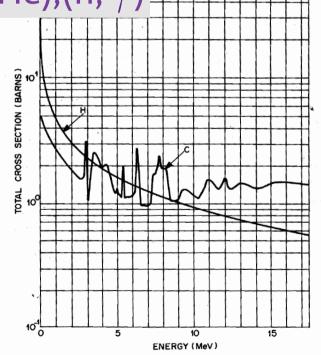
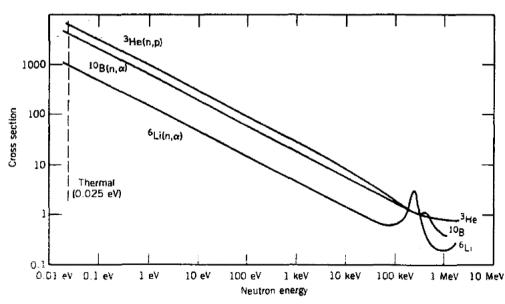


FIGURE 9.2. Total cross sections for neutrons with hydrogen and carbon as functions of energy.

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**Figure 12.5** Neutron cross sections for  ${}^{3}\text{He}(n,p)$ ,  ${}^{10}\text{B}(n,\alpha)$ , and  ${}^{6}\text{Li}(n,\alpha)$ . The cross section shows the 1/v behavior for E < 1 keV, but begins to show resonances above 100 keV.

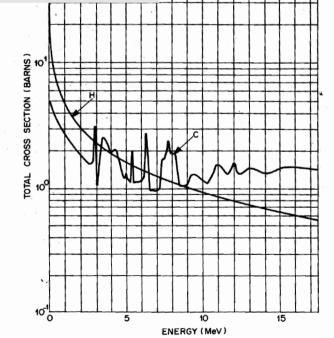


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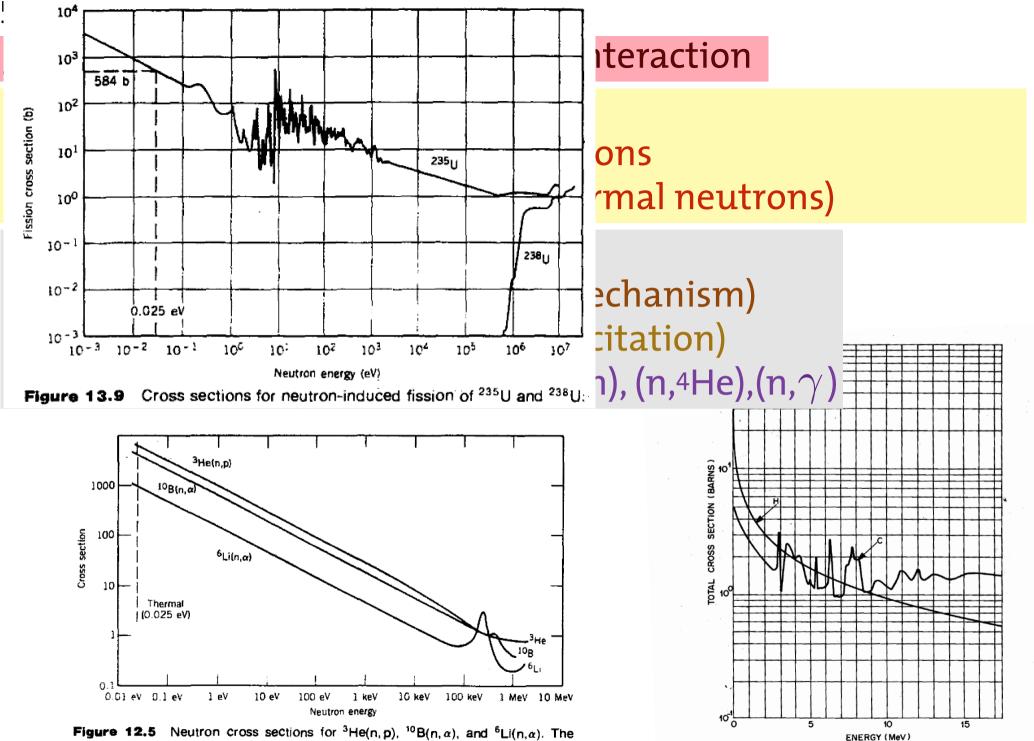


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## 6. Simulating particle transport in matter

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 radiation transport in materials: topic of interest in numerous fields 6. Simulating particle transport in matter

• radiation transport in materials:

topic of interest in numerous fields

- electronic and positronic surface spectrosopy
- electronic microscopy
- microanalysis with electronic probe
- design and use of radiation detectors
- dosimetry
- radiotherapy
- ...

6. Simulating particle transport in matter
radiation transport in materials: topic of interest in numerous fields
high energy particles in materials: suffer many interactions, transfer energy and produce secondary particles 6. Simulating particle transport in matter

radiation transport in materials:

topic of interest in numerous fields

 high energy particles in materials: suffer many interactions, transfer energy and produce secondary particles

- formation of particle showers

6. Simulating particle transport in matter
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•radiation transport in materials:

solve the Boltzmann equation of transport

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 radiation transport in materials: solve the Boltzmann equation of transport

- difficult in limited geometries
- aplied to infinite or semi-infinite means

6. Simulating particle transport in matter radiation transport in materials: topic of interest in numerous fields • high energy particles in materials: suffer many interactions, transfer energy and produce secondary particles • radiation transport in materials: solve the Boltzmann equation of transport transport proccesses: intrinsically random

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parallel to computer development

6. Simulating particle transport in matter radiation transport in materials: topic of interest in numerous fields • high energy particles in materials: suffer many interactions, transfer energy and produce secondary particles • radiation transport in materials: solve the Boltzmann equation of transport •transport proccesses: intrinsically random Monte Carlo simulation: powerful alternative to solve transport problems •increasing use: parallel to computer development • first MC simulation of photon transport:

Hayward and Hubbell (1954): 67 histories

## 6. Simulating particle transport in matter

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•<u>detailed simulation</u>: is exact!

6. Simulating particle transport in matter<a href="https://detailed.simulation">detailed simulation</a>: is exact!

simulation of all the interaction events experienced by a particle in cronological sequence 6. Simulating particle transport in matter•detailed simulation: is exact!

detailed simulation: is exact!
 detailed simulation of photons:
 relatively simple
 in cronological sequence

simulation of all the

-photons interact just a few times before being absorbed

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 complicated due to the continuous interaction
 small energy losses in each interaction
 many interactions before being absorbed
 only possible for low energy particles

or thin geometries

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 complicated due to the continuous interaction

- small energy losses in each interaction
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• solving the problem:

multiple scattering theories

- -the effect of a large number of interaction events is simulated globally
- approximate theories: possible problems
- details depend on the particular code

## 6. Simulating particle transport in matter

6. Simulating particle transport in matter

**Scattering model and probability distributions** 

Scattering model and probability distributions

•simulate the particle transport in material media

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•simulate the particle transport in material media

homogeneous media with uniform density

- gases, liquids or amorphous solids

- random scattering of particles

Scattering model and probability distributions

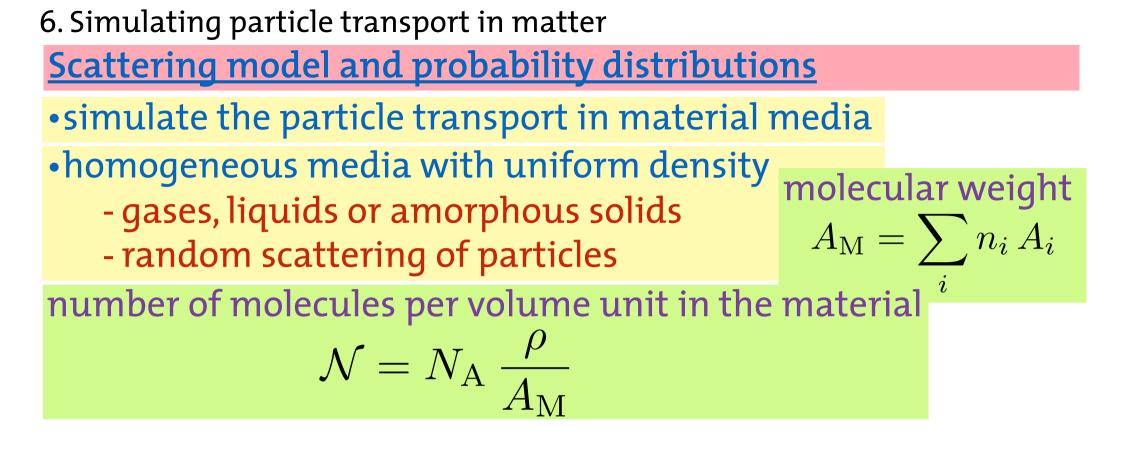
•simulate the particle transport in material media

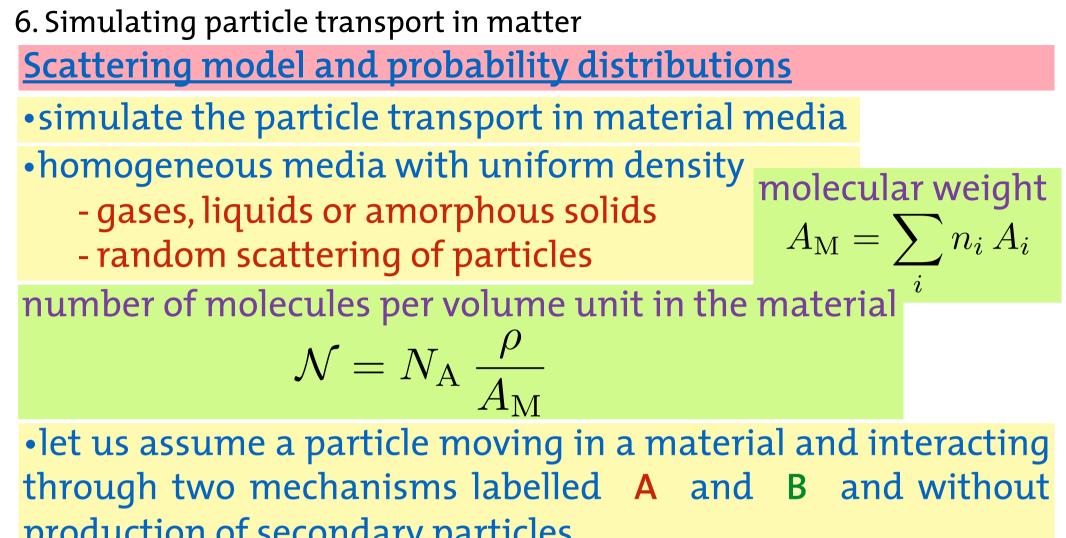
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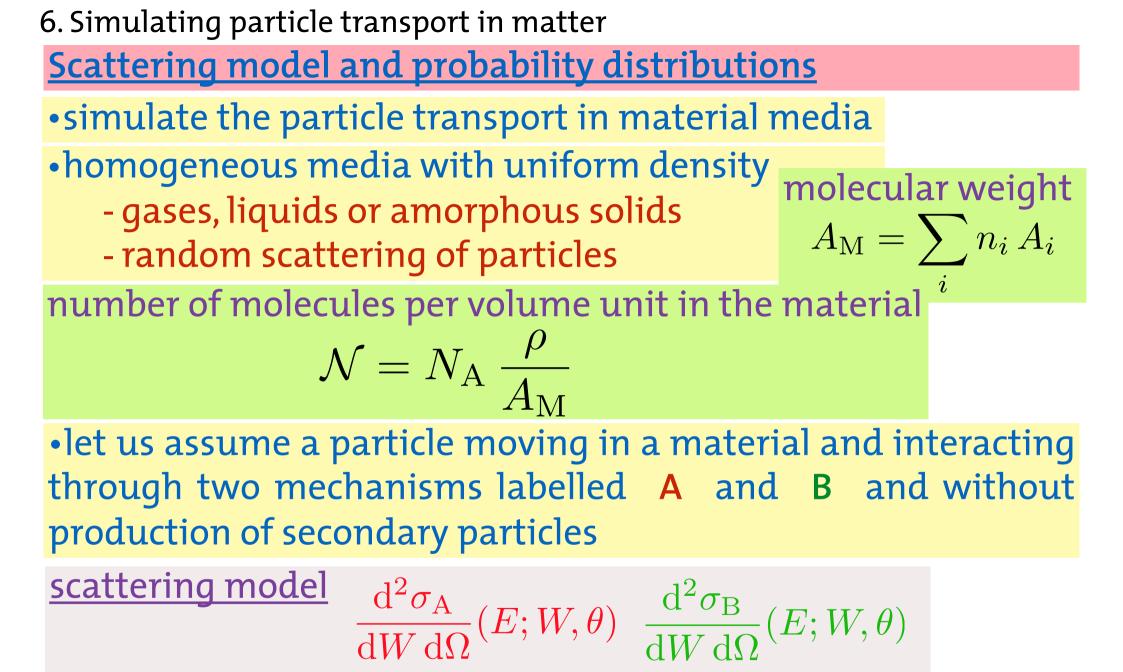
number of molecules per volume unit in the material

$$\mathcal{N} = N_{\rm A} \, \frac{\rho}{A_{\rm M}}$$





production of secondary particles



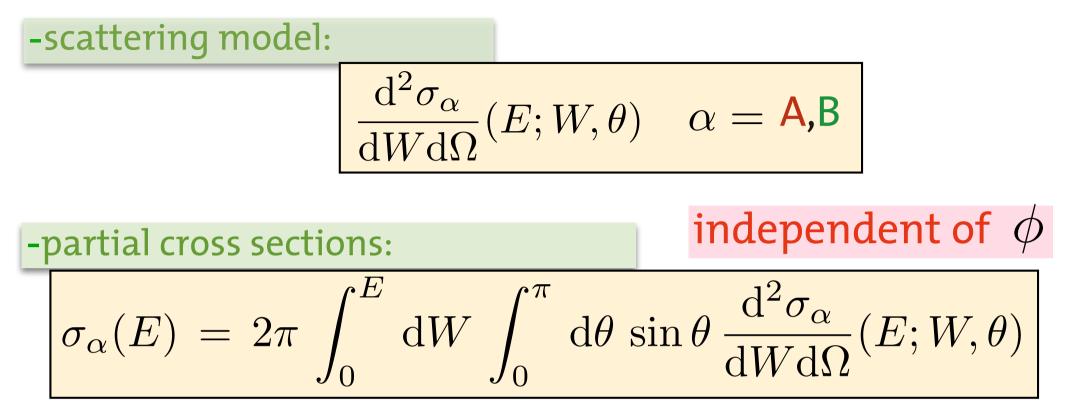
 $W \equiv$  energy lost after an interaction  $\Omega \equiv$  solid angle in the new direction

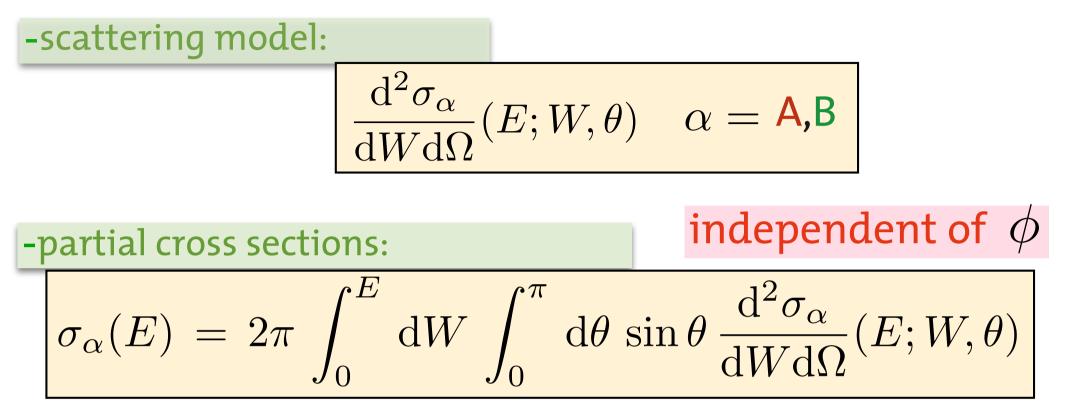
6. Simulating particle transport in matter

-scattering model:  $\frac{\mathrm{d}^2\sigma_\alpha}{\mathrm{d}W\mathrm{d}\Omega}(E;W,\theta) \quad \alpha=\mathsf{A},\mathsf{B}$ 

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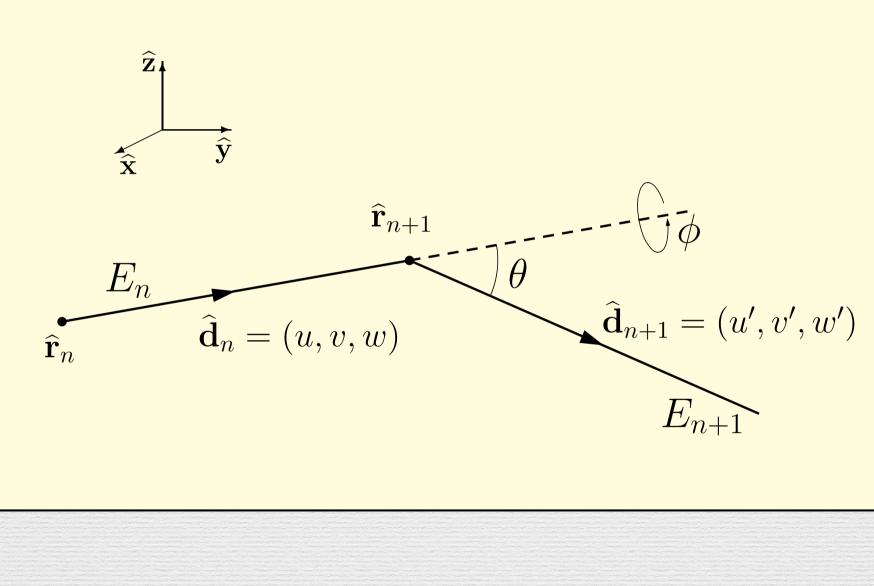
# -partial cross sections: $\sigma_{\alpha}(E) = 2\pi \int_{0}^{E} dW \int_{0}^{\pi} d\theta \sin \theta \frac{d^{2}\sigma_{\alpha}}{dWd\Omega}(E; W, \theta)$





-total scattering cross section:

$$\sigma_{\rm T}(E) = \sum_{\alpha \equiv \mathbf{A}, \mathbf{B}} \sigma_{\alpha}(E)$$



-probability density function of the track length followed by the particle between its actual position and that of the next interaction: f(s)

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-probability that the particle follows a track of length S without interacting  $\int_{-\infty}^{\infty}$ 

$$\mathcal{F}(s) = \int_{s}^{\infty} \mathrm{d}\xi f(\xi)$$

-probability density function of the track length followed by the particle between its actual position and that of the next interaction: f(s)

-probability that the particle follows a track of length *S* without interacting  $\int_{-\infty}^{\infty} |f_{n}f(t)|^{2}$ 

$$\mathcal{F}(s) = \int_{s} d\xi f(\xi)$$

-probability that the particle interacts after going through a distance in the range (s, s + ds)

 $\mathcal{F}(s)\mathcal{N}\sigma_{\mathrm{T}}\mathrm{d}s = f(s)\,\mathrm{d}s$ 

-probability density function of the track length followed by the particle between its actual position and that of the next interaction: f(s)

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-probability that the particle interacts after going through a distance in the range (s, s + ds) $\mathcal{F}(s)\mathcal{N}\sigma_{\mathrm{T}}ds = f(s) ds$ 

$$f(s) = \mathcal{N}\sigma_{\mathrm{T}}\mathcal{F}(s) = \frac{1}{\lambda_{\mathrm{T}}} \int_{s}^{\infty} \mathrm{d}\xi f(\xi) \qquad f(\infty) = 0$$

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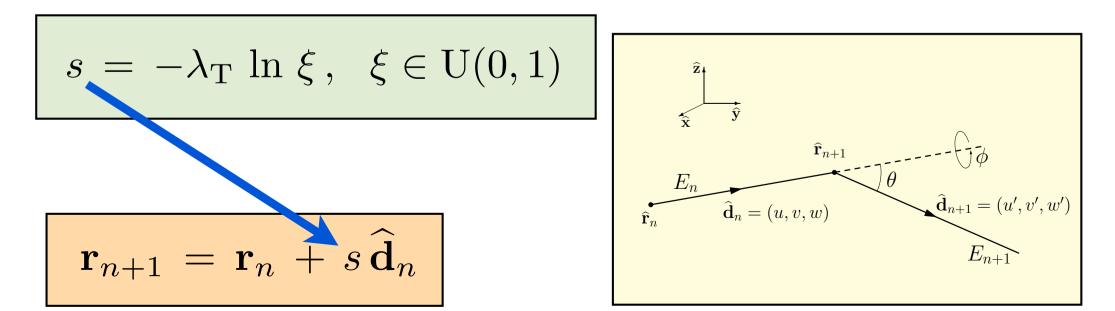
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$$\mathbf{r}_{n+1} = \mathbf{r}_n + s \,\widehat{\mathbf{d}}_n$$

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-which type of interaction? A o B?

6. Simulating particle transport in matter

$$\mathbf{r}_{n+1} = \mathbf{r}_n + s \,\widehat{\mathbf{d}}_n$$

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-determine W,  $\theta$  and  $\phi$  $p_{\alpha}(E; W, \theta) = 2\pi \sin \theta \frac{1}{\sigma_{\alpha}(E)} \frac{d^{2}\sigma_{\alpha}}{dWd\Omega}(E; W, \theta)$   $\phi \in U[0, 2\pi]$   $\sigma_{\alpha}(E) = 2\pi \int_{0}^{E} dW \int_{0}^{\pi} d\theta \sin \theta \frac{d^{2}\sigma_{\alpha}}{dWd\Omega}(E; W, \theta)$ 

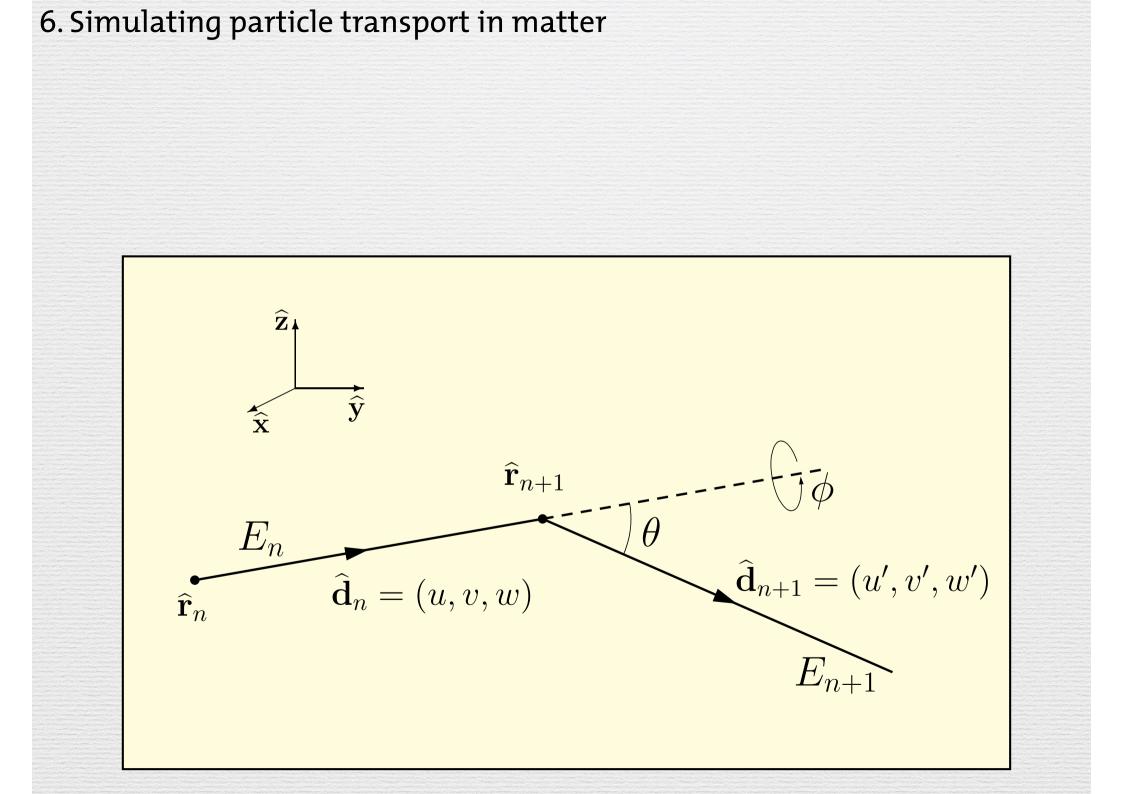
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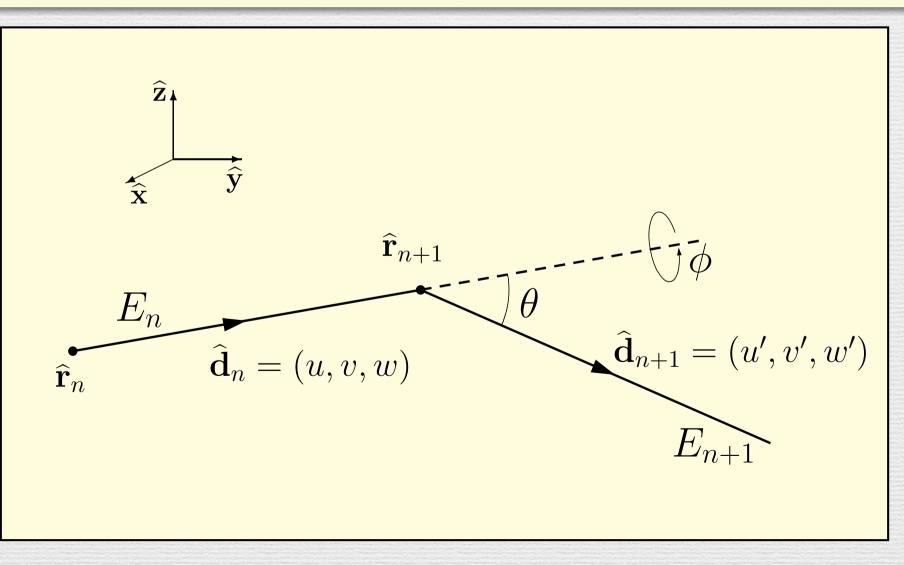
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- •the algorithm is iteratively repeated until
  - -the particle energy is below the absorption energy (fixed by the user)
  - -the particle moves out of the material medium



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# • problems to be solved:

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# • problems to be solved:

# -interface crossing

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-evaluation of statistical estimators

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# • problems to be solved:

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-variance reduction techniques